

**OPTICAL PROPERTIES OF TURBID COASTAL  
WATER AND THEIR REMOTE SENSING  
RETRIEVALS BASED ON RADIATIVE  
TRANSFER APPROACHES**

*Xiaolong Yu*



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TRANSFER APPROACHES**

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## Summary

Radiative transfer theory describes how light, regarding its spectral radiance, travels and varies along any specific path at a specified point in the water, and ties together the water apparent optical properties (AOPs) and inherent optical properties (IOPs) through the radiative transfer equations (RTEs). In this study, the author used the analytical 2SeaColor model, which is based on the approximate solution of RTEs using two-stream approach, to retrieve IOPs, diffuse attenuation coefficient ( $K_d$ , in  $\text{m}^{-1}$ ), and the concentration of suspended particulate matter ( $C_{\text{SPM}}$ , in  $\text{mg/L}$ ) in the turbid Yangtze estuary. The application of 2SeaColor to the GOCI geostationary sensors has enabled a better understanding of the diurnal cycle of suspended sediment in the study area. The known linear relationship between  $C_{\text{SPM}}$  and particulate backscattering coefficient ( $b_{bp}$ , in  $\text{m}^{-1}$ ), was investigated, using 2SeaColor retrievals and field measurements. Although the validity of the linear relationship was confirmed for relatively clear waters, the author showed that in turbid waters, nonlinearity governs the relationship between  $C_{\text{SPM}}$  and  $b_{bp}$ . The two respective relationships between  $C_{\text{SPM}}$  and  $b_{bp}$  for relatively clear and turbid waters were further integrated using a sigmoid function to retrieve wide-range  $C_{\text{SPM}}$  without the need for a switching scheme. Apart from SPM, colored dissolved organic matter (CDOM) also significantly reduces light transparency and contributes to the complexity of the optical properties in the Yangtze estuary. Thereby, the author has also investigated the seasonal variability of CDOM absorption properties. Finally, the author proposed a locally calibrated model to estimate the concentration of dissolved organic carbon (DOC). The summary of the major results obtained from this dissertation is briefly described in the following.

The thesis is structured in six chapters, of which the first chapter describes the background of ocean color remote sensing in optically complex waters and the second chapter explains the physical principals of the 2SeaColor model. In Chapter 3, the author first proposed improved IOPs parameterizations to the analytical 2SeaColor model to account for the complex optical properties in the Yangtze estuary. The improved 2SeaColor model was later employed to derive the diffuse attenuation coefficient ( $K_d$ , in  $\text{m}^{-1}$ ) from remote sensing reflectance ( $R_{rs}$ , in  $\text{sr}^{-1}$ ). Retrievals from the improved model were then validated with three *in-situ* datasets and compared with an empirical (the Zhang model) and a semi-analytical model (the Lee model). Statistics from the validation results show that the 2SeaColor model provides the best estimates of  $K_d$  for the full range of

observations, with the largest determination coefficient ( $R^2 = 0.935$ ) and the smallest root mean squared error (RMSE =  $0.078 \text{ m}^{-1}$ ). For clear waters, where  $K_d(490) < 0.2 \text{ m}^{-1}$ , the Zhang model provides the most accurate  $K_d$  estimations, but results from the Lee model and the 2SeaColor are comparable. For turbid waters, where  $K_d(490) > 0.2 \text{ m}^{-1}$ , the 2SeaColor model is found to be more accurate, with an RMSE of  $0.186 \text{ m}^{-1}$ , compared to RMSEs of  $0.279 \text{ m}^{-1}$  and  $0.388 \text{ m}^{-1}$  for the Zhang model and the Lee model, respectively. The improved 2SeaColor model is finally applied to the GOCI (Geostationary Ocean Color Imager) level 2 product to produce  $K_d$  maps over the Yangtze estuary, resulting in a reasonable distribution and expected range of  $K_d$  for the Yangtze estuary, while the Zhang model and the Lee model show the potential of underestimation and overestimation, respectively. The consistently stable and accurate  $K_d$  estimates in both clear and turbid waters show great potentials of the 2SeaColor model for estimating  $K_d$  over optically complex waters.

In Chapter 4, the author proposed a novel model to quantitatively estimate the concentration of SPM ( $C_{\text{SPM}}$ , in mg/L) from remotely sensed particulate backscattering ( $b_{\text{bp}}$ , in  $\text{m}^{-1}$ ).  $b_{\text{bp}}$  of both clear and turbid waters were first retrieved by the improved 2SeaColor model, and then the relationships between  $b_{\text{bp}}$  and  $C_{\text{SPM}}$  were calibrated for a wide-range of  $C_{\text{SPM}}$  from 0.4 to 2068.8 mg/L. Calibration results show that  $b_{\text{bp}}$  and  $C_{\text{SPM}}$  are linearly correlated for relatively clear waters, but their relationship becomes significantly curved for turbid water. To avoid the discontinuity issue resulted from the two respective  $b_{\text{bp}}-C_{\text{SPM}}$  relationships for clear and turbid waters, an index of  $b_{\text{bp}}$  ( $S_{\text{index}}$ ) was introduced based on a sigmoid function to retrieve  $C_{\text{SPM}}$ . The  $S_{\text{index}}-C_{\text{SPM}}$  model was calibrated using a calibration (Cal) dataset and then validated with the validation (Val) dataset. Validation results show that the  $S_{\text{index}}-C_{\text{SPM}}$  model has overall better retrieval accuracy (rMAD = 33.4%) compared with the re-calibrated  $He$  model (rMAD = 185.2%) and the GOCI SPM model (rMAD = 50.2%). For extremely turbid waters, these three models show comparable performance with rMAD of 15.29%, 12.72%, and 17.22% for the  $S_{\text{index}}-C_{\text{SPM}}$  model, the GOCI SPM model, and the  $He$  model, respectively. The  $S_{\text{index}}-C_{\text{SPM}}$  model, without tuning the empirical constants, was also applied to an independent SeaSWIR dataset collected from the estuaries of Europe and Argentina with measured  $C_{\text{SPM}}$  up to 1400 mg/L. Validation results show a very promising retrieval accuracy with rMAD of 30.3 % without tuning the empirical coefficients, indicating the potential good transferability of the  $S_{\text{index}}-C_{\text{SPM}}$  model among global turbid coastal waters. Further, the three models were implemented on the GOCI image acquired

at 02:16 UTC, March 8<sup>th</sup> of 2013 to generate the SPM mapping products for the Yangtze estuary, where discontinuities are observed for the switching-based GOCI SPM model at the front of the turbid plume, while the  $S_{index}$ - $C_{SPM}$  model produced a smooth and realistic SPM mapping product. At last, the  $S_{index}$ - $C_{SPM}$  model was applied to 8 GOCI images acquired on March 8<sup>th</sup> of 2013 to investigate the diurnal variation of  $C_{SPM}$ , and the comparison with tide information shows that the predicted  $C_{SPM}$  can be well explained by hydrologic processes, which further enhanced our confidence in the reliability and repeatability of the  $S_{index}$ - $C_{SPM}$  model.

In Chapter 5, the author investigated the seasonal variability of CDOM absorption properties and its relationship between salinity and dissolved organic carbon (DOC). The CDOM absorption coefficient at 355 nm ( $a_g(355)$ ) was found to be inversely correlated with salinity, with Pearson's coefficients  $r$  of -0.901 and -0.826 for summer and winter observations, respectively. Analysis results of the relationships between salinity and CDOM optical properties (*i.e.*, absorption coefficient and spectral slope) suggested that terrigenous inputs dominated CDOM sources in the Yangtze estuary, but the proportion of terrigenous CDOM declined with increasing salinity. Resuspension of bottom sediments could be an important source of CDOM in the winter. We further developed an empirical model to estimate DOC concentration from  $a_g(355)$  and spectral slope  $S_{275-295}$  using the calibration dataset by non-linear regression. The developed model was validated with the validation dataset, resulting in an acceptable error with a determination coefficient ( $R^2$ ) of 0.746, a root mean squared error (RMSE) of 20.99  $\mu\text{mol/L}$  and a mean relative absolute error (rMAD) of 14.46%. This empirical model could be potentially adopted in retrieving DOC concentration using ocean color remote sensing over the Yangtze estuarine and coastal waters.

## *Summary*

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## Samenvatting

De theorie van stralingsoverdracht beschrijft hoe licht, in eenheden van spectrale radiantie, zich voortplant and varieert over elke specifieke weg in het water, en verbindt de schijnbare optische eigenschappen (apparent optical properties of AOP's) van water met de inherente optische eigenschappen (inherent optical properties of IOP's) via de vergelijkingen voor stralingstransport. In deze studie heeft de auteur het analytische model 2SeaColor gebruikt, dat is gebaseerd op een vereenvoudigde oplossing van de stralingstransportvergelijkingen met de tweestromen benadering, om IOP's and concentraties van opgeloste deeltjes ( $C_{SPM}$ , in mg/L) te bepalen in het zeer troebele Yangtze estuarium. De toepassing van 2SeaColor op de sensoren van de GOCI geostationaire satelliet heeft een beter begrip mogelijk gemaakt van de dagelijkse cyclus in het zwevend sediment in het studiegebied. Het bekende lineaire verband tussen  $C_{SPM}$  en de deeltjes backscattercoëfficiënt ( $b_{bp}$ , in  $m^{-1}$ ) is onderzocht door middel van 2SeaColor bepalingen en veldmetingen. Hoewel de geldigheid van dit lineaire verband werd bevestigd voor relatief heldere wateren, toont de auteur aan dat in troebel water het verband gekenmerkt wordt door niet-lineariteit. De twee respectievelijke verbanden tussen  $C_{SPM}$  en  $b_{bp}$  voor relatief heldere en troebele wateren zijn verder geïntegreerd met een sigmoïde functie om  $C_{SPM}$  over een breed bereik te kunnen bepalen zonder gebruik te maken van een schakelmechanisme. Naast SPM, reduceren ook gekleurde opgeloste stoffen (colored dissolved organic matter, of CDOM) de transparantie voor licht en dragen deze bij aan de complexiteit van de optische eigenschappen in het Yangtze estuarium. Daarom heeft de auteur ook de seizoensafhankelijke variabiliteit in absorptie-eigenschappen van deze stoffen in het Yangtze estuarium onderzocht. Tenslotte formuleerde de auteur een lokaal geijkt model om de concentratie van opgeloste organische koolstof (dissolved organic carbon, of DOC) te schatten. Een overzicht van de voornaamste resultaten verkregen in deze dissertatie wordt hieronder beschreven.

De dissertatie is gestructureerd in zes hoofdstukken, waarvan het eerste hoofdstuk de achtergronden beschrijft van ocean color remote sensing in optisch complexe wateren en het tweede hoofdstuk de fysische principes verklaart van het 2SeaColor model. In Hoofdstuk 3 stelt de auteur eerst verbeterde IOP parametrisaties voor in het analytische 2SeaColor model teneinde rekening te kunnen houden met de complexe optische eigenschappen in het Yangtze estuarium. Het verbeterde 2SeaColor model werd later ingezet om de diffuse uitdovingscoëfficiënt ( $K_d$ , in  $m^{-1}$ ) af te leiden uit de remote sensing reflectie ( $R_{rs}$ ,

in  $\text{sr}^{-1}$ ). Bepalingen met het verbeterde model zijn vervolgens gevalideerd met drie in-situ datasets en vergeleken met een empirisch model (het Zhang model) en met een semi-analytisch model (het Lee model). De statistiek van de validatieresultaten laat zien dat het 2SeaColor model de beste schattingen geeft van  $K_d$  over het gehele bereik van de waarnemingen, met de hoogste determinatiecoëfficiënt ( $R^2 = 0.935$ ) en de kleinste root mean squared error (RMSE =  $0.078 \text{ m}^{-1}$ ). Voor heldere wateren, met  $K_d(490) < 0.2 \text{ m}^{-1}$ , geeft het Zhang model de nauwkeurigste  $K_d$  schattingen, maar de resultaten van het Lee model en 2SeaColor zijn redelijk vergelijkbaar. Voor troebel water, waar  $K_d(490) > 0.2 \text{ m}^{-1}$ , werd gevonden dat 2SeaColor nauwkeuriger is, met een RMSE van  $0.186 \text{ m}^{-1}$ , vergeleken met RMSE's van  $0.279 \text{ m}^{-1}$  en  $0.388 \text{ m}^{-1}$  voor respectievelijk het Zhang model en het Lee model. Het verbeterde 2SeaColor model is uiteindelijk toegepast op het GOCI (Geostationary Ocean Color Imager) level 2 product om  $K_d$  kaarten te maken van het Yangtze estuarium, wat resulteerde in een redelijke verdeling en een verwacht bereik van  $K_d$  voor dit gebied, terwijl het Zhang model en het Lee model respectievelijk onderschatting en overschatting kunnen laten zien. De consequent stabiele en nauwkeurige  $K_d$  schattingen in zowel helder als troebel water tonen een groot potentieel van het 2SeaColor model voor het schatten van  $K_d$  in optisch complexe wateren.

In Hoofdstuk 4 introduceert de auteur een nieuw model voor het kwantitatief schatten van de concentratie van SPM ( $C_{\text{SPM}}$ , in  $\text{mg/L}$ ) uit de met remote sensing bepaalde deeltjes backscatter coëfficiënt ( $b_{bp}$ , in  $\text{m}^{-1}$ ). De  $b_{bp}$  van zowel heldere als troebele wateren zijn eerst bepaald met het verbeterde 2SeaColor model, en vervolgens zijn de verbanden tussen  $b_{bp}$  en  $C_{\text{SPM}}$  geijkt over een breed bereik van  $C_{\text{SPM}}$  van 0.4 tot  $2068.8 \text{ mg/L}$ . Calibratieresultaten laten zien dat  $b_{bp}$  en  $C_{\text{SPM}}$  lineair zijn gecorreleerd voor relatief heldere wateren, maar het verband wordt sterk kromlijng voor troebel water. Om discontinuïteit als gevolg van de twee respectievelijke  $b_{bp}$ - $C_{\text{SPM}}$  verbanden voor helder en troebel water te vermijden is een index van  $b_{bp}$  ( $S_{\text{index}}$ ) geïntroduceerd die is gebaseerd op een sigmoïdefunctie om  $C_{\text{SPM}}$  te bepalen. Het  $S_{\text{index}}$ - $C_{\text{SPM}}$  model is gecalibreerd met een Cal dataset en vervolgens gevalideerd met een Val dataset. Validatieresultaten laten zien dat het  $S_{\text{index}}$ - $C_{\text{SPM}}$  model over het geheel een betere schattingsnauwkeurigheid heeft (relative mean absolute deviation rMAD = 33.4%) vergeleken met het gehercalibreerde He model (rMAD = 185.2%) en het GOCI SPM model (rMAD = 50.2%). Voor extreem troebele wateren laten deze drie modellen vergelijkbare prestaties zien, met respectievelijk een rMAD van 15.29%, 12.72%, en 17.22% voor het  $S_{\text{index}}$ - $C_{\text{SPM}}$  model, het GOCI SPM model, en het He model. Het  $S_{\text{index}}$ -



$C_{SPM}$  model, zonder aanpassingen aan de empirische coëfficiënten, is ook toegepast op een totaal onafhankelijke SeaSWIR dataset verzameld in estuaria in Europa en Argentinië met gemeten  $C_{SPM}$  tot 1400 mg/L. De validatieresultaten tonen een veelbelovende schattingsnauwkeurigheid met een rMAD van 30.3 % zonder bijstelling van de empirische coëfficiënten, wat wijst op een goede overdraagbaarheid van het  $S_{index}$ - $C_{SPM}$  model naar de troebele kustwateren wereldwijd. Verder zijn de drie modellen toegepast op een GOCI beeld opgenomen om 2:26 uur UTC op 8 maart 2013 om SPM kaartproducten te genereren voor het Yangtze estuarium, waar discontinuïteiten worden waargenomen met het geschakelde GOCI SPM model bij het front van een troebele pluim, terwijl het  $S_{index}$ - $C_{SPM}$  model een geleidelijk en realistisch SPM kaartproduct genereerde. Tenslotte is het  $S_{index}$ - $C_{SPM}$  model toegepast op acht GOCI beelden, ook opgenomen op 8 maart 2013, om de dagelijkse variaties in  $C_{SPM}$  te onderzoeken, en de vergelijking met getijde-informatie laat zien dat de voorspelde  $C_{SPM}$  goed kan worden verklaard door hydrologische processen, hetgeen ons vertrouwen in de betrouwbaarheid en de herhaalbaarheid van het  $S_{index}$ - $C_{SPM}$  verder versterkte.

In Hoofdstuk 5 heeft de auteur de seizoensafhankelijke variabiliteit in de absorptie-eigenschappen van CDOM onderzocht en de relaties hiervan met zoutgehalte en opgeloste organische koolstof (DOC). Voor de CDOM absorptiecoëfficiënt bij 355 nm ( $a_g(355)$ ) werd gevonden dat deze negatief gecorreleerd was aan het zoutgehalte, met Pearson coëfficiënten  $r$  van -0.901 en -0.826 voor waarnemingen in respectievelijk de zomer en winter. Resultaten van de analyse van de relaties tussen zoutgehalte en CDOM optische eigenschappen (d.w.z. de absorptiecoëfficiënt en de spectrale helling) suggereren dat CDOM bronnen worden gedomineerd door materiaal van terrigene oorsprong in het Yangtze estuarium, maar het aandeel van terrigeen CDOM nam af met toenemend zoutgehalte. Hersuspensie van bodemsedimenten zou een belangrijke bron van CDOM in de winter kunnen zijn. We hebben verder een empirisch model ontwikkeld om de DOC concentratie te schatten uit  $a_g(355)$  en de spectrale helling  $S_{275-295}$  door gebruik te maken van de calibratie dataset via niet-lineaire regressie. Het ontwikkelde model is gevalideerd met de validatie dataset, wat resulteerde in een acceptabele fout met een determinatiecoëfficiënt ( $R^2$ ) van 0.746, een RMSE van 20.99  $\mu\text{mol/L}$  en een rMAD van 14.46%. Dit empirische model zou potentieel geadopteerd kunnen worden voor het schatten van DOC concentraties met ocean color remote sensing boven het Yangtze estuarium en de nabije kustwateren.



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“All endings are new beginnings.”

At the time I complete this thesis, I do realize it is just the beginning of a new chapter of my life and career. There will undoubtedly be more challenges ahead, but thanks to my families, supervisors, and friends, I am well prepared now.

*For Mom (ChuanXia Zhang) and Dad (Yu Zhang)*



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## **Lists of abbreviations and symbols**

### *Abbreviations*

AC	Atmospheric correction
AOPs	Apparent Optical Properties
CDM	Combination of CDOM and NAP
CDOM	Colored Dissolved Organic Matter
Chl-a	Chlorophyll a
COASTLOOC	Coastal Surveillance through Observation of Ocean Color
CZCS	Coastal Zone Color Scanner
DOC	Dissolved Organic Carbon
DOM	Dissolved Organic Matter
F100	Percentage of observations with relative error < 100%
F25	Percentage of observations with relative error < 25%
GDPS	GOCI Data Processing System
GOCI	Geostationary Ocean Color Imager
HyperSAS	Hyperspectral Surface Acquisition System
IOCCG	International Ocean Colour Coordinating Group
IOPs	Inherent Optical Properties
K-M	Kubelka-Munk Model
LUTs	Lookup Tables for parameters estimated by MODTRAN
MERIS	Medium Resolution Imaging Spectroradiometer
MODIS	Moderate Resolution Imaging Spectroradiometer
NASA	The National Aeronautics and Space Administration
NAP	Non-algal particles
NIR	Near-infrared band
NN	Neural Network
NOMAD	NASA Bio-Optical Marine Algorithm Data set
OLCI	Ocean Land Colour Instrument
POC	Particulate Organic Carbon
PSU	Practical Salinity Unit
QIOPs	Quasi-Inherent Optical Properties
QSSA	Quasi-Single-scattering Approximation
rMAD	Relative Mean Absolute Difference
RMSE	Root Mean Squared Error
RTE	Radiative Transfer Equation

SeaWiFS	Sea-Viewing Wide Field of View Sensor
SDGs	Sustainable Development Goals
SPM	Suspended Particulate Matter
SSA	Single-scattering Approximation
SRF	Spectral Response Function
SZA	Solar Zenith Angle
TOA	Top of the atmosphere
UTC	Coordinated Universal Time

### ***Symbols***

$a$	Total absorption coefficient, $\text{m}^{-1}$
$a_{dg}$	Combined absorption coefficients of CDOM and NAP, $\text{m}^{-1}$
$a_g$	Absorption coefficient of CDOM, $\text{m}^{-1}$
$a_g^*(355)$	DOC normalized $a_g$ at 355 nm, $\text{m}^{-1} \mu\text{mol}^{-1} \text{L}$
$a_{NAP}$	Absorption coefficient of NAP, $\text{m}^{-1}$
$a_{NAP}^*$	Specific-absorption coefficient of NAP, $\text{m}^2 \text{mg}^{-1}$
$a_p$	Absorption coefficient of total particles, $\text{m}^{-1}$ ,
$a_\phi$	Absorption coefficient of phytoplankton pigment, $\text{m}^{-1}$
$b$	Total scattering coefficient, $\text{m}^{-1}$
$b_b$	Total backscattering coefficient, $\text{m}^{-1}$
$b_{bp}$	Backscattering coefficient of NAP, $\text{m}^{-1}$
$b_{bw}$	Backscattering coefficient of pure water, $\text{m}^{-1}$
$b_f$	Total forward scattering coefficient, $\text{m}^{-1}$
$b_p^*$	Specific-scattering coefficient of total particles, $\text{m}^2 \text{mg}^{-1}$
$b_w$	Scattering coefficient of pure water, $\text{m}^{-1}$
$c$	Total attenuation coefficient, $\text{m}^{-1}$
$C_{SPM}$	SPM concentration, $\text{mg/L}$
$[DOC]$	The concentration of DOC, $\mu\text{mol/L}$
$E^-$	Diffuse downward irradiance, $\text{W m}^{-2} \text{nm}^{-1}$
$E^+$	Diffuse upward irradiance, $\text{W m}^{-2} \text{nm}^{-1}$
$E_d$	Total downward planar irradiance, $\text{W m}^{-2} \text{nm}^{-1}$
$E_s$	Solar irradiance, $\text{W m}^{-2} \text{nm}^{-1}$
$F_0$	Extraterrestrial solar irradiance, $\text{W m}^{-2} \text{nm}^{-1}$
$J$	Jacobian matrix
$k$	Extinction coefficient, $\text{m}^{-1}$
$K_d$	Diffuse attenuation coefficient, $\text{m}^{-1}$

$L_i$	Sky radiance, $\text{W sr}^{-1} \text{m}^{-2} \text{nm}^{-1}$
$L_t$	Sea-surface radiance, $\text{W sr}^{-1} \text{m}^{-2} \text{nm}^{-1}$
$L_{TOA}$	Total radiance received at the TOA, $\text{W sr}^{-1} \text{m}^{-2} \text{nm}^{-1}$
$L_w$	Water-leaving radiance, $\text{W sr}^{-1} \text{m}^{-2} \text{nm}^{-1}$
$n_w$	Refraction index of water
$Q$	Ratio of upwelling irradiance to upwelling radiance just beneath the surface
$R$	Irradiance reflectance
$r_\infty$	Bi-hemispherical infinite reflectance
$R^2$	Determination coefficient
$\rho_{nw}$	Normalized water-leaving reflectance
$R_{rs}$	Above-surface remote sensing reflectance, $\text{sr}^{-1}$
$r_{rs}$	Below-surface remote sensing reflectance, $\text{sr}^{-1}$
$r_{sd}^\infty$	Directional-hemispherical infinite reflectance
$S_{index}$	An index of bbp based on a sigmoid function
$S_{275-295}$	Spectral slope of $a_g(\lambda)$ between 275 and 295 nm, $\text{nm}^{-1}$
$S_{440}$	Spectral slope of $a_g(\lambda)$ calculated using 440 nm as reference, $\text{nm}^{-1}$
$S_{CDOM}$	Spectral slope of $a_g(\lambda)$ , $\text{nm}^{-1}$
$S_{NAP}$	Spectral slope of $a_{NAP}(\lambda)$ , $\text{nm}^{-1}$
$\theta_{s'}$	Under-water zenith angle, degree
$\theta_s$	Above-water solar zenith angle, degree
$u$	$= b_b / (a + b_b)$
$\mu_w$	Cosine of under-water zenith angle
$x$	$= b_b / a$
$Y$	The power-law exponent of $b_{bp}(\lambda)$ , $\text{nm}^{-1}$
$\eta$	Forward-backward asymmetry parameter for scattering coefficient
$z$	Depth, zero at surface and decrease in downward direction



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# **Chapter 1    Introduction**

## 1.1 Background

The propagation of light in natural waters is determined by how photons travel and change energy with depth in a water body. The properties of water that affect the light penetration are defined as the optical properties, which are conveniently divided into two mutually exclusive classes: the inherent optical properties (IOPs) and the apparent optical properties (AOPs) (Mobley, 1994). Inherent optical properties are those properties that depend only upon the medium, such as the absorption and scattering coefficients. IOPs of natural waters are determined by the optically significant constituents in the form of dissolved and particulate matter, which are essentially attributable to four components based upon the similarity in their optical properties: pure water, colored dissolved organic matter (CDOM), non-algae particles (NAP) and phytoplankton pigment (Chlorophyll). Apparent optical properties, on the other hand, depend on both the medium (the IOPs) and the geometric (directional) structure of the ambient light field, such as various reflectance (Kirk, 1994; Mobley, 1994). Two of the most important AOPs in ocean color remote sensing is the normalized water-leaving radiance ( $nL_w$ , in  $\text{W}\cdot\text{sr}^{-1}\cdot\text{m}^{-2}\cdot\text{nm}^{-1}$ ) and the remote sensing reflectance ( $R_{rs}$ , in  $\text{sr}^{-1}$ ). Note that  $R_{rs}$  is the ratio of  $nL_w$  to the extra-terrestrial solar irradiance ( $F_0$ , in  $\text{W}\cdot\text{m}^{-2}\cdot\text{nm}^{-1}$ ) (Gordon and Wang, 1994a). Theoretically, the composition of in-water constituents determines how much light is absorbed and scattered at a specific wavelength and thus affects the magnitude and shape of the spectral distribution of radiance that leaving from the water surface, such as  $nL_w$ . Reversely, with proper inversion, IOPs and the concentrations of in-water constituents can be derived from  $nL_w$  or  $R_{rs}$  measurements (Carder et al., 1999; Garver and Siegel, 1997; Hoge and Lyon, 1996; Lee et al., 2007; Lee et al., 2002; Lee et al., 2009; Maritorena et al., 2002; Salama and Verhoef, 2015; Werdell et al., 2013).

Traditional investigation of optical properties and their retrievals from AOPs is mainly based on discrete sampling, which is time-consuming and shows less success in interpreting the variations of optical properties at large spatial-temporal scales. Since the launch of Coastal Zone Color Scanner (CZCS) in 1978, researches on water optical properties have been seen great leap forward (IOCCG, 2000, 2006). The followed ocean color satellites, with greatly improved spatial and temporal resolutions and enhance radiometric measurements, allows simultaneous observation of the optical properties and interpretation of in-water constituents with much advanced accuracy. Among these satellites are the Sea-Viewing Wide Field of View Sensor (SeaWiFS, launched in 1999), the Moderate

Resolution Imaging Spectroradiometer (MODIS, launched in 2002), the Medium Resolution Imaging Spectroradiometer (MERIS, launched in 2002), the Geostationary Ocean Color Imager (GOCI, launched in 2010), the Visible Infrared Imaging Radiometer Suite (VIIRS, launched in 2012), the Ocean and Land Colour Instrument (OLCI, launched in 2016) and et cetera. However, as optical properties of natural waters have been reported with great temporal and spatial variabilities among rivers, lakes, coastal waters, and Open Ocean, it is still a challengeable task to obtain accurate products of IOPs and in-water constituents from remote sensing imagery.

The variabilities of optical properties among global waters are closely related to the concentration of dissolved and particulate substances, such as the CDOM, Chl-a, suspended sediments. Further, the in-water constituents could also have strong relationships with biogeochemical parameters, such as the strong correlations between CDOM and the dissolved organic carbon (DOC), and between SPM and the particulate organic carbon (POC) (Asmala et al., 2012; Ferrari, 2000; Fichot and Benner, 2011, 2012; Matsuoka et al., 2012a; Rochelle-Newall et al., 2014; Spencer et al., 2007). Therefore, remotely sensed optical properties could further be employed to retrieve biogeochemical parameters. Comparing to conventional investigation methods, e.g., discrete sampling, remote sensing of biogeochemical parameters provides synoptic coverage at large temporal-spatial scales allowing thereby a synoptic view of the aquatic system, and therefore represents a compelling alternative (Doerffer and Fischer, 1994; Gregg et al., 2017; Joshi et al., 2017; Kirk, 1994; Lee et al., 1999; Longhurst and Harrison, 1989; Mobley, 1994; Sathyendranath et al., 1989).

The fundamental task for remote sensing of the water optical properties, also referred to as ocean color, is to retrieve accurate products, including both the IOPs and the associated concentrations of in-water constituents, from airborne or spaceborne radiometric measurements. Ocean color remote sensing from satellite imagery includes two critical steps. The first step is to remove the atmosphere's contributions from the airborne or spaceborne measurements to obtain reliable  $nL_w$  or  $R_{rs}$ . This procedure is the so-called "atmospheric correction" (Gordon and Wang, 1994a). The second step is to quantitatively derive the optical properties and the concentrations of the water constituents from the water-leaving radiance or the  $R_{rs}$ . This step employs the "inverse algorithm" to retrieve IOPs and the in-water constituents from AOPs. However, due to the variability in the optical properties among different waters, numerous algorithms for atmospheric correction and IOPs retrieval have been developed (IOCCG, 2000, 2006).

### **1.1.1 Development of the atmospheric correction algorithm**

About 90% of the radiance seen by a satellite at the top of the atmosphere (TOA) is contributed by the atmosphere (including aerosols, water vapor, and clouds), and about 10% from the background signal from the ocean (Huot et al., 2001; Morel, 1980). In order to obtain target IOPs signals from satellite sensors, satellite imagery must be corrected to remove the radiance contributions from the atmosphere. Atmospheric effects are mainly attributed to multiple scattering by air molecules (Rayleigh scattering) and aerosols only, and the interaction term between molecular and aerosol scattering (Gordon and Wang, 1994b).

An effective method to remove the atmosphere's contribution from the total radiance received at TOA ( $L_{TOA}$ ) was proposed by Gordon and Wang (1994b), where the assumption was made that sea-water absorbs all the light in the red and near-infrared (NIR) region of the spectrum in clear ocean waters. This atmospheric correction scheme is referred to as GW94 and this assumption is well known as the black-pixel assumption. The GW94 allows an accurate estimation of the atmosphere's contribution for SeaWiFS products and can be easily adapted to other satellite products, especially for open ocean waters. However, the black-pixel assumption does not work well for optically complex waters, mostly in coastal regions and inland waters, because the non-negligible water-leaving radiance in the NIR will lead to an overestimation of the atmospheric radiance in the whole visible spectrum (IOCCG, 2000; Ruddick et al., 2000; Salama and Shen, 2010a; Siegel et al., 2000). To account for the complex optical properties, many algorithms have been developed to extend the GW94 model to process the atmospheric correction for the coastal waters (Bailey et al., 2010; Kuchinke et al., 2009; Ruddick et al., 2000; Stumpf et al., 2003).

For productive waters or turbid coastal waters, satellite-derived  $R_{rs}$  from the NIR atmospheric correction can be significantly underestimated and can even become negative in the blue bands (Wang et al., 2012). An improved scheme was proposed to remove the atmospheric effects for MODIS ocean colour products by using shortwave infrared (SWIR) bands, where the water can still be considered as black for turbid waters due to the strong absorption of water at SWIR bands. This algorithm is hereafter denoted as the SWIR atmospheric correction (Wang, 2007; Wang and Shi, 2005; Wang et al., 2009b; Wang et al., 2007). Furthermore, the SWIR atmospheric correction was proved to be more effective and flexible for turbid coastal waters compared with the NIR atmospheric correction, mainly because it can avoid masking out some of the

turbid water pixels as cloud pixels (Wang and Shi, 2006). However, the applicability of the SWIR atmospheric correction could be limited for many ocean color satellite sensors due to lacking the SWIR bands, such as the MERIS, OLCI, and GOCI. Alternative atmospheric correction algorithms were developed to optimize the atmospheric and oceanic signals at the NIR bands and then extrapolate to the visible bands by the neural network (NN) inversion method, which worked well for MERIS products over optically complex waters (Brajard et al., 2012; Schroeder et al., 2007).

For extremely turbid waters, a practical model for atmospheric correction based on radiative transfer theory was developed using MODTRAN simulations (Adler-Golden et al., 1999; Berk et al., 1999; Verhoef and Bach, 2003). With the assumption of a hypothetical homogeneous Lambertian surface of the water, a lookup table (LUT) can be generated from MODTRAN simulations for various atmospheric conditions (visibility and aerosol type) and viewing geometries and then applied to solve the inversion problem. Atmospherically corrected  $R_{rs}$  from MERIS based on the MODTRAN simulation was found to be promising when validated with matched *in-situ* measurements in the extremely turbid coastal waters in the Yangtze estuary (Shen et al., 2010b). Similar to the assumption of black water at NIR band in clear ocean waters, water-leaving radiance at ultraviolet (UV) wavelengths could be neglected compared with longer wavelengths in extremely turbid coastal waters due to strong scattering by the particles in the atmosphere and strong absorption by CDOM in the water (He et al., 2013; He et al., 2012). The implementation of the UV atmospheric correction on MODIS and GOCI images shows great improvements in the accuracy of satellite-derive  $R_{rs}$  with respect to the NIR and SWIR atmospheric correction (Wang and Shi, 2007). The UV atmospheric correction is also found more promising than the standard atmospheric correction approach of GOCI (Ahn et al., 2012) over the Yangtze estuary and the Hangzhou Bay (He et al., 2013; He et al., 2012).

### 1.1.2 Retrieval of IOPs and the concentrations of in-water constituents

While the atmospherically corrected spectra ( $nL_w$  or  $R_{rs}$ ) are of some interests in themselves, their main significance remains in that they can be used to remotely interpret the optical properties of the water body, such as the inherent optical properties. The two fundamental IOPs are the absorption coefficient ( $a(\lambda)$ , in  $\text{m}^{-1}$ ) and volume scattering function ( $\beta(\lambda)$ , in  $\text{m}^{-1} \cdot \text{sr}^{-1}$ ), which describe how light is

absorbed and scattered in the water. The scattering coefficient ( $b(\lambda)$ , in  $\text{m}^{-1}$ ) and the backscattering coefficients ( $b_b(\lambda)$ , in  $\text{m}^{-1}$ ) can be derived from the volume scattering function by integrating the  $\beta(\lambda)$  over all scattering angles and over only the backward scattering angles, respectively. Another ecologically essential IOP, or Quasi-IOP as some scholars may define, is the diffuse attenuation coefficient ( $K_d$ , in  $\text{m}^{-1}$ ), which describes the penetration of incident solar radiation in the water column. Many efforts have been made to remotely retrieve IOPs,  $K_d$ , and the concentrations of in-water constituents, such as chlorophyll-a, CDOM, and suspended particulate matter (SPM), from earth observation data since the launch of CZCS in 1978. Numerous inverse algorithms have been developed, which can be inclusively classified into three types, *i.e.*, the empirical approach, the theoretical approach, and the semi-analytical approach that employs both empirical relationships and analytical solutions.

#### *1.1.2.1 Empirical approach*

Empirical models are mostly employed in the retrievals of the concentration of optically active compositions (*e.g.*, Chl-*a* and SPM). It is not possible to analytically estimate the concentrations of Chl-*a* or SPM because their relationships with inherent optical properties are dependent on the physical characteristics of the particles, such as the types, composition, grain size, and et cetera. Developing an empirical model requires *a priori* established relationships between remotely measured reflectance, at a specific wavelength or combination of two or more wavelengths, and the concentrations of in-water constituents determined from *in-situ* measurements. Chl-*a* is one of the first in-water constituents that has been retrieved from remotely sensed data. Empirical models have been traditionally developed to retrieve Chl-*a* concentration using the blue-green band ratio of the  $nL_w$  or  $R_{rs}$  in the optically clear waters and using two-band and three-band algorithms for optically complex waters (Dall'Olmo et al., 2003; Gitelson, 1992; Gordon et al., 1983; Le et al., 2013; Le et al., 2009; O'Reilly et al., 1998; Shen et al., 2010c). Empirical models have also been extensively employed to retrieve the concentrations of CDOM and suspended sediments in coastal waters (D'Sa et al., 2007; Doerffer and Fischer, 1994; Liu et al., 2014; Loisel et al., 2014; Siswanto et al., 2011; Zhang et al., 2010). Other biogeochemical parameters, such as dissolved organic carbon (DOC), through its relationship with CDOM absorption properties, can also be empirically retrieved from earth observation data (Cao et al., 2018; Zhao and Song, 2018).  $K_d$ , as a QIOP, has also been empirically related to reflectance measurements for both



clear waters and optically complex waters (Austin and Petzold, 1981; Mueller, 2000; Werdell and Bailey, 2005; Zhang and Fell, 2007).

However, it is worthy to point out that the empirical models, although they are simple and easy to implement, have a major drawback regarding the model transferability. The empirical coefficients derived from statistical regression are highly dependent on the calibration dataset, and therefore only appropriate for waters with optical properties similar to those used in the model development. For example, most empirical models have satisfying performance in optically simple waters, usually referred to as the Case 1 water (Morel and Prieur, 1977), where absorption by Chl-*a*, apart from pure water, dominates the inherent optical properties. However, the performance of these models degrades noticeably when applied to optically complex waters, where strong absorptions by CDOM and NAP exist at the blue bands and the scattering from NAP cannot be negligible, without re-calibrating the empirical coefficients (Carder et al., 1989; Gitelson et al., 2007; Gons, 1999).

#### 1.1.2.2 *Semi-analytical methods*

To account for the relatively large and independent variations of Chl-*a*, CDOM, and NAP in optically complex waters, apart from the empirical models, semi-analytical and analytical models have been proposed. Analytical models are based on theoretical modelling of the underwater light field, such as the radiative transfer equations (RTEs) (Mobley, 1994). Due to the complexity of the shape of the volume scattering function and variations in the environmental factors that affect the light field in natural waters, exact analytical solutions of the RTEs exist only for a few idealized and unphysical situations such as no scattering. However, it is possible to arrive at approximate analytical solutions relating the AOPs and IOPs with proper assumptions and restrictive conditions, such as the single-scattering approximation (SSA) (Gordon, 1973).

Gordon and Morel (1983) proposed one of the most widely used analytical model based on simplified solutions of the RTEs, where the remote sensing reflectance just beneath the surface ( $r_{rs}$ , in  $\text{sr}^{-1}$ ) can be theoretically linked to inherent optical properties, namely the ratio of total backscattering coefficient ( $b_b$ ) to the sum of total absorption and backscattering coefficients ( $a + b_b$ ). The ratio is hereafter referred to as  $u(\lambda)$  (Gordon et al., 1988). The conversion of  $r_{rs}$  to the above-surface remote sensing reflectance ( $R_{rs}$ ) was also given as a function of the Fresnel surface reflectance (Gordon et al., 1988). With the proposed relationship between

$R_{rs}$  and  $u(\lambda)$ , knowledge of the absorption coefficient will enable the retrieval of backscattering or vice versa.

Many semi-analytical algorithms were developed on the basis of Gordon's equations with  $u(\lambda)$  first being derived from  $R_{rs}$ , and then decomposed (or inverted)  $u(\lambda)$  into IOPs of different components with proper empirical approach or iterative approach, including mainly the spectral optimization, spectral deconvolution, and bulk inversion (Hoge and Lyon, 1996; Werdell et al., 2013). Among these semi-analytical models are the Carder's algorithm for MODIS (Carder et al., 1999), the QAA algorithm (Lee et al., 2007; Lee et al., 2002; Lee et al., 2009) and the GSM algorithm (Garver and Siegel, 1997; Maritorena et al., 2002). Many other efforts have been made using neural network algorithms (Buckton et al., 1999; Doerffer and Schiller, 2007; Schiller and Doerffer, 1999) or nonlinear optimization (Doxaran et al., 2006; Salama et al., 2009; Salinas et al., 2007) to obtain the optimal solutions of the Gordon's equations to retrieve IOPs. It is worthy to note that the absorption coefficients by CDOM and NAP share similar spectral features and are difficult to be derived directly from the semi-analytical algorithms. CDOM and NAP absorption coefficients are usually treated as a combined parameter (CDM). Derived CDM from semi-analytical algorithms are further partitioned into individual absorption of CDOM and NAP using empirical models (Dong et al., 2013; Lee et al., 2002). With retrieved total absorption and backscattering coefficients,  $K_d$  can be then remotely retrieved (Doron et al., 2007; Lee et al., 2005a; Lee et al., 2005b; Morel and Loisel, 1998; Wang et al., 2009a).

For the retrieval of SPM concentrations, semi-analytical models have also been proposed based on radiative transfer (Doxaran et al., 2006; Shen and Verhoef, 2010; Shen et al., 2014; Volpe et al., 2011). In some other models, SPM concentrations are inverted from remotely retrieved IOPs with the assumption of a constant specific absorption or scattering coefficient, which is the ratio of  $a$  or  $b$  to SPM concentration. However, the relationship between IOPs and SPM could be highly dependent on the type, grain size, and composition of suspended particles. Therefore, the assumption of constant specific IOPs results in larger uncertainties when applied to a new region of interest and pre-determining the specific IOPs from locally calibrated data is necessary to improve model accuracy (Binding et al., 2005; Doxaran et al., 2006).

Semi-analytical models can retrieve IOPs for both optically simple and complex waters with satisfying accuracy. Compared with empirical models, although the

retrieval accuracy from semi-analytical models could be sometimes worse than for locally calibrated empirical models, semi-analytical models can be applied to different types of water without *a priori* calibrations on the empirical coefficients (IOCCG, 2006).

### 1.1.2.3 Analytical models

As mentioned in Section 1.1.2.2, most of the semi-analytical models adopted Gordon's solution of RTEs, while little has been done to improve the forward model for radiative transfer in water. Kubelka and Munk (1931) proposed a generic two-stream solution of the RTEs for turbid media, which has been validated for extremely turbid waters with very promising estimation of SPM concentration (Ambarwulan et al., 2012; Shen et al., 2010b). Duntley (1942) also introduced an approximate solution of the RTEs based on a two-stream approach that considers both direct and diffuse incident radiation. The Duntley model was later incorporated with an inverse scheme to analytically retrieve IOPs and  $K_d$ , namely, the 2SeaColor model (Salama and Verhoef, 2015). The 2SeaColor model includes the solution of the Duntley equation as the forward model and an inversion scheme using spectral optimization. Yu et al. (2016) revised the 2SeaColor model with improved IOPs parameterizations, which produced more accurate and reasonable estimations of the IOPs and  $K_d$  for the Yangtze estuary.

In this dissertation, the author first employed the 2SeaColor model for the retrieval of IOPs and  $K_d$ , then utilized locally calibrated models to estimate SPM concentration from derived particulate backscattering coefficient and derive DOC concentration from CDOM absorption coefficients, respectively. Empirical and semi-analytical models from other publications were also adopted in this study for model inter-comparison.

## 1.2 The Yangtze estuary

The Yangtze River is third largest in length, fifth largest in water discharge, and fourth largest in sediment load, after the Amazon, Zaire and Orinoco, with average annual water and sediment discharge into the estuary up to  $8.96 \times 10^{11} \text{ m}^3$  and  $3.9 \times 10^{11} \text{ kg}$  from 1950 to 2010, respectively. It originates from Geladaindong Peak on the Tibetan Plateau and runs into the East China Sea at Shanghai city. The rich water and sediments discharges make the Yangtze estuary one of the world's most turbid estuaries. The estuary, mainly the inshore area and the river mouth, is affected by a strong semidiurnal tide where tidal amplitudes are average between 4.5 and 5.0 m (Huang, 1992). The Yangtze delta area is the most

populous and economically developed region in China, which accounts for more than 20% of China's Gross Domestic Product.

The Yangtze estuary, same as world's major estuaries, functions as the exchange hub of the marine ecosystem and land ecosystem where massive material exchanges between land and ocean. The material exchanges, especially the change of carbon as dissolve substance, promote the global carbon cycle and maintain the sustainability of the environment. Besides, the estuary ecosystem is heavily affected by human activities and very sensitive to global climate change. Therefore, monitoring in the optical properties of the Yangtze estuary could be of great scientific, economic, and social benefits. Furthermore, optical measurements can provide insights into the aquatic ecosystem, such as the health status of the water body, biological productivity, sediment load, pollutants, bathymetry, and carbon stocks et cetera. (Doerffer and Fischer, 1994; Gregg et al., 2017; Joshi et al., 2017; Kirk, 1994; Lee et al., 1999; Longhurst and Harrison, 1989; Mobley, 1994; Sathyendranath et al., 1989).

### 1.3 Challenges

Water optical properties in the Yangtze estuary are frequently changed due to seasonal riverine discharges, tidal currents, and wind-driven resuspensions of bottom sediments, fresh and salt water mixing, resulting in one of the world's most optically-complex regions. For turbid waters, one of the challenges for ocean color remote sensing is the scarcity of "ground-true" data for model validation, which is because current commercial instruments for IOPs measurements, such as Wet Labs ac-s meter and bb9 meter, are prone to saturate in high turbidity environments. The other challenges is that due to the large variation of the IOPs and concentrations of in-water constituents, most semi-analytical models have been less successful in retrieving IOPs or  $K_d$  (IOCCG, 2000, 2006; Salama et al., 2011; Wang et al., 2017; Zibordi et al., 2009), which is mainly attributed to the strong backscattering by the suspended particles. Therefore, developed models for the retrieval of SPM and Chl-a concentrations in the Yangtze estuary are all empirical or semi-analytical based (Shen et al., 2010b; Shen et al., 2010c; Shen et al., 2013; Sokoletsky et al., 2014). Studies based on radiative transfer approach, on the other hand, are hardly discussed in published literature. Despite the lack of *in-situ* measured IOPs for model validation, the comparison between analytical model and semi-analytical models could still provide compelling insights into the models' performance in turbid waters, which will be discussed in this dissertation.

SPM concentration was reported to be less than 1 mg/L in the offshore waters of the Yangtze estuary and beyond 5000 mg/L in the Hangzhou Bay (He et al., 2013). Due to the large variation of SPM concentration in the Yangtze estuary, the currently proposed SPM models have to include a switching scheme to account for the less turbid and extremely turbid waters (Shen et al., 2010b; Shen et al., 2013), which in return results in discontinuity issue with derived SPM. To avoid the discontinuity issue and produce reasonable and realistic mapping products of SPM, dedicated efforts are required to improve the SPM retrieval models.

## **1.4 Research objectives**

The Yangtze estuary is rich in both particulate and dissolved substances that originate from terrestrial sources discharged from the Yangtze River. The particulate substance, *i.e.*, SPM, is mainly due to land erosion in the river basin while the dissolved substance, *i.e.*, CDOM, is related to agriculture activities. The abundance of SPM and CDOM, both significantly restrict light penetration in water, make the optical properties of the Yangtze estuary much more complex than the majority of the world's estuaries. Thereby, investigating the optical properties of the Yangtze estuary is of great interest for IOP retrieval models' development and evaluation in optically-complex waters.

The main objective of this dissertation is to investigate the optical properties in the Yangtze estuary and develop fairly accurate models to retrieve both the optical properties and the concentrations of water constituents. The main objective can be fulfilled by the three sub-objectives described below.

- 1) Improve the analytical 2SeaColor model to enhance its accuracy in the Yangtze estuary in retrieving inherent optical properties and diffuse attenuation coefficient.
- 2) Explore the relationship between particulate backscattering properties and SPM concentration in the Yangtze estuary and develop a novel model to retrieve the wide-range SPM concentration.
- 3) Investigate the seasonal variability of CDOM absorption properties in the Yangtze estuary and propose a regionally applicable model for the estimation of dissolved organic carbon concentration from CDOM.

## **1.5 Outline of this thesis**

This thesis consists of six chapters. First two chapters are the general introduction and detailed description of the analytical 2SeaColor model, followed by three

core chapters and a last chapter of synthesis. Each core chapter addresses one of the sub-objectives mentioned above and has been prepared as a standalone research paper that has been published in or submitted to the peer-reviewed journal. The six chapters are arranged as follows:

Chapter 1 describes the background of ocean color remote sensing and some state-of-art developments of models in optically-complex water that remotely retrieve inherent optical properties and the concentrations of in-water constituents.

Chapter 2 introduces the general concept of radiative transfer theory and provide a comprehensive description of the two-stream radiative transfer model, *i.e.*, the 2SeaColor model, which includes both the original 2SeaColor model and its updated version with improved parametrizations.

Chapter 3 describes the retrieval of the diffuse attenuation coefficient ( $K_d$ ) in the Yangtze estuary using the 2SeaColor model with improved parameterization and compares the model's performance with empirical and semi-analytical  $K_d$  retrieval models.

In Chapter 4, the derived particulate backscattering coefficient ( $b_{bp}$ ) from the 2SeaColor model is used to quantitatively estimate the concentration of suspended particulate matter ( $C_{SPM}$ ). The proposed model utilizes a  $b_{bp}$  index based on sigmoid function to accurately retrieve wide-range  $C_{SPM}$  between 0.4 and 2000 mg/L without the need for a switching scheme.

Chapter 5 presents the seasonal variability of the light absorption properties of colored dissolved organic matter (CDOM) in the Yangtze estuary and the robust relationship between dissolved organic carbon and CDOM absorption properties, includes the absorption coefficient and the spectral slope of the absorption coefficient.

Chapter 6 is a Synthesis of the results obtained in this dissertation. It includes the main conclusions of Chapter 3, 4, and 5 and recommendations for future efforts in ocean color remote sensing in turbid coastal waters.

## **Chapter 2    Principal Basics and Methods**

## 2.1 Radiative transfer in the water

Light can be described by various physical and mathematical quantities in terms of the radiance. The radiative transfer theory explains how radiance varies with distance along any specific path at a specified point in the medium (Mobley, 1994). The radiative transfer theory is widely used to explain different kinds of science questions, among them is the radiative transfer in atmosphere and water. Figure illustrates how light propagates and interacts with the medium within the atmosphere and water in the coastal environment. For remote sensing applications, the atmospheric correction algorithm focuses on the light propagation between just above the water surface and TOA and utilizes the radiative transfer equations (RTE) for the atmosphere, while retrieval model of inherent optical properties targets at the light transfer inside the water column.

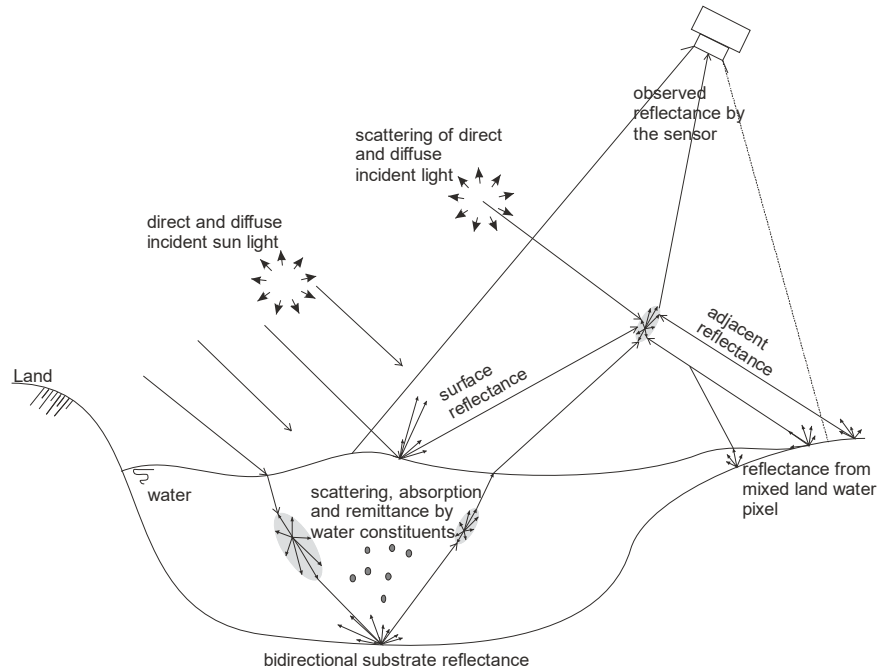


Figure 2.1. Light pathways in the atmosphere and water in coastal regions (adapted from Su et al., 2011).

The radiative transfer equation, commonly called the RTE, expresses the conservation of energy written for a collimated beam of radiance traveling through an absorbing, scattering and emitting medium. Therefore, radiative transfer equations can tie together the radiance and the IOPs, which provide the theoretical framework for the optical oceanography and ocean-color remote sensing. A useful expression can be derived from the RTEs to relate the IOPs to



the average cosine and the vertical attenuation coefficient for net downward irradiance. With IOPs and boundary conditions go into the equations, radiance can be predicted. However, solutions of RTEs must be done numerically except for a few trivial cases such as non-scattering media, which can be hardly found in natural waters.

Due to the complexity of the shape of the volume scattering function and variations in the environmental factors that affect the light field in natural waters, retrieval of optical properties from exact analytical solutions of the RTEs is not possible for natural waters, as the conditions and boundaries for RTEs are hard to be exactly quantified. However, with pre-assumed boundary conditions, the equations arising from the RTE theory can be solved with simplified solutions. One of these approximate solutions is the single-scattering approximation (SSA). This solution is developed by assuming that the water is homogeneous, the sea surface is level, the sun is a point source in a black sky, there are no internal sources, and only single-scattering of photons is considered (Gordon and Wang, 1994b). The SSA is generally accepted and validated in optical simple waters, where numerous models, analytical and semi-analytical, were developed for the retrieval of IOPs (Gordon et al., 1988; Gordon and Morel, 1983; Lee et al., 1999; Lee et al., 2002). The other approximate solution of RTEs is based on the two-stream approach in which direct and diffuse incident radiation is considered with the assumption of water as an optically deep medium (Duntley, 1942). Salama and Verhoef (2015) later developed the 2SeaColor model that includes the solution of the Duntley equation and an inversion scheme using spectral optimization for the retrieval of IOPs. In this dissertation, the author focuses on the improvements of the 2SeaColor model and its application of retrieving IOPs in turbid coastal waters.

## 2.2 The 2SeaColor model

### 2.2.1 Forward model

Optical properties of water are usually described by means of the absorption coefficient  $a(\lambda)$  and the backscattering coefficient  $b_b(\lambda)$  (wavelength dependency are omitted in this Section for brevity). Radiative transfer inside the water column can be described by the Duntley equations as follows (adapted from Duntley (1942)):

$$\frac{dE_s}{dz} = kE_s \quad (2.1)$$

$$\begin{aligned}\frac{dE^-}{dz} &= -s'E_s + \alpha E^- - \sigma E^+ \\ \frac{dE^+}{dz} &= sE_s + \sigma E^- - \alpha E^+\end{aligned}$$

where  $E_s$  is the direct solar flux, and  $E^-$ ,  $E^+$  are the diffuse downward and upward fluxes,  $z$  is supposed to increase in an upward direction and to be zero at the top,  $k$  is the extinction coefficient for direct sunlight and can be calculated by total absorption ( $a$ ) and scattering coefficient ( $b$ ),

$$k = \frac{a + b}{\cos \theta_s'}, \quad (2.2)$$

where  $\theta_s'$  is the under-water zenith angle, and can be found with Snell's law from above-water solar zenith angle ( $\theta_s$ ),  $\theta_s' = \arcsin(\sin \theta_s / n_w)$ , and  $n_w$  is the refraction index of water, which is about 1.33.

The  $s$  and  $s'$  in system (2.1) are the backscattering and forward scattering coefficients for direct sunlight and given by,

$$s = \frac{b_b}{\cos \theta_s'}, \quad (2.3)$$

and

$$s' = \frac{b_f}{\cos \theta_s'}, \quad (2.4)$$

where  $b_b$  and  $b_f$  are scattering coefficients in the backward and forward direction. Total scattering is then given by  $b = b_b + b_f = (1 + \eta)b_b$ , and the  $\eta$  is the forward-backward asymmetry parameter.

The  $\alpha$  and  $\sigma$  in system (2.1) are the attenuation coefficient and diffuse backward scattering coefficient and are given by  $\alpha = 2(a + b_b)$  and  $\sigma = 2b_b$ . The eigenvalue  $m$  is defined as  $m = \sqrt{\alpha^2 - \sigma^2}$ .

For the solutions to the diffuse fluxes in system (2.1), the total downward planar irradiance of diffuse and direct light is given by,

$$E_d(z) = E^-(z) + E_s(0)e^{kz} \quad (2.5)$$

The diffuse upward irradiance is expressed as,

$$E^+(z) = r_\infty E^-(z) + r_{sd}^\infty E_s(z) \quad (2.6)$$

where  $r_\infty$  is the bi-hemispherical reflectance for the semi-infinite medium and  $r_{sd}^\infty$  is the directional-hemispherical reflectance of the semi-infinite medium.

The downwelling attenuation coefficient of diffuse and direct light is the analytical differentiation of the downwelling irradiance with depth,

$$K_d = \frac{1}{E_d} \frac{dE_d}{dz} \quad (2.7)$$

The downward diffuse irradiance is derived as:

$$E^-(z) = E^-(0)e^{mz} + (s' + \sigma r_{sd}^\infty)E_s(0)J_1(k, z) \quad (2.8)$$

with,

$$J_1(k, z) = \begin{cases} \frac{e^{mz} - e^{kz}}{k - m} & (k \neq m) \\ -ze^{kz} & (k = m) \end{cases} \quad (2.9)$$

The term  $r_\infty$  can be written as,

$$\begin{aligned} r_\infty &= \frac{\sigma}{\alpha + m} = \frac{2b_b}{2(a + b_b) + \sqrt{4(a + b_b)^2 - 4b_b^2}} \\ &= \frac{b_b}{a + b_b + \sqrt{a(a + 2b_b)}} \end{aligned} \quad (2.10)$$

The term  $r_{sd}^\infty$  is given by,

$$r_{sd}^\infty = \frac{s + s'r_\infty}{k + m} \quad (2.11)$$

Substitution of water inherent optical properties under the assumption  $\eta = 1$  gives,

$$r_{sd}^\infty = \frac{(1 + r_\infty)b_b/\mu_w}{(a + 2b_b)/\mu_w + m} \quad (2.12)$$

where  $\mu_w = \cos \theta_s'$ . Similarly, we can further write,

$$\begin{aligned} r_{sd}^\infty &= \frac{\sqrt{a + 2b_b}(\sqrt{a + 2b_b} - \sqrt{a})}{(a + 2b_b) + 2\mu_w\sqrt{a(a + 2b_b)}} \\ &= \frac{\sqrt{1 + 2b_b/a} - 1}{\sqrt{a + 2b_b/a} + 2\mu_w} \end{aligned} \quad (2.13)$$

By introducing the ratio  $x = b_b/a$  into Eq. (2.10) and Eq. (2.13), one can finally write,

$$r_\infty = \frac{x}{1 + x + \sqrt{1 + 2x}} \quad (2.14)$$

$$r_{sd}^{\infty} = \frac{\sqrt{1 + 2x} - 1}{\sqrt{1 + 2x} + 2\mu_w} \quad (2.15)$$

In order to obtain remote sensing reflectance ( $R_{rs}$ ) from the forward model, we first combine Eq. (2.5) and Eq. (2.6) to calculate the irradiance reflectance just beneath the surface as,

$$R = \frac{E^+(0)}{E_d(0)} = \frac{r_{\infty}E^-(z) + r_{sd}^{\infty}E_s(z)}{E_d(0)} \quad (2.16)$$

The  $R$  is later translated to  $R_{rs}$  following Mobley (1994),

$$R_{rs} = \frac{0.52R}{Q - 1.7R} \quad (2.17)$$

where  $Q$  is the ratio of the water upwelling irradiance to upwelling radiance just beneath the surface.  $Q$  is approximate to 3.25 for clear sky condition (Morel and Gentili, 1993; Lee and Carder et al., 1998).

### 2.2.2 Inverse Scheme

Remote sensing reflectance is fundamentally a measure of inherent optical properties, i.e., the bulk absorption coefficient and backscattering coefficient, and therefore can be inverted from known  $a$  and  $b_b$ . Following Eq. (2.14) and Eq. (2.15), the inherent optical properties  $x$  ( $=b_b/a$ ) can be linked with the 2SeaColor model predicted reflectance results  $r_{\infty}$  and  $r_{sd}^{\infty}$ . However,  $r_{\infty}$  holds only for diffuse incident radiation, and under sunny conditions, it is better to use  $r_{sd}^{\infty}$  as a basis for estimation of water optical properties. Therefore,  $r_{sd}^{\infty}$  in Eq. (2.15) can be inverted to yield the inherent optical properties  $x$  by,

$$x = \frac{r_{sd}^{\infty}(1 + 2\mu_w)[1 + r_{sd}^{\infty}(\mu_w - 0.5)]}{(1 - r_{sd}^{\infty})^2} \quad (2.18)$$

The  $r_{sd}^{\infty}$  can be approximated by  $Qr_{rs}$ , where  $Q$  is approximate to 3.25 for clear sky condition (Morel and Gentili, 1993; Lee and Carder et al., 1998) and  $r_{rs}$  is the remote sensing reflectance just beneath the surface. Further, an empirical relationship is given to convert  $r_{rs}$  to  $R_{rs}$  (Lee et al., 2002),

$$r_{rs} = \frac{R_{rs}}{0.52 + 1.7R_{rs}} \quad (2.19)$$

Note that the bulk absorption and backscattering coefficients can be parameterized as the total of optical properties by different water constituents, which gives,

$$a(\lambda) = a_w(\lambda) + a_{\phi}(\lambda) + a_g(\lambda) + a_{NAP}(\lambda) \quad (2.20)$$

and

$$b_b(\lambda) = b_{bw}(\lambda) + b_{bp}(\lambda) \quad (2.21)$$

where the subscripts  $w$ ,  $\phi$ ,  $g$ , and  $NAP$  in Eq. (2.20) represent pure water, phytoplankton pigment, CDOM, and NAP, respectively.  $b_{bw}$  and  $b_{bp}$  are referred to as the backscattering coefficients of pure water and NAP, with the assumption that backscattering coefficients of CDOM and phytoplankton pigment are negligible. Absorption and backscattering coefficients of pure water are adopted from (Mobley, 1994; Morel, 1974; Smith and Baker, 1981).

In the original 2SeaColor model, water molecules are assumed to be the only contributor to absorption in the near-infrared (NIR) band (Salama and Verhoef, 2015). Therefore,  $b_b$  at the NIR bands can be first estimated from the ratio  $x$  and the pure water absorption coefficient ( $a_w$ ). Then  $b_b(\text{NIR})$  can be extrapolated to shorter wavelengths by,

$$b_{bp}(\lambda) = b_{bp}(\lambda_0) \left( \frac{\lambda_0}{\lambda} \right)^Y \quad (2.22)$$

where  $\lambda_0$  is the reference wavelength and  $Y$  is the power law exponent for the particle backscattering coefficient. Further, total absorption coefficient  $a(\lambda)$  is derived from  $b_b(\lambda)$  and the ratio  $x(\lambda)$ .

The parameterizations of the  $a(\lambda)$  and  $b_b(\lambda)$  were later improved to account for optically complex waters where residual absorptions by NAP cannot be neglected in NIR bands (Yu et al., 2016). As implemented in Salama and Shen (2010b), absorption and backscattering coefficients of water constituents are parameterized by (Bricaud et al., 1981; Lee et al., 1998; Sathyendranath et al., 2001),

$$a_\phi(\lambda) = (a_0(\lambda) + a_1(\lambda) \ln[a_\phi(440)]) a_\phi(440) \quad (2.23)$$

$$a_{dg}(\lambda) = a_{dg}(\lambda_0) e^{-S(\lambda - \lambda_0)} \quad (2.24)$$

with  $a_0$  and  $a_1$  given in Lee et al. (1998). The subscript  $dg$  stands for the combination of CDOM and NAP, and  $S$  is the spectral slope of the combined absorption coefficient.

For remote sensing applications, we proceed firstly by inverting the measured spectra of  $R_{rs}$  or satellite  $R_{rs}$  products into inherent optical properties. From Eq. (2.18) to Eq. (2.24), one can simulate  $R_{rs}$  for given  $a_\phi(440)$ ,  $a_{dg}(\lambda_0)$ ,  $b_{bp}(\lambda_0)$ ,  $S$ , and  $Y$ . Therefore, the five unknowns can be retrieved from measured  $R_{rs}$  spectra as long as  $R_{rs}$  data has more than 5 spectral bands. In this study, initial values of  $a_\phi(440)$ ,  $a_{dg}(\lambda_0)$ ,  $b_{bp}(\lambda_0)$  and  $Y$  are estimated from band-ratios of measured

$R_{rs}$  (Lee et al., 1999).  $S$  is set to  $0.015 \text{ nm}^{-1}$  (Lee et al., 2007; Lee et al., 2002). A spectral optimization program using nonlinear curve fitting (lsqcurvefit in Matlab®) was applied to compute the sum of squared differences between the model-simulated  $R_{rs}$  curve and the field measured  $R_{rs}$  curve and iterates until the minimum is reached, at which point the values of  $a_\phi(440)$ ,  $a_{dg}(\lambda_0)$ ,  $b_{bp}(\lambda_0)$  and  $Y$  are considered to be retrieved. These unknowns can further be used to estimate the spectral absorption and backscattering coefficient of water constituents according to Eq. (2.20) to Eq. (2.24). The  $K_d$  is meanwhile calculated from the derived  $a(\lambda)$  and  $b_{bp}(\lambda)$  according to the analytical forward model described in Section 2.2.1.

### 2.3 Models sensitivities to SPM concentration

The reason that most IOPs retrieval models show less success in optical-complex waters, especially in turbid waters, lies mainly in the limitations of the radiative transfer equation used in these models. In this section, we evaluate the sensitivities of a couple of well-known IOPs retrieval models to increasing SPM concentrations. These models are,

- (1) the QAA (Lee et al., 2002)

$$r_{rs} = \sum_{i=1}^2 g_i \left( \frac{b_b}{b_b + a} \right)^i \quad (2.25)$$

where  $g_1=0.0895$  and  $g_2=0.1247$ .

- (2) The GSM model (Garver and Siegel, 1997; Maritorena et al., 2002),

$$R_{rs} = \frac{t^2}{n_w^2} \sum_{i=1}^2 g_i \left( \frac{b_b}{b_b + a} \right)^i \quad (2.26)$$

where  $t$  is sea–air transmission factor and  $n_w$  is the index of refraction of the water.  $g_1=0.0949$  and  $g_2=0.0794$ .

- (3) The Kubelka-Munk model (Kubelka and Munk, 1931),

$$r_{rs}(\lambda) = \frac{\sqrt{1 + 2x} - 1}{\sqrt{1 + 2x} + 1} \quad (2.27)$$

where  $x = b_b/a$  and the model is referred to as the K-M model for brevity.

- (4) The Duntley model, which is expressed in Eq. (2.15).

To understand how SPM concentration can affect these models, we made some assumptions that first only water absorbed in these models and thus  $a$  equals  $a_w$  over the whole spectral range. Second, the backscattering coefficients of CDOM and chlorophyll are negligible in the turbid water, which gives  $b_b$  as a total of  $b_{bw}$  and  $b_{bp}$ . Moreover,  $b_{bp}$  is assumed to be linearly related to the SPM concentration ( $C_{SPM}$ ) and at a reference wavelength. Then we define,

$$x \equiv \frac{b_b}{a} = \frac{b_{bw} + b_{bp}}{a_w} \rightarrow b_{bp} = a_w \times x - b_{bw} = f \times C_{SPM} \quad (2.28)$$

where  $f$  is the empirical coefficient of the linear relationship between  $C_{SPM}$  and  $b_{bp}$ .

Based on the above assumptions and Eq. (2.28), the variation of under-water irradiance reflectance predicted by these four models with increasing SPM concentration is presented in Figure 2.2. Note that the assumptions may not hold in natural waters, especially in optically-complex waters, but the simplification of the IOPs could still provide some insights into how different models respond to increasing water turbidity.

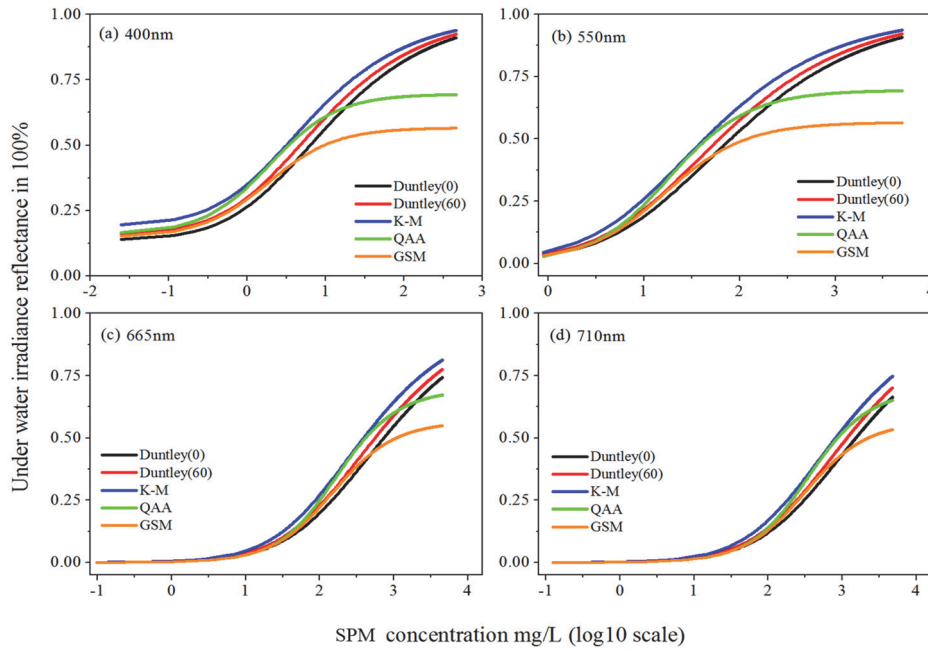


Figure 2.2. Model sensitivities to SPM concentrations at different reference wavelengths ('0' and '60' for the Duntley model refer to the under-water solar zenith angle).

As shown in Figure 2.2, under-water irradiance reflectance predicted by the four IOPs demonstrate different trends with the increasing of  $C_{SPM}$ . The predictions of QAA and GSM model present a similar pattern, as both of them employ the quadratic relationship and their major difference lies in the coefficients of the quadratic equations. These two models present an earlier saturation of the under-water irradiance reflectance compared with the K-M model and the Duntley model, which indicates they will be less sensitive to very high  $C_{SPM}$ . Therefore, using the K-M model or Duntley model could be more suitable to formulate the relationship between IOPs and reflectance in optically-complex waters. As expressed in Eq. (2.15) and Eq. (2.27), the prediction from the Duntley model would be equal to the K-M model when the  $\mu_w$  in Eq. (2.15) equals 0.5, which corresponds to an under-water zenith angle of  $60^\circ$ . However, this assumption can never happen for sunlight under water since the maximum angle of refraction is  $48.6^\circ$ . Thus, the results of the K-M model may lead to an overestimation in the calculation of under-water irradiance reflectance. It is worthy to note that results in Figure 2.2 also indicate that the saturation effects could shift from blue bands to longer wavelengths with increasing SPM, which is also confirmed from above water measured  $R_{rs}$  (Knaeps et al., 2012; Knaeps et al., 2015; Luo et al., 2018; Shen et al., 2010b).



## **Chapter 3      Retrieval of the diffuse attenuation coefficient using the 2SeaColor model: A case study in the Yangtze Estuary\***

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\* This Chapter is based on:

Yu, X., Salama, M.S., Shen, F., & Verhoef, W. (2016a). Retrieval of the diffuse attenuation coefficient from GOCI images using the 2SeaColor model: A case study in the Yangtze Estuary. *Remote Sensing of Environment*, 175, 109-119

### 3.1 Introduction

The diffuse attenuation coefficient ( $K_d$ , in  $\text{m}^{-1}$ ) is an important optical property of water bodies that describes the penetration of incident solar radiation in the water column.  $K_d$  is classified as a quasi-inherent optical property (QIOP), since it depends on both the optically active water constituents and the distribution of the ambient light field (Kirk, 1984, 1991; Mobley, 1994; Smith and Baker, 1978).

Based on the fact that 90% of the diffusely reflected light from a water body comes from a surface layer of water within a depth of  $1/K_d$  m, the diffuse attenuation coefficient shows potential effects on the functioning of ecological processes and biogeochemical cycles in oceanic and coastal ecosystems. Previous researches have shown that  $K_d$  can be an effective indicator of turbidity in water and sediment transportation and resuspension (Kirk, 1994; Majozi et al., 2014; Zhang et al., 2006). Meanwhile,  $K_d$  is a key parameter in studies such as heat transfer in the upper water layer (Stramska and Zuzewicz, 2013; Wu et al., 2007), phytoplankton photosynthesis and photobleaching (Loiselle et al., 2009; McClain et al., 1996), and the estimation of primary production in natural waters (Bergamino et al., 2010; Lee et al., 2011).

Algorithms for the remote retrieval of  $K_d$  have been developed for clear waters (Austin and Petzold, 1981; Mueller, 2000), but they always generate huge errors in turbid waters due to the complexity in optical properties and limitations in the algorithms (IOCCG, 2000, 2006; Salama et al., 2011; Zibordi et al., 2009). To retrieve  $K_d$  in optically complex waters, several methods were adopted. The empirical models were developed by statistical fitting based on *in-situ* measurements (Werdell and Bailey, 2005; Zhang and Fell, 2007). The semi-analytic models usually calculated the  $K_d$  from remotely retrieved inherent optical properties (IOPs) of water constituents (Doron et al., 2007; Lee et al., 2005a; Lee et al., 2005b; Morel and Loisel, 1998; Wang et al., 2009a).

Additionally, Salama and Verhoef (2015) developed the analytical 2SeaColor model to estimate  $K_d$  at different spectral bands and depths, based on the solution of the radiative transfer equation using the two-stream approximation. The two-stream approximation was previously adopted in the retrieval of suspended matter (Ambarwulan et al., 2012; Shen et al., 2010b). However, the assumption in the original 2SeaColor model that only water molecules contribute to the total absorption in the near-infrared (NIR) bands may not work for highly turbid waters, which leads to large uncertainties for observations from turbid waters. Therefore,

an improvement in the 2SeaColor model is strongly desired for  $K_d$  estimation over turbid waters.

The objective of this study is to improve the 2SeaColor model and validate the respective performance of the 2SeaColor model and other empirical and semi-analytical models. We first improve the parameterizations in the inversion scheme of the 2SeaColor model. Then the 2SeaColor model is validated with three *in-situ* datasets and compared with empirical and semi-analytic models. At last, we produce  $K_d$  maps over the Yangtze Estuary by implementing the 2SeaColor model on the level 2 product (L2P) of the Geostationary Ocean Color Imager (GOCI).

## 3.2 Data and methods

### 3.2.1 Datasets

Three *in-situ* datasets are employed in this study for model validation and inter-comparison, which are collected from a variety of waters covering from Open Ocean to coastal regions. They are the NOMAD (NASA Bio-Optical Marine Algorithm Data set) (Werdell and Bailey, 2005), the COASTLOOC (Coastal Surveillance through Observation of Ocean Color) (Babin et al., 2003) and *in-situ* measurements in Yangtze estuarine and coastal water, hereafter denoted as the Yangtze dataset. Figure 3.1 presents the locations of the *in-situ* observations of these three datasets. Solar zenith angle (SZA, in degree) of each observation is calculated according to location (latitude and longitude and time zone), date and time (C++ code is available at <http://www.psa.es/sdg/sunpos.htm>). The percentage distribution of SZA for each dataset is presented in Figure 3.2.

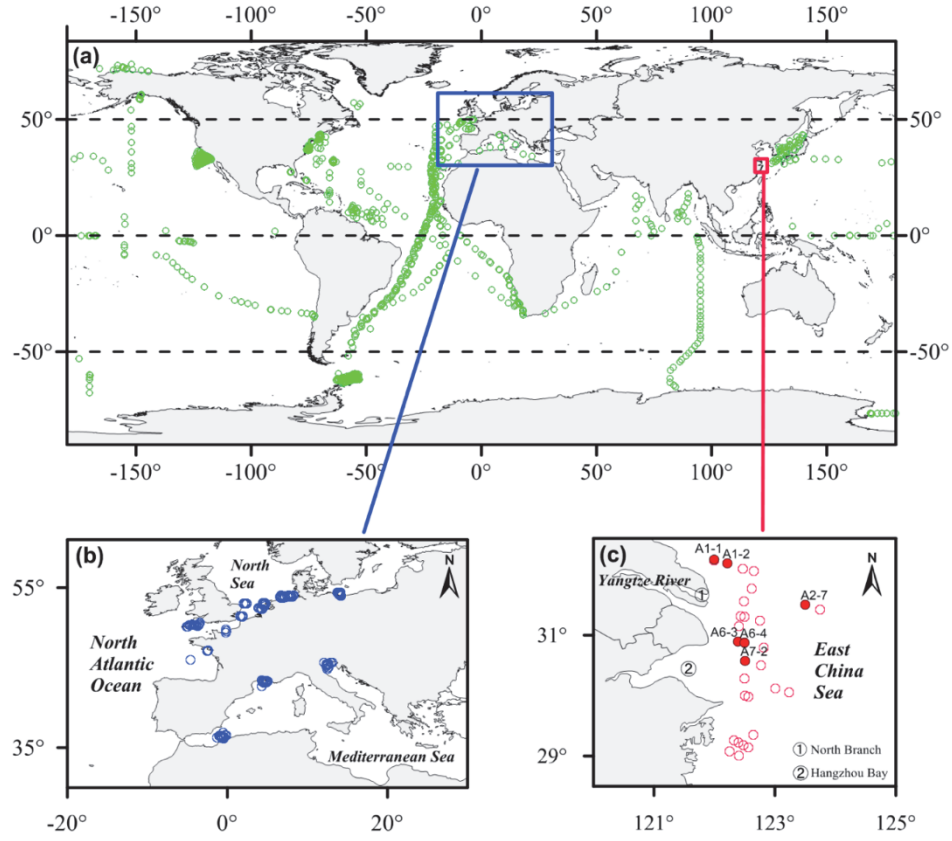


Figure 3.1. Location of the field data sets. (a) The NOMAD, (b) the COASTLOOC and (c) the Yangtze dataset (solid circles are matchups sites within 1h of GOCI capturing images).

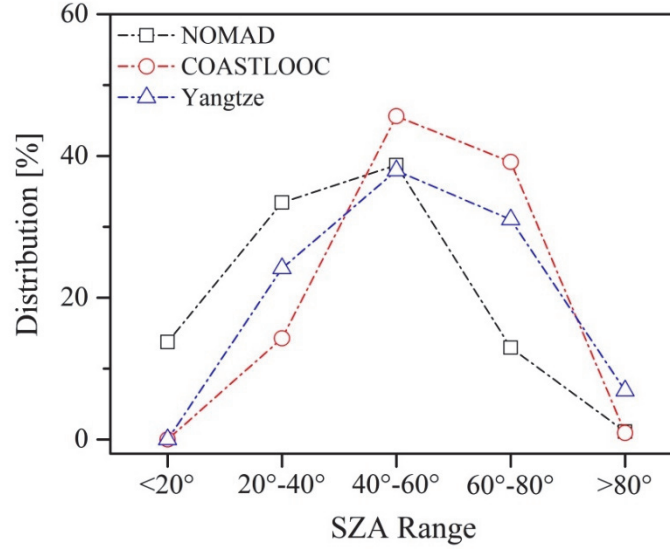


Figure 3.2. Percentage distribution of solar zenith angle (SZA) in the three in-situ datasets.

#### 3.2.1.1 NOMAD dataset

The NOMAD dataset (Werdell and Bailey, 2005) contains 4459 observations of bio-optical parameters collected worldwide from Open Ocean to coastal waters. Observations were selected where the upwelling water radiance ( $L_w$ ) and the spectral downwelling irradiance ( $E_s$ ) in the bands at 411, 443, 489, 510, 555, and 665 nm are all valid, as well as the diffuse attenuation coefficient at 443 and 489 nm. As a result, 1518 concomitant observations were selected (sampled sites are shown in Figure 3.1a). The diffuse attenuation coefficient at 489nm,  $K_d(489)$ , of the selected observations ranges from  $0.0085 \text{ m}^{-1}$  to  $0.9113 \text{ m}^{-1}$  with an average value of  $0.0852 \text{ m}^{-1}$  and a standard deviation of  $0.105 \text{ m}^{-1}$ . Remote sensing reflectance ( $R_{rs}$ ) is calculated as the ratio of upwelling water radiance  $L_w$  to the downwelling irradiance  $E_s$ , and  $K_d(489)$  is considered equivalent to  $K_d(490)$  for model validation and inter-comparison in this study.

#### 3.2.1.2 COASTLOOC dataset

The COASTLOOC dataset (Babin et al., 2003) consists of 422 observations of bio-optical parameters collected from European coastal waters. The irradiance reflectance just below the sea surface  $R(0^-)$  is available in the COASTLOOC data set instead of  $R_{rs}$ . Therefore, the  $R(0^-)$  was converted to  $R_{rs}$  using the formulation used in Zhang and Fell (2007) where  $R_{rs} = 0.133R(0^-)$ . It is worthy to mention that this formulation uses a constant value for the upwelling irradiance-to-

radiance ratio ( $Q = 4$ ), which will generate uncertainties to the calculated  $R_{rs}$ . However, the effect of the  $Q$  value on the derived  $K_d$  is limited when the ratio of  $R_{rs}$  or the spectral optimization technique is used, which will cancel out much of the dependencies. The COASTLOOC dataset was also filtered and 217 concomitant observations were selected where  $R_{rs}$  and  $K_d$  in the bands at 411, 443, 489, 510, 555, 665 and 705 nm are all valid, as well as the solar zenith angle  $\theta_s$  (sampled sites are shown in Figure 3.1b). The  $K_d(490)$  of the selected observations ranges from  $0.025 \text{ m}^{-1}$  to  $2.976 \text{ m}^{-1}$  with an average value of  $0.490 \text{ m}^{-1}$  and a standard deviation of  $0.487 \text{ m}^{-1}$ .

### 3.2.1.3 Yangtze dataset

The Yangtze dataset was collected during a cruise campaign in the Yangtze estuarine and coastal water from 4<sup>th</sup> to 19<sup>th</sup> of March 2013. This dataset consists of 29 measurements of  $R_{rs}$  and  $K_d$  (sampled sites are shown in Figure 3.1c). Radiometric measurements were recorded by the Hyperspectral Surface Acquisition System (HyperSAS, Satlantic Inc., Canada), consisting of the sea-surface radiance ( $L_t$ ), sky radiance ( $L_i$ ) and solar irradiance ( $E_s$ ) from 350 to 858 nm in 136 spectral channels. Details about sensors mounting and observation geometries can be found in Shen et al. (2014).  $R_{rs}$  was calculated by Mobley (1999),

$$R_{rs} = \frac{L_w}{E_s} = \frac{L_t - \rho L_i}{E_s} \quad (3.1)$$

where  $\rho$  is the sea surface reflectance factor which is a function of solar zenith angle and surface roughness. The Ruddick-Hojerslev model (RHM) (Ruddick et al., 2006) is adopted here to compute the  $\rho$  (See Table 6 in Sokoletsky and Shen (2014)).

Vertical profiles of underwater downwelling planar irradiance,  $E_d(\lambda, z)$ , were recorded by the Hyperspectral Ocean Color Radiometer (HyperOCR, Satlantic Inc., Canada), which provides 136 channels of factory calibrated optical data from 350 to 858 nm.  $K_d$  is calculated as the absolute slope of a linear regression fitting between  $\ln[E_d(\lambda, z)]$  and depth  $z$  (Darecki and Stramski, 2004; Kirk, 1994). For the calculation of  $K_d$  in the surface layer,  $E_d(\lambda)$  between 0.6 and 2m with a 0.1m interval are used in the linear regression, and results with determination coefficients larger than 0.99 are accepted. Field measured  $R_{rs}$  and  $K_d$  are presented in Figure 3.3. The  $K_d$  spectrum is smoothed using the adjacent-averaging method (explanation at <http://www.originlab.com/doc/Origin->

Help/Smooth-Algorithm).  $K_d(490)$  in the Yangtze data set ranges from  $0.179 \text{ m}^{-1}$  to  $3.462 \text{ m}^{-1}$  with an average value of  $1.229 \text{ m}^{-1}$  and a standard deviation of  $0.968 \text{ m}^{-1}$ .

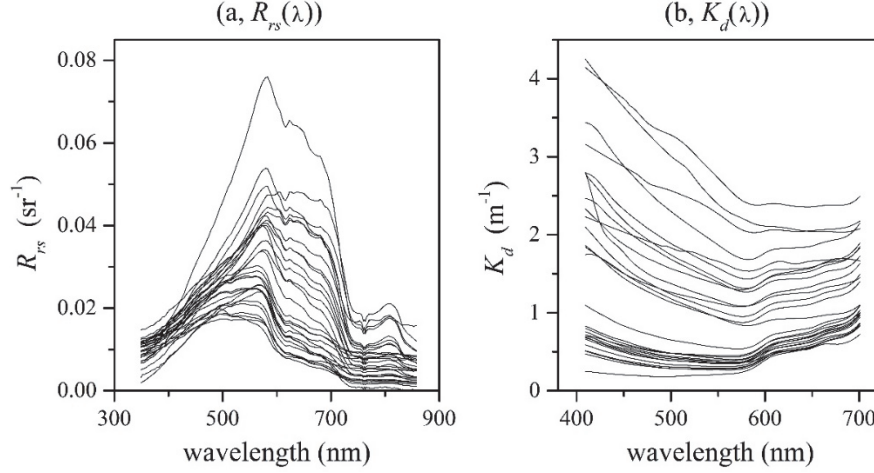


Figure 3.3. Field measured spectra of  $R_{rs}(\lambda)$  and  $K_d(\lambda)$  of the Yangtze dataset.

### 3.2.2 GOCI images and atmospheric correction

GOCI (2010 - present) is a geostationary-orbiting ocean color instrument with a spatial resolution of 500m. It acquires data in eight spectral bands (six visible and two NIR), and eight images are available for the Yangtze Estuary during daytime from 00:15 to 07:45 UTC (Ryu et al., 2012). Table 3-1 briefly introduces the spectral band, bandwidth, and the primary applications for the eight GOCI bands which are centered at 412, 443, 490, 555, 660, 680, 745, and 865nm, respectively. A maximum of three images is free for download before the year 2014 from the official website ([http://hero911.dothost.co.kr/list\\_GOCI.php?dir=/GOCI/1.0](http://hero911.dothost.co.kr/list_GOCI.php?dir=/GOCI/1.0)), while all images are free to access since 2014.

Table 3-1. Introduction of the GOCI 8 spectral bands and their primary application.

Band	Central Wavelengths (nm)	Bandwidth (nm)	Primary Application
1	412	20	Yellow substance and turbidity
2	443	20	Chlorophyll absorption maximum
3	490	20	Chlorophyll and other pigments
4	555	20	Turbidity, suspended sediment
5	660	20	Baseline of fluorescence signal, Chlorophyll, suspended sediment
6	680	10	Atmospheric correction and fluorescence signal
7	745	20	Atmospheric correction and baseline of fluorescence signal
8	865	40	Aerosol optical thickness, vegetation, water vapor reference over the ocean

Six *in-situ* measurements were sampled within 1h of GOCI overpass and selected as matchups in this study. A match-up set is the pair of a 3-by-3 satellite pixel box and an *in-situ* site matching in location and sampling time. The average value of the 3-by-3-pixel box is used as the matched value of  $R_{rs}$  and  $K_d$ . The matched images were acquired on 03:04 UTC and 04:16 UTC of March 4<sup>th</sup> (the matchups sites are A1-1 and A1-2 in Figure 3.1c), 02:16 UTC of March 5<sup>th</sup> (A2-7 in Figure 3.1), 02:16 UTC of March 8<sup>th</sup> (A6-3 and A6-4 in Figure 3.1c) and 04:16 UTC of March 9<sup>th</sup> (A7-2 in Figure 3.1c).

In this study, we corrected the GOCI image using the standard atmospheric correction algorithm (Ahn et al., 2012) which is implemented in GDPS (GOCI Data Processing System) version 1.3. This algorithm is theoretically based on the SeaWiFS method (Gordon and Wang, 1994b) with modifications for turbid waters where the black pixel assumption is invalid (Goyens et al., 2013; Hu et al., 2000; Salama et al., 2004; Salama and Shen, 2010a; Wang and Shi, 2007; Wang et al., 2009b). To improve the accuracy of retrieved water-leaving reflectance in turbid waters, GDPS version 1.3 adopts method of (Hu et al., 2000) to select the nearest non-turbid pixels for deducing NIR water leaving reflectance in turbid waters, and uses the NIR  $R_{rs}$  empirical relationship of (Ruddick et al., 2000) for constraining the AC. A brief algorithm description of GDPS version 1.3 can be found at [http://222.236.46.37/GDPS/ATBD\\_master\\_v9.pdf](http://222.236.46.37/GDPS/ATBD_master_v9.pdf). The GOCI level 2 product (L2P) from GDPS including remote sensing reflectance ( $R_{rs}$ ), which is used in this study to produce maps of diffuse attenuation coefficient.



### 3.2.3 The 2SeaColor model

The 2SeaColor model is based on the solution of the two-stream radiative transfer equations including direct sunlight, as described in Duntley (1942, 1963). Both the analytical forward model and the inversion scheme are provided in detail in Salama and Verhoef (2015). To analytically derive  $K_d$  from  $R_{rs}$ , the first step is to derive the total absorption coefficient ( $a$ ) and backscattering coefficient ( $b_b$ ) from  $R_{rs}$  in the inversion scheme using spectral optimization, after which  $K_d$  is calculated from the derived IOPs using the analytical forward model. Descriptions of the analytical forward model and the inverse scheme are included in Section 2.2.1 and Section 2.2.2, respectively.

### 3.2.4 The Zhang model and the Lee model

To evaluate the performance of the 2SeaColor model, we compare the 2SeaColor model with previous reported empirical and semi-analytical models. The candidate empirical model is the Zhang model (ratio-based) (Zhang and Fell, 2007), and the semi-analytical model is the Lee model (IOPs-based) (Lee et al., 2005a; Lee et al., 2005b).

(i) The Zhang model:

If  $\frac{R_{rs}(490)}{R_{rs}(555)} \geq 0.85$  then

$$K_d(490) = 10^{(-0.843 - 1.459X - 0.101X^2 - 0.811X^3)} + 0.016 \quad (3.2)$$

where  $X = \log_{10}\left(\frac{R_{rs}(490)}{R_{rs}(555)}\right)$  ;

If  $\frac{R_{rs}(490)}{R_{rs}(555)} < 0.85$  then

$$K_d(490) = 10^{(0.094 - 1.302X + 0.247X^2 - 0.021X^3)} + 0.016 \quad (3.3)$$

where  $X = \log_{10}\left(\frac{R_{rs}(490)}{R_{rs}(665)}\right)$  ;

(ii) The Lee model

Total absorption and backscattering coefficient,  $a(\lambda)$  and  $b_b(\lambda)$ , are first derived from remote sensing reflectance by the 2SeaColor model using the spectral optimization.  $K_d(490)$  is then calculated according to the following equation (Lee et al., 2005a),

$$K_d(490) = (1 + 0.005\theta_s)a(490) + 4.18(1 - 0.52e^{-10.8a(490)})b_b(490) \quad (3.4)$$

where  $\theta_s$  is the sun zenith angle.

### 3.2.5 Data analysis

Five statistical parameters are used to quantitatively evaluate the goodness-of-fit between derived and known values. They are the slope and intercept of the regression line, the determination coefficient  $R^2$ , the mean of absolute relative difference (rMAD) and the root mean square error (RMSE). The rMAD and RMSE are defined as follows,

$$\text{rMAD} = \sum |1 - \text{derived}/\text{known}|/N \times 100\% \quad (3.5)$$

$$\text{RMSE} = \sqrt{\sum (\text{derived} - \text{known})^2 / N} \quad (3.6)$$

where  $N$  is the number of observations. Additional parameters such as the percentage of observations within 25% and 100% relative error (F25 and F100, respectively) are also provided.

## 3.3 Results

### 3.3.1 Improvement on the parameterizations

Considering  $R_{rs}$  in the NIR bands ( $>750$  nm) are absent in both the NOMAD and the COASTLOOC datasets, the comparison between the original 2SeaColor and the improved version proposed in this study is made on the Yangtze dataset and the IOCCG (International Ocean Colour Coordinating Group) synthetic dataset. The IOCCG dataset is available at [http://www.ioccg.org/groups/OCAG\\_data.html](http://www.ioccg.org/groups/OCAG_data.html), and the synthetic data for the sun at  $30^\circ$  from zenith is used in this study. The comparison is presented in Figure 3.4. As shown in Figure 3.4a, the original 2SeaColor seems to overestimate  $K_d$  systematically for observations with high  $K_d$ . It is more obvious in Figure 3.4b that the improved version provides more accurate  $K_d$  estimations for the Yangtze dataset compared to the original one.

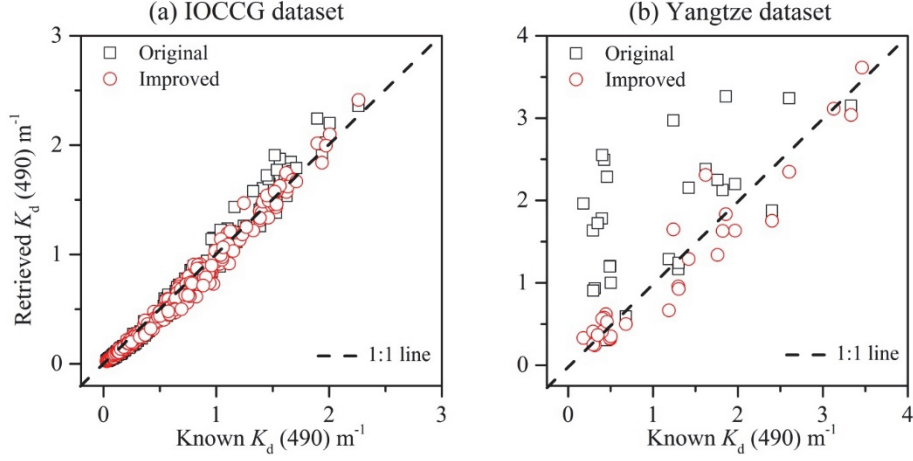


Figure 3.4. Comparison between the original 2SeaColor model and the 2SeaColor model with improved parameterizations for (a) IOCCG dataset and (b) Yangtze dataset.

### 3.3.2 Validation and inter-comparison for $K_d$ retrieval models

Three in-situ datasets described in Section 3.2.3 are used to validate the 2SeaColor model. The performance of the 2SeaColor model is compared with the empirical Zhang model and the semi-analytical Lee model with IOPs derived from the 2SeaColor model. Figure 3.5 shows the known versus derived  $K_d(490)$  for all  $K_d$  retrieval models applying on three in-situ datasets. In order to get a better understanding of these models and their limitations, we divide the three datasets into two subsets for observations from clear and turbid water, respectively. The threshold used in the Zhang model is adopted here, which is corresponding to  $K_d(490) = 0.2 \text{ m}^{-1}$ . This threshold is also considered as an upper limit for simple empirical algorithms (Lee et al., 2013; Lee et al., 2005a).

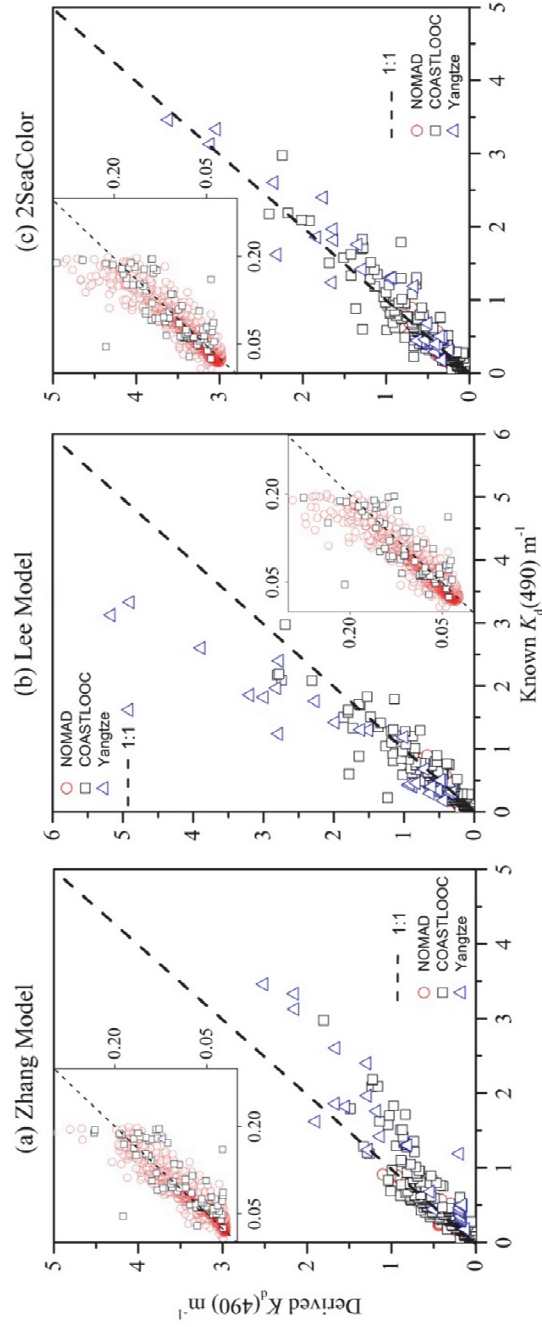


Figure 3.5. Known versus derived  $K_d(490)$  for (a) Zhang model, (b) Lee model and (c) 2SeaColor model, applied to all three data sets. The insets with smaller markers show the results for low  $K_d(490)$  observations where known  $K_d(490) \leq 0.2 \text{ m}^{-1}$ .

The overall performance of the 2SeaColor model is better than that of the other two models with data points well distributed around the 1:1 line (see Figure 3.5c). The Lee model seems to overestimate  $K_d$  in turbid waters as shown in Figure 3.5b, especially for the Yangtze dataset, while the Zhang model seems to systematically underestimate  $K_d$  in turbid waters as shown in Figure 3.5a. From the insets in Figure 3.5, we can find that both the Lee model and the 2SeaColor model generate larger errors for some observations from clear water where  $K_d(490) \leq 0.2 \text{ m}^{-1}$ . Statistical results for the  $K_d(490)$  estimation from all three models are listed in Table 3-2.

As presented in Table 3-2, we can observe that derived  $K_d(490)$  from the 2SeaColor for the full range of observations is most accurate with the highest  $R^2$  ( $= 0.935$ ) and the smallest RMSE ( $= 0.078 \text{ m}^{-1}$ ), compared to the Zhang model ( $R^2 = 0.898$  and RMSE  $= 0.113 \text{ m}^{-1}$ ) and the Lee model ( $R^2 = 0.885$  and RMSE  $= 0.157 \text{ m}^{-1}$ ). For observations from turbid waters where  $K_d(490) > 0.2 \text{ m}^{-1}$ , the 2SeaColor model provides most accurate  $K_d$  estimations, where the RMSE of the Zhang model and the Lee model are 50 % and 109% greater than that of 2SeaColor. The highest  $R^2$  ( $= 0.891$ ) and the smallest rMAD ( $= 21.56\%$ ) are also found for retrievals from the 2SeaColor model for turbid water observations. Predictions from the 2SeaColor are stable with 64.18% of the estimations confined within a 25% relative error margin and 99.29% of the estimations confined within a 100% relative error margin. For clear waters where  $K_d(490) \leq 0.2 \text{ m}^{-1}$ , the Zhang model provides the best estimates with almost unity slope, the smallest rMAD ( $= 21.26\%$ ) and RMSE ( $= 0.017 \text{ m}^{-1}$ ). However, the performance of the 2SeaColor, in this case, are rather comparable in terms of both RMSE ( $= 0.020 \text{ m}^{-1}$ ) and rMAD ( $= 22.58\%$ ).

Table 3-2. Statistical results of three  $K_d$  retrieval models with respect to observations divided into two subsets with a threshold of  $K_d(490) = 0.2 \text{ m}^{-1}$ .

	All (N=1764)			$K_d(490) \leq 0.2 \text{ m}^{-1}$ (N=1481)			$K_d(490) > 0.2 \text{ m}^{-1}$ (N=283)		
	Zhang	Lee	2SeaColor	Zhang	Lee	2SeaColor	Zhang	Lee	2SeaColor
Slope	0.769	1.325	0.929	1.099	1.162	1.125	0.686	1.468	0.908
Intercept	0.014	-0.03	0.012	-0.013	-0.002	-0.001	0.093	-0.203	0.023
$R^2$	0.898	0.885	0.935	0.879	0.845	0.846	0.824	0.835	0.891
RMSE	0.113	0.157	0.078	0.017	0.021	0.020	0.279	0.388	0.186
rMAD	21.78	23.32	22.43	21.26	22.99	22.58	24.52	25.03	21.56
F25	63.27	67.74	67.46	64.55	67.99	68.13	56.38	66.43	64.18
F100	99.72	98.47	99.09	99.66	98.71	99.05	100.00	97.17	99.29

Compared to the 2SeaColor model, the slopes from the linear regression of derived and known  $K_d(490)$  of the Zhang model and the Lee model are far off unity, especially in turbid waters where the slope is more than 30% off unity for the Zhang model and almost 50% off unity for the Lee model. The difference in the slope indicates that predictions from the Zhang model will be underestimated with increasing  $K_d$ , while predictions from the Lee model will be overestimated with increasing  $K_d$ .

### 3.3.3 Estimating $K_d$ from GOCI images

#### 3.3.3.1 Atmospheric correction for GOCI image

To demonstrate the performance of the GOCI standard atmospheric correction algorithm, we compared the field measured  $R_{rs}$  and the  $R_{rs}$  product of GOCI at six matchups sites (Figure 3.6). The GOCI standard atmospheric correction algorithm implemented in GDPS 1.3 generates an acceptable error with RMSE = 0.005  $\text{sr}^{-1}$  and rMAD = 27.30%, which is in accordance with the expected error (Salama and Stein, 2009; Wang et al., 2009b). Note that only seven bands were adopted in the comparison because  $R_{rs}(865)$  is not available in the Yangtze dataset.

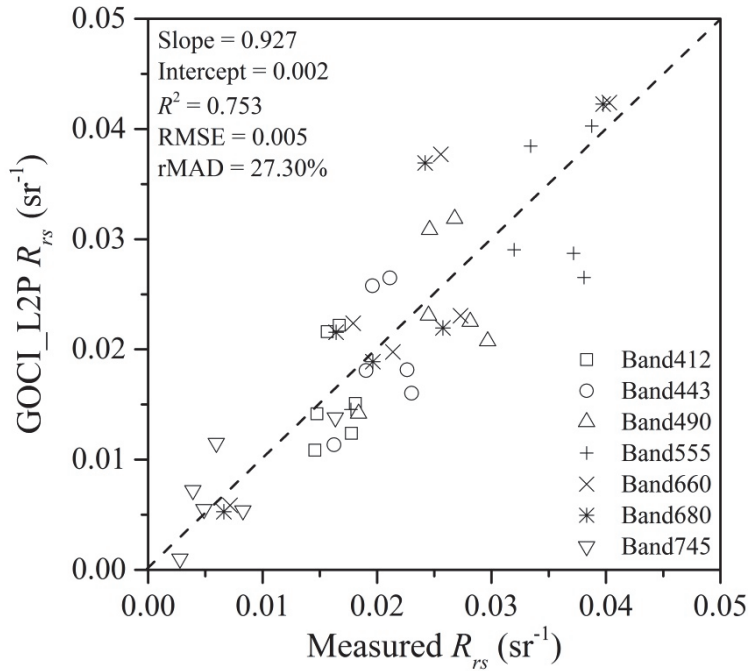


Figure 3.6. Comparisons between field-measured  $R_{rs}$  and GOCI  $R_{rs}$  product for six matchups at each GOCI bands except band 865nm.

It is worthy to mention that the GOCI standard atmospheric correction masks out highly turbid pixels in some of the matched images, even in some of the clear water pixels in this study, which was also reported in previously published results (Hu et al., 2012; Shen et al., 2014). However, developing a more accurate atmospheric correction or recovering the masked pixels (Nordkvist et al., 2009) is out of the scope of this study. The GOCI standard atmospheric correction is considered appropriate for us to focus our study on the retrieval and validation of the diffuse attenuation coefficient.

### *3.3.3.2 Producing $K_d$ maps using the 2SeaColor model*

To better present the distribution of  $K_d$  over the Yangtze estuary, we selected an image with the minimum number of masked pixels from all matched images, which was captured at 02:16 UTC, on 8<sup>th</sup> March 2013. The 2SeaColor model was applied to the L2P product of this image to produce  $K_d$  maps over the Yangtze estuary at wavelengths of 440, 490, 550 and 640nm, respectively. The produced  $K_d$  maps are shown in Figure 3.7.

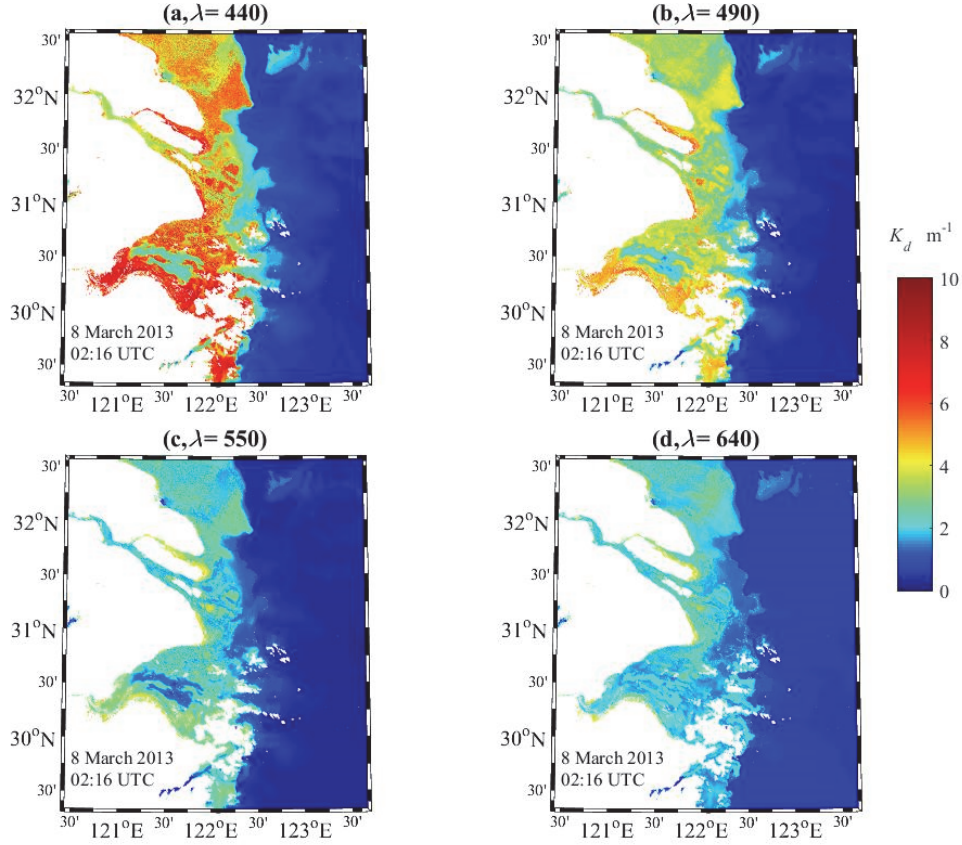


Figure 3.7.  $K_d$  derived from atmospherically corrected GOCI image using the 2SeaColor model. Four wavelengths are shown: (a) 440nm, (b) 490nm, (c) 550nm and (d) 640nm.

The spatial distribution of the diffuse attenuation coefficient at the four chosen wavelengths shows the same pattern. A wide range of  $K_d$  is observed as high  $K_d$  values are present in nearshore waters, especially in Hangzhou Bay and the North Branch of the Yangtze Estuary (see Figure 3.1c), while  $K_d$  in offshore waters is about 1~2 orders of magnitude lower. For instance, in Figure 3.7b,  $K_d(490)$  reaches approximately  $10 \text{ m}^{-1}$  in some pixels in Hangzhou Bay and the North Branch but keeps steadily low at around  $0.1 \text{ m}^{-1}$  in offshore waters. The  $K_d$  maps confirm that both Hangzhou bay and the Yangtze Estuary (especially the North Branch) are extremely turbid waters, which is highly consistent with the known distribution of suspended matter in this area (Shen and Verhoef, 2010; Shen et al., 2010b; Shen et al., 2013; Shen et al., 2014).

From Figure 3.7 we can observe that  $K_d$  decreases from 440nm to 640nm in nearshore waters, but  $K_d$  at 640nm in offshore waters is higher than at other



wavelengths. The difference between the  $K_d$  spectra in clear and turbid waters can also be found in other published studies (Lee et al., 2013) as well as in field measurements in the Yangtze dataset (see Figure 3.3b). The nature of this behavior is the difference in the constituents between clear and turbid waters. For turbid water in the Yangtze estuary, where suspended particulate matter (SPM) and CDOM dominate the inherent optical properties (IOPs), namely a higher absorption coefficient as well as a higher backscattering coefficient at short wavelengths, which leads to a higher  $K_d$  value in this range. For clear water where chlorophyll and water molecules dominate the IOPs, the behavior of the  $K_d$  spectrum is mainly determined by the tremendous increase of the absorption coefficient of water molecules beyond 600nm.

## 3.4 Discussion

### 3.4.1 The improved 2SeaColor model

The improved 2SeaColor model with adopted parameterization provides a better estimation of retrieved  $K_d(490)$ , which proves that the assumption in the original parameterization of the 2SeaColor model is not valid for turbid waters. This is because the absorption coefficient in turbid waters is far more complex than the backscattering coefficient, and the 2SeaColor model will generate uncertainty if no *prior* knowledge of the absorption coefficient is available. Water molecules are not the only absorbing material in the NIR region in turbid waters, as laboratory experiments by Röttgers and Gehnke (2011) showed that natural water samples with high amounts of detrital material still absorb strongly in the NIR region after removing the contribution from the water molecules. Compared to the original 2SeaColor model that requires NIR spectral data, the improved parameterizations will enable the 2SeaColor model to be applied to the datasets without NIR spectral data such as the NOMAD and COASTLOOC datasets.

On the other hand, total absorption and backscattering coefficients ( $a$  and  $b_b$ ) derived from the improved 2SeaColor model are within acceptable uncertainties for both clear and turbid waters. Figure 3.8 shows the validation of  $a(490)$  and  $b_b(490)$  derived from the improved 2SeaColor model using matched observations from the NOMAD and COASTLOOC datasets. For the NOMAD dataset, the accuracy of derived  $a(490)$  from the 2SeaColor model is promising with  $R^2$  of 0.708 and RMSE of 0.708  $\text{m}^{-1}$  (Figure 3.8a). The accuracy of derived  $b_b(490)$  (Figure 3.8c) is also robust with  $R^2$  of 0.558 and RMSE of 0.001  $\text{m}^{-1}$ . As for the COASTLOOC dataset, the comparison between derived  $a(490)$  from the

2SeaColor model and known  $a(490)$  (Figure 3.8b) also shows good accordance with  $R^2$  of 0.794 and RMSE of  $0.201\text{m}^{-1}$ . Note that no matched  $b_b(490)$  is available in the COASTLOOC dataset.

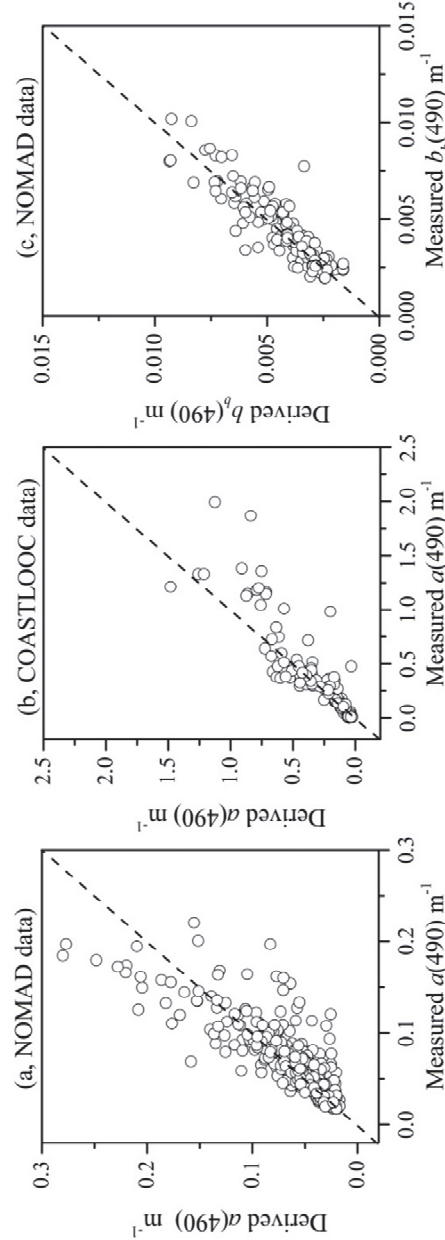


Figure 3.8. Validation of 2SeaColor-derived total absorption and backscattering coefficient at 490nm,  $a(490)$  and  $b_b(490)$ , using matched observations in the NOMAD and COASTLOOC datasets. Note the difference in scales for  $a(490)$  in both cases.

### 3.4.2 Assessments of $K_d$ retrieval models

Empirical models based on simple blue/green  $R_{rs}$  ratios are always deemed unsuitable for turbid coastal waters and considerably underestimate  $K_d$  by a factor of 2 ~ 3 or more in turbid waters (Chen et al., 2014; Sun et al., 2014a; Wang et al., 2009a). This is mainly because remote sensing reflectance in the blue and green bands is mainly affected by detrital material and gelbstoff in turbid waters. These two constituents absorb strongly in the blue region, which makes the water-leaving signal continually low in the blue region. Therefore, variations in the mixture of detrital material and gelbstoff may not affect the blue/green ratios of  $R_{rs}$  but have huge effects on the diffuse attenuation coefficient. Both the Zhang model and the Lee model have often been applied in model inter-comparisons and were validated against various field measurements, showing that they were suitable for  $K_d$  retrieval in both clear and turbid waters (Jamet et al., 2012; Lee et al., 2005b; Zhang and Fell, 2007). Results in this study show that derived  $K_d$  from the 2SeaColor model is comparable with these two models in clear waters, while the 2SeaColor has better performance over highly turbid waters.

The good performance for the Zhang model and the Lee model in clear water can be attributed to the fact that the empirical coefficients in these models have already been fine-tuned to clear waters during model development. For clear waters, where in most cases chlorophyll and water molecules dominate the optical properties, the variability of optical properties is rather limited compared to those of turbid waters. Therefore, the fine-tuned empirical or semi-analytical model should always work well for clear waters. However, when it comes to turbid waters, the empirical models generate large uncertainty since the fine-tuned coefficients are not suitable for optically complex waters, and a fine-tuning of the empirical coefficients would be required (Mitchell et al., 2014; Qiu et al., 2013). This is one of the main reasons that the empirical Zhang model systematically underestimates  $K_d$ , and the semi-analytical Lee model overestimates  $K_d$  in turbid waters.

Additionally,  $K_d$  is considered as sun angle dependent, and the ratio-based models are most likely to represent only the diffuse attenuation at a fixed angle (Lee et al., 2013). As for the three datasets, the percentage distribution (Figure 3.2) shows a large variation in sun zenith angle, the mismatch in sun zenith angle would result in large errors for the Zhang model. As presented in Figure 3.2, most observations were measured at sun zenith angles larger than 40°, especially for the COASTLOOC and the Yangtze datasets (the percentage are 85.71% and

75.86, respectively). A large solar zenith angle usually results in a higher  $K_d$  due primarily to the increased photon path length in water (Lee et al., 2013). The mismatch can be corrected for the IOPs-based Lee model and the 2SeaColor model where the solar zenith angle is used as an input parameter. However, for the Zhang model, the mismatch in sun zenith angle could contribute to the large uncertainty of derived  $K_d$  in turbid waters.

As for the Lee model, the contribution rate of  $b_b$  to  $K_d$  is empirically modeled, which means estimations from the Lee model in turbid waters will also involve larger uncertainties without adapting the empirical coefficients. Another reason for the large uncertainty in turbid waters for the Lee model is that the scattering coefficient in turbid waters tends to become more significant in the forward direction, while the Lee model ignores this forward scattering. The forward scattering may have little effect on the reflectance, but it has a large effect on  $K_d$ .

In contrast, the success of the 2SeaColor for clear and turbid waters should be attributed to both the analytical forward model and its inversion scheme. The analytical forward model takes forward scattering into account by applying the similarity transformation, and estimates  $K_d$  based on radiative transfer physics, while the inversion scheme, with improved parameterizations, can accurately derive the absorption and backscattering coefficients using the spectral optimization technique. In general, the advantages of the 2SeaColor model are: (i) no empirical relationships are involved and a fine-tuning of empirical coefficients is avoided; (ii) no *prior* knowledge on optical properties of the regions of interest is required; (iii) only the spectrum of remote sensing reflectance is necessary; (iv) it provides consistently accurate estimations for both clear and turbid waters. With these advantages, we can state that 2SeaColor has little limitations due to regional variability in optical properties. Thus, 2SeaColor forms an attractive optional solution for the retrieval of  $K_d$  from Earth Observation data for both clear and turbid waters. Furthermore, to produce the best diffuse attenuation coefficient for an image at a large spatial scale covering both clear and turbid waters, a blended algorithm combining the Zhang model for clear waters with the 2SeaColor model for turbid waters would be an ideal approach.

### **3.4.3 Variation of $K_d$ maps in Yangtze Estuary**

We reproduced the  $K_d(490)$  maps from the selected GOCI images acquired on 8 March of 2013 using the Zhang model, the Lee model, and the 2SeaColor model

and the results are presented in Figure 3.9. We can observe that the  $K_d(490)$  maps over the Yangtze Estuary from all three models present the similar patterns in spatial distribution, also shown in previous studies (Shi and Wang, 2010; Shi et al., 2011; Sokoletsky et al., 2014; Wang et al., 2009a). As shown in Figure 3.9, results from all three models in clear waters are comparable, but significant differences can be observed in the extremely turbid regions, such as the Hangzhou Bay and North Branch of Yangtze Estuary. The  $K_d$  map from the 2SeaColor model results in an approximate range from 0.04 to 9.82  $\text{m}^{-1}$  (Figure 3.9c), while the ranges of  $K_d(490)$  from the Zhang model and the Lee model are 0.02 to 3.5  $\text{m}^{-1}$  (Figure 3.9a) and 0.18 to 12.63  $\text{m}^{-1}$  (Figure 3.9b), respectively. The magnitudes of  $K_d(490)$  also vary largely compared to previously reported results for the Yangtze estuary.  $K_d$  maps from Shi et al. (2011) and Shi and Wang (2010) are produced from Moderate Resolution Imaging Spectroradiometer (MODIS-Aqua), yielding a range of derived  $K_d(490)$  approximately from 0 ~ 4  $\text{m}^{-1}$  and 0 ~ 5  $\text{m}^{-1}$ , respectively. The highest  $K_d(490)$  even reaches 250  $\text{m}^{-1}$  when Sokoletsky et al. (2014) simply calculated  $K_d(490)$  as the sum of total absorption and backscattering coefficient.

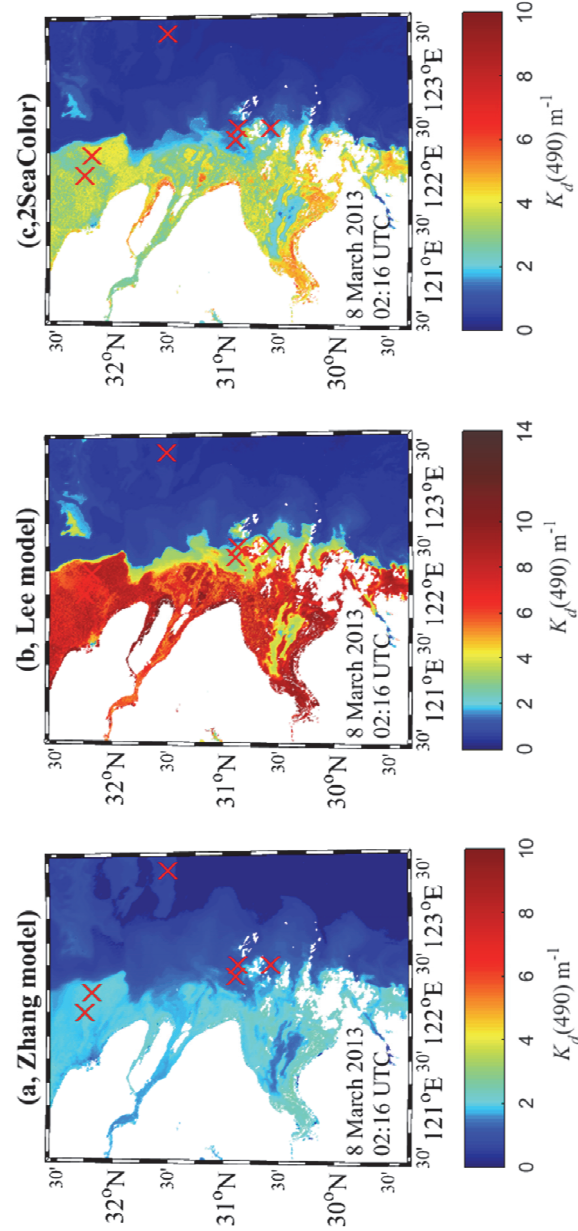


Figure 3.9.  $K_d$  maps derived from the GOCI level 2 product (L2P) at wavelength 490 nm using (a) Zhang model, (b) Lee model and (c) 2SeaColor model. The image was captured at 02:16 UTC, 8 March 2013 and 'X' marks are matchup sites.

We compared the field measured  $K_d(490)$  with GOCI-derived  $K_d(490)$  at six matchup sites for all three models, as presented in Figure 3.10. The results are in

accordance with the comparison from the Yangtze dataset (Figure 3.8c) since the error from the GOCI standard atmospheric correction is acceptable (Figure 3.6). The 2SeaColor model ( $R^2 = 0.937$  and  $RMSE = 0.317 \text{ m}^{-1}$ ) is able to provide accurate and stable estimations, while the Zhang model ( $R^2 = 0.817$  and  $RMSE = 0.640 \text{ m}^{-1}$ ) tends to underestimate  $K_d(490)$  and the Lee model ( $R^2 = 0.927$  and  $RMSE = 2.347 \text{ m}^{-1}$ ) tends to overestimate  $K_d(490)$  for turbid conditions. Note that there are also six data points for the Zhang model in Figure 3.10, but two data points overlapped because both the measured and derived values are extremely close. These two matchup sites are geographically close to each other (A6-3 and A6-4 in Figure 3.1c).

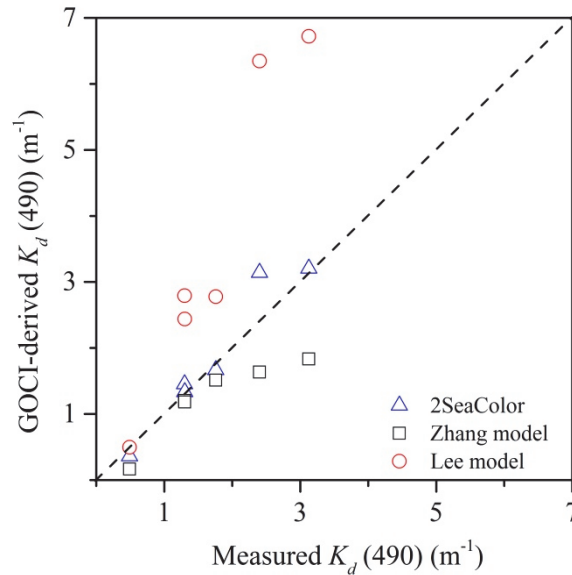


Figure 3.10. Comparison between field-measured  $K_d(490)$  and GOCI-derived  $K_d(490)$  using the Zhang model, the Lee model and the 2SeaColor model for the six matchups.

As we can see from Figure 3.9, all matchup sites are located in clear and medium turbid regions. It is worth mentioning that field measurements in the extremely turbid regions are absent for model validation in this study, making validation of derived  $K_d$  in extremely turbid water impossible in this study. However, The  $K_d(490)$  range predicted by the 2SeaColor model is consistent with field measurements from the East China Sea and Yangtze estuarine waters by with  $K_d(490)$  ranging from 0.03 to 10.3  $\text{m}^{-1}$  with a mean value of  $0.92 \pm 1.59 \text{ m}^{-1}$  (Qiu et al., 2013). In addition, the concentration of suspended particulate matter (SPM) in the Hangzhou Bay and the North Branch of Yangtze Estuary is 1 ~ 2 orders of magnitude greater than the one at the matchup sites (Shen et al., 2010b), which

should result in a relatively large range of  $K_d$ . However, as shown in Figure 3.9a, the difference in  $K_d$  derived from the Zhang model is rather small for pixels in the Hangzhou Bay and the two matchup sites at the top of the image that are considered in medium turbid waters, indicating that predictions from the Zhang model are underestimated in turbid waters.

On the other hand,  $K_d(490)$  derived from the Lee model is above  $6 \text{ m}^{-1}$  in most coastal pixels as shown in Figure 3.9b, which is much larger than the average value reported by Qiu et al. (2013). In accordance with the validation of GOCI-derived  $K_d$  as shown in Figure 3.10, predictions from the Lee model could be systematically overestimated in medium and extremely waters. Therefore, we argue that the GOCI-derived  $K_d$  map from the 2SeaColor model is more reliable, but further validation in extremely turbid waters is still strongly recommended.

### **3.5 Conclusions**

In this study, we proposed improvements on the parameterizations in the inverse scheme of the analytical 2SeaColor model. The 2SeaColor model, with improved parametrizations, can provide stable and accurate estimations of  $K_d$  for wide ranges of water turbidity without the need for tuning coefficients from field measurements. The 2SeaColor model is more appropriate for  $K_d$  retrieval in turbid waters, compared to the empirical Zhang model and the semi-analytical Lee model. Reasonable  $K_d$  maps are also derived when applying the 2SeaColor model to atmospherically corrected GOCI images, reproducing the expected magnitude and range of  $K_d$  over the Yangtze Estuary. The successful application of 2SeaColor to GOCI images enables our further study of variations in the  $K_d$  distribution at large spatial and temporal scales.



## **Chapter 4      Retrieval of wide-range suspended particulate matter concentration from remotely sensed particulate backscattering coefficients\***

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\* This Chapter is based on:

Yu, X., Salama, M.S., Shen, F., & Verhoef, W. Retrieval of wide-range suspended particulate matter concentration from remotely sensed particulate backscattering coefficients in the Yangtze estuary. Preparing submission to *Remote Sensing of Environment*.

## **4.1 Introduction**

Suspended particulate matter (SPM) has long been recognized as a fundamental component in determining coastal ecological and biological functioning, as SPM limits available light in the water column and restrains phytoplankton growth and primary production (Cloern, 1987; Redalje and Fahnenstiel, 1994). Knowledge of the SPM dynamics could be beneficial for a better understanding of the transportation of riverine sediments, heavy metals and particulate organic carbon (Bunt et al., 1999; Chen et al., 2001; Liu et al., 2007). Investigating the variation of SPM could also provide practical suggestions for shipping channels maintenance, offshore constructions, and other administrative services. To better monitor SPM dynamics in coastal regions, remotely sensed imagery could be an efficient approach, which allows detection of SPM with synoptic spatial distributions and with high temporal-spatial resolutions (IOCCG, 2000, 2006).

Many efforts have been carried out to remotely retrieve SPM concentrations ( $C_{\text{SPM}}$ , in mg/L) from satellite images in the coastal waters, and numerous models have been developed empirically (Dogliotti et al., 2015; Zhang et al., 2010) or based on radiative transfer (Doxaran et al., 2006; Shen and Verhoef, 2010; Shen et al., 2014; Volpe et al., 2011). Due to remote sensing reflectance ( $R_{rs}$ , in  $\text{sr}^{-1}$ ) tends to saturate at high SPM concentrations at the visible wavebands (Lee et al., 2016; Nechad et al., 2010; Shen et al., 2010a), most empirically models are calibrated by the least-squares-fit between  $C_{\text{SPM}}$  and the reflectance ratio of NIR and visible bands (Doxaran et al., 2002b, 2003). The calibrated relationships, in the meantime, are also expected to switch from linear in less turbid waters (Binding et al., 2005; Hu et al., 2004) to noticeably curved for high-turbidity waters (Doxaran et al., 2002a; He et al., 2013; Mao et al., 2012).

The validity of  $R_{rs}$ - $C_{\text{SPM}}$  models could be fundamentally attributed to the fact that SPM concentration is one of the major factors that determine the water absorption and scattering capability in aquatic environment, while  $R_{rs}$  is a measure of bulk absorption ( $a$ , in  $\text{m}^{-1}$ ) and backscattering coefficients ( $b_b$ , in  $\text{m}^{-1}$ ) (Lee et al., 2002; Salama and Verhoef, 2015). The correlation between  $C_{\text{SPM}}$  and the water scattering properties is more profound as suspended particles are the major water constituent that scatters the light. For remote sensing applications, SPM is primarily related to the backscattering coefficient of suspended particles ( $b_{bp}$ , in  $\text{m}^{-1}$ ), as the forward scattering signal cannot be directly received by above-water radiometric measurements or satellites (Babin et al., 2003; Neukermans et al., 2012; Sutherland et al., 2000). In many ocean color practices,  $R_{rs}$  at near-infrared

(NIR) bands is expressed simply as a function of  $b_{bp}$ , assuming negligible contributions from other optically-active constituents to the inherent optical properties (IOPs) apart from the absorption by pure water (Chen et al., 2013; Doxaran et al., 2003).

The relationship between  $b_{bp}$  and  $C_{SPM}$  is mostly considered as linear in clear and moderate turbid waters, with the assumption of constant mass-specific backscattering coefficients  $b_{bp}^*$  (the ratio of  $b_{bp}$  to  $C_{SPM}$ ) (Dogliotti et al., 2015; Doxaran et al., 2006; Forget et al., 1999; Knaeps et al., 2015; Neil et al., 2011). However, the  $b_{bp}$ - $C_{SPM}$  relationship in highly turbid waters is less investigated and documented. One of the reasons is probably due to the absence of reliable  $b_{bp}$  measurements. Previous studies have demonstrated that theoretical approaches can improve the retrieval accuracy of IOPs from  $R_{rs}(\lambda)$  in optically complex waters (Brewin et al., 2013; IOCCG, 2000; Yu et al., 2016), which allows us to evaluate the  $b_{bp}$ - $C_{SPM}$  relationship in highly turbid waters with remotely retrieved  $b_{bp}$ .

In this study, we proposed a  $C_{SPM}$  retrieval model based on derived  $b_{bp}$  from  $R_{rs}(\lambda)$  using the inverse scheme of the 2SeaColor model (Salama and Verhoef, 2015; Yu et al., 2016). An index of  $b_{bp}$  based on sigmoid function ( $S_{index}$ ) is introduced to calibrate the  $C_{SPM}$  retrieval model. The proposed model is found to be promising for both clear and turbid waters without the need for a switching scheme. At last, the proposed  $S_{index}$ - $C_{SPM}$  model is applied to atmospherically corrected GOCI (Geostationary Ocean Color Imager, 2010 - present) images to investigate the spatial distribution and diurnal dynamics of SPM in the Yangtze estuary.

## 4.2 Data and methods

Two cruises were carried out in the Yangtze estuary in February 2011 from 24<sup>th</sup> to 25<sup>th</sup>, and in March 2013 from 4<sup>th</sup> to 19<sup>th</sup>, with shipborne measurements of remote sensing reflectance and IOPs. Water samples were collected for laboratory analysis to measure SPM concentration. The sampling sites are presented in Figure 4.1. Note that during the cruise in February 2011, as shown in Figure 4.1b, we followed the route from site C to site D, and then to site E, and repeated the route five times.

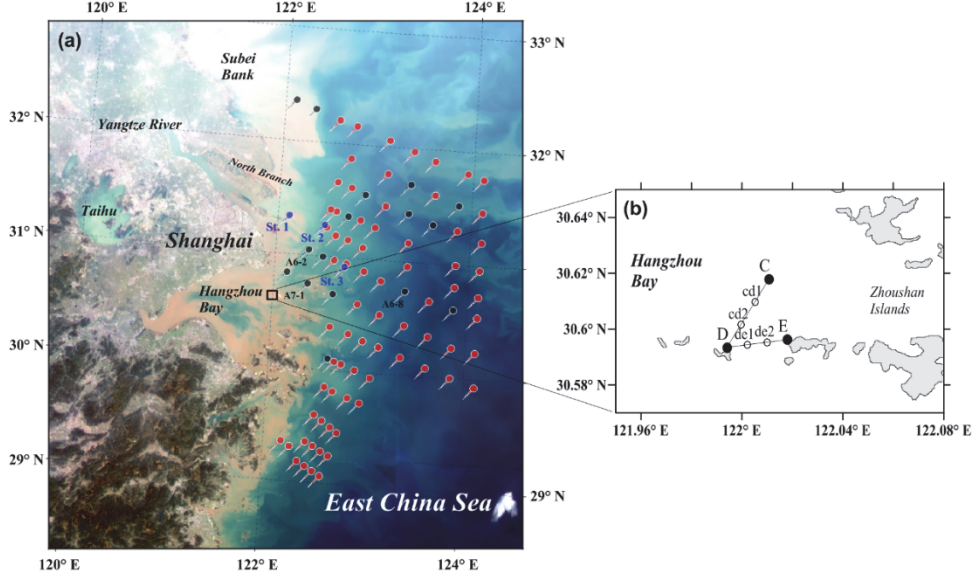


Figure 4.1. Location of the sampling stations for the two cruises in (a) March 2013 and (b) February 2011. The color image in (a) (composited by GOCI bands at 680 nm, 555 nm, and 443 nm) was acquired at 01:16 UTC on March 8<sup>th</sup>, 2013. The black pins in (a) indicate the stations with GOCI match-ups and the blue pins in (a) are the hydrologic stations.

GOCI is a geostationary-orbiting ocean color instrument with a spatial resolution of 500m. It acquires data in 8 spectral bands (6 visible and 2 NIR), and eight images are available for the Yangtze Estuary during daytime from 00:15 UTC to 07:45 UTC (Ryu et al., 2012). The central wavelengths are 412, 443, 490, 555, 660, 680, 745 and 865 nm. The match-up sites are defined as *in-situ* measurements were conducted within one hour of GOCI capturing image, resulting in a total of 16 match-ups. GOCI  $R_{rs}$  matchup is calculated as the averaged  $R_{rs}$  of a 3-by-3 satellite pixel box that matches *in-situ* measurement in location and sampling time.

#### 4.2.1 *In-situ* measurements of IOPs and $R_{rs}$

Radiometric measurements were recorded by two radiance sensors and one irradiance sensor of the Hyperspectral Surface Acquisition System (HyperSAS, Satlantic Inc.). The viewing-illumination geometry of the mounted-sensors was confirmed to the NASA protocol (Fargion and Mueller, 2000). The HyperSAS measures sea-surface radiance ( $L_t$ ), sky radiance ( $L_i$ ) and down-welling solar irradiance ( $E_s$ ) from 350 to 858 nm in 136 spectral channels. The remote sensing reflectance was then calculated by,

$$R_{rs} = (L_t - \rho L_i) / E_s \quad (4.1)$$

where  $\rho$  is the sea surface reflectance factor and was computed according to Ruddick et al. (2006). Figure 4.2 shows the spectra of measured  $R_{rs}$  from the two cruises in February 2011 (33 measurements) and March 2013 (49 measurements).

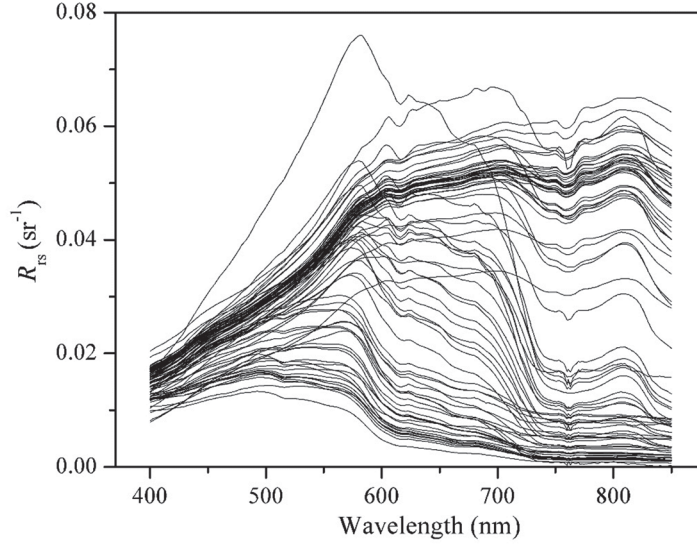


Figure 4.2. Spectra of field-measured  $R_{rs}$  during the two cruises.

In this study, we also employed an independent dataset, namely the SeaSWIR dataset (Knaeps et al., 2018), for model validation. The SeaSWIR data was collected from three turbid estuarine sites of Scheldt (Belgium), Gironde (France), and Río de La Plata (Argentina) with 137 measurements of  $R_{rs}$  from 350 to 1300 nm by ASD (Analytical Spectral Devices, Inc.). SPM concentrations, paired with the  $R_{rs}$  data, were also analyzed and varied between 48.33 and 1400.54 mg/L.

To account for band equivalent effect, all measured  $R_{rs}$  were convoluted to GOCI center bands according to the GOCI relative spectral response function (SRF) by,

$$R_{rs}^{convoluted}(B_i) = \frac{\int R_{rs}(\lambda) SRF_i(\lambda) d\lambda}{\int SRF_i(\lambda) d\lambda} \quad (4.2)$$

where  $SRF_i(\lambda)$  is the spectral response function at band  $i$  ( $B_i$ ) and  $R_{rs}(\lambda)$  is the field measured hyperspectral data. Note that due to the bandwidth of the eighth GOCI band exceed the range of HyperSAS measurement,  $R_{rs}(865)$  of the HyperSAS was empirically calculated from  $R_{rs}(858)$  based on linear relationship from the SeaSWIR data (*i.e.*,  $R_{rs}(865) = 0.97 R_{rs}(858)$ ,  $R^2=1$ ).

During the cruise in March 2013, a package of in-water instruments was deployed to measure the surface IOPs, which are defined as the averaged IOPs in the upper

0.5 m of water depth. Light attenuation and absorption coefficients of the water body, except for water molecular, were measured by the WET Labs ac-s meter at 83 spectral bands from 400 nm to 744 nm. Particulate backscattering coefficient was measured by the WET Labs ECO-bb9 (at wavelengths of 412, 440, 488, 510, 532, 595, 650, 676, and 715 nm, and a scattering angle of 117°).

#### **4.2.2 Field sampling and Lab measurements**

Water samples were collected from the upper 0.5 m, simultaneously with *in-situ* IOPs measurements. SPM concentration samples were filtered through pre-dried, and pre-weighed 0.45  $\mu\text{m}$  pore-size membranes, and then were rinsed three times with milli-Q water to remove dissolved salt before being stored at -20° during the cruise. Membranes were then dried at 60° for 24h and re-weighed in the Lab, with  $C_{\text{SPM}}$  calculated as the difference between the re-weighed particle-loaded membrane and the pre-weighed blank membrane. The measured  $C_{\text{SPM}}$  in surface waters varied from 433.67 to 2068.8 mg/L in February 2011, and from 0.4 to 591 mg/L in March 2013, respectively.

CDOM absorption samples were filtered through 0.22  $\mu\text{m}$  polycarbonate membranes and stored in borosilicate glass vials at -20° during the cruise. The CDOM absorption coefficient ( $a_g(\lambda)$ , in  $\text{m}^{-1}$ ) was measured, at room temperature, using the spectrophotometric method. Details about the shipborne filtration and laboratory measurements can be found in Section 5.2.2.

Prior to the filtration of particulate absorption samples, the 47mm GF/F (Whatman®) glass fiber membrane was pre-soaked in Milli-Q water for 15 mins, and the filtration system was rinsed three times by the Milli-Q water. The volume of seawater filtered varied from 10 ml to 2 L depending on the turbidity of water, which ensured the maximum of the measured optical density would not exceed 0.5 (Röttgers and Gehrke, 2011). The particle-loaded membrane was then moved to the plastic petri dish and immediately frozen in liquid nitrogen. The reference samples were prepared every day following the same procedure whereby seawater was replaced by the Milli-Q water. Immediately back to the lab, the membranes were unfrozen to room temperature in a dark environment. Particulate absorption coefficient ( $a_p(\lambda)$ , in  $\text{m}^{-1}$ ) was then determined by the absorption method (Röttgers and Gehrke, 2011), whereby the membranes were placed inside the 150 mm integrating sphere of the Lambda 1050 UV/VIS spectrophotometer. The scan was conducted between 300 to 950nm with a 1 nm

interval, and  $a_p(\lambda)$  was calculated from the measured optical density of the sample membrane ( $D_s$ ) and the reference membrane ( $D_r$ ) by,

$$a_p(\lambda) = \frac{2.303A}{V\beta} (D_s(\lambda) - D_r(\lambda)) \quad (4.3)$$

where  $A$  (in  $\text{m}^2$ ) is the membrane clearance area,  $V$  (in  $\text{m}^3$ ) is the filtered volume, and  $\beta$  is the path-length amplification factor. The recommended value of  $\beta = 4.5$  was used for  $(D_s(\lambda) - D_r(\lambda)) < 0.1$ . For the cases where  $(D_s(\lambda) - D_r(\lambda)) > 0.1$  (maximum 0.5),  $\beta$  can be calculated by the polynomial function,

$$\begin{aligned} \beta = & 6.475(D_s(\lambda) - D_r(\lambda))^2 \\ & - 6.474(D_s(\lambda) - D_r(\lambda)) + 4.765 \end{aligned} \quad (4.4)$$

#### 4.2.3 Atmospheric correction of GOCI images

Developing an accurate AC algorithm is beyond the scope of this study, the standard atmospheric correction algorithm of Ahn et al. (2012) was therefore adopted. The GOCI standard AC is based on the SeaWiFS method (Gordon and Wang, 1994b) with a modified algorithm for turbid waters (Ruddick et al., 2000). For the extremely turbid waters (generally  $C_{\text{SPM}} > 100$  mg/L), the GOCI standard AC iteratively restores the normalized water-leaving reflectance at NIR bands ( $\rho_{\text{wn}}(\text{NIR})$ ) empirically from the red band ( $\rho_{\text{wn}}(\text{Red})$ ). However, the CAL/VAL activities of GOCI are mainly performed in the Korean coastal waters where the  $C_{\text{SPM}}$  is much lower than  $C_{\text{SPM}}$  in the Yangtze estuary, results in larger errors when applying the empirical relationships between  $\rho_{\text{wn}}(\text{NIR})$  and  $\rho_{\text{wn}}(\text{Red})$  to the Yangtze estuary. Figure 4.3 presents the validation of GOCI-derived  $R_{\text{rs}}$  against the field measurements for match-up pairs, resulting rMAD of 31.6%, 25.5%, 18.7%, 13%, 19%, 21.9%, 54.6%, and 60.7% for GOCI  $R_{\text{rs}}$  at 412, 443, 490, 555, 660, 680, 745, and 865 nm, respectively.

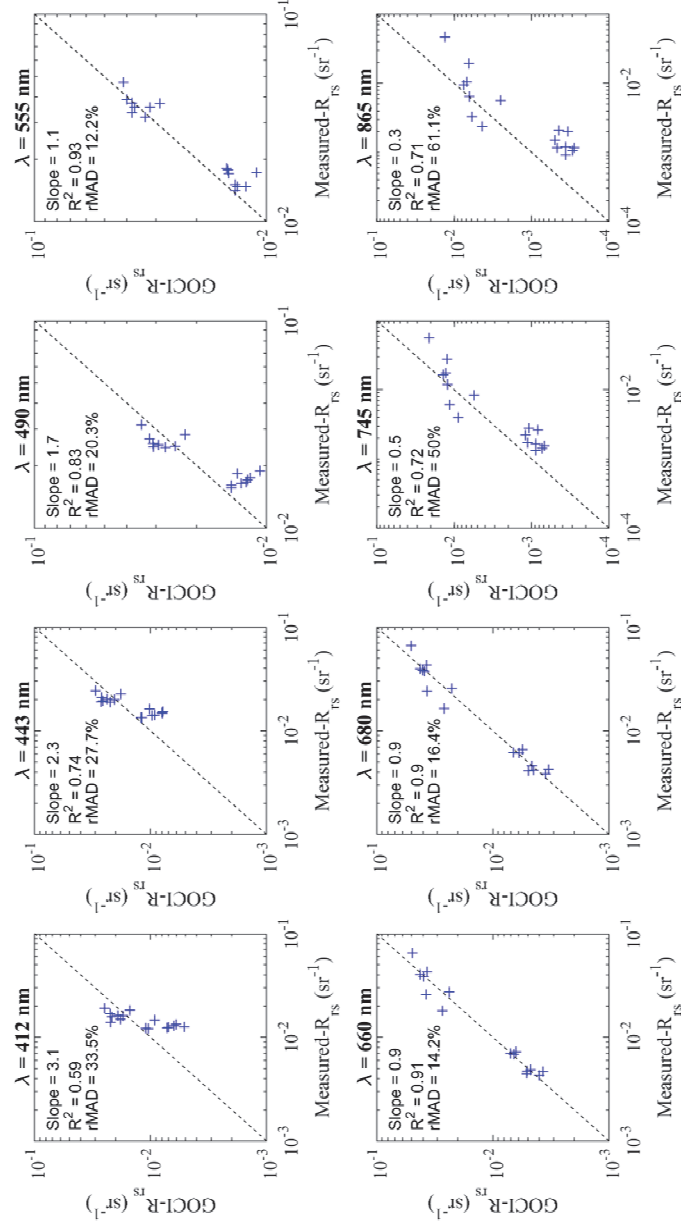


Figure 4.3. Validation of  $R_{rs}$  product of the GOCI standard AC against matched *in-situ* measurements for GOCI 8 wavebands.

The GOCI standard AC could provide satisfactory reflectance results at blue to red bands but has larger uncertainties  $R_{rs}$  at NIR bands (745 nm and 865 nm), especially for extremely turbid waters in the Yangtze estuary. The underestimation of GOCI-derived  $R_{rs}$ (NIR) is also reported by Pan et al. (2017).



Figure 4.4 shows an example of  $R_{rs}$  mapping products from GOCI standard AC at the eight GOCI center bands. Compared with field-measured  $R_{rs}$  spectra (Figure 4.2), GOCI  $R_{rs}$  at the two NIR bands are underestimated for extremely turbid waters, e.g., the Hangzhou Bay.

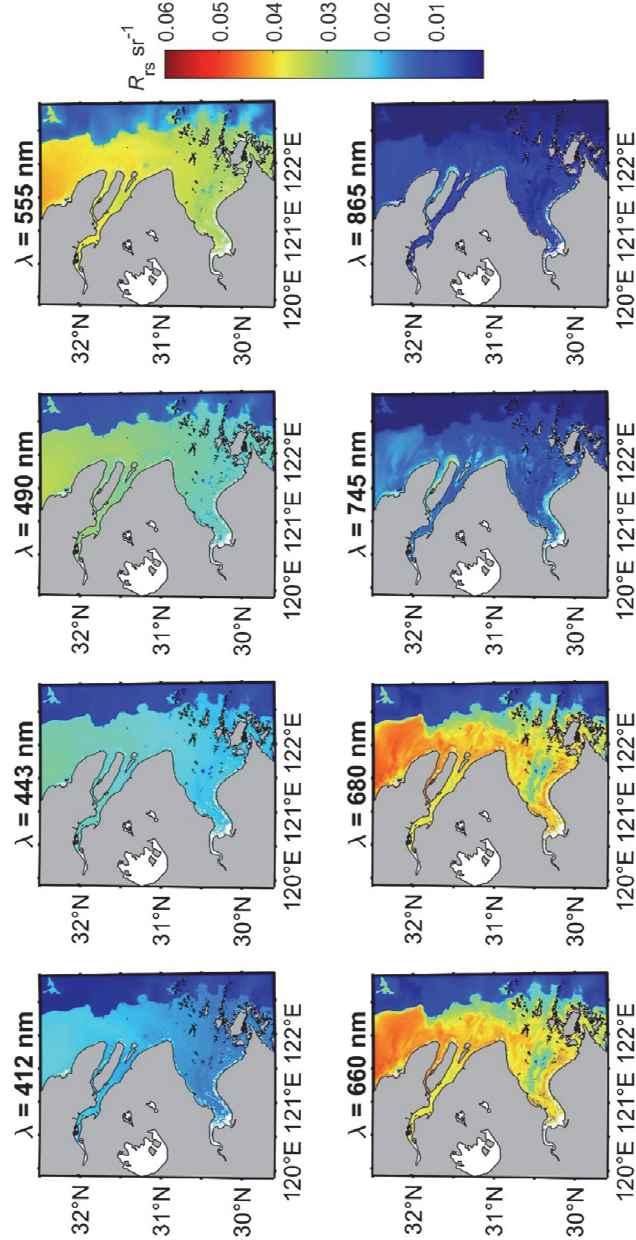


Figure 4.4.  $R_{rs}$  mapping products of GOCI standard AC over the Yangtze estuary for the image acquired on 02:16 UTC, March 8 of 2013.

To address the significant underestimation of the  $R_{rs}$ (NIR) by the GOCI standard AC in extremely turbid waters, we followed the method proposed by Zhang et al. (2010) and adopted regionally calibrated relationships between  $R_{rs}$ (NIR) and  $R_{rs}$ (660) to correct the  $R_{rs}$ (NIR) of GOCI level 2 product (L2P) for extremely turbid pixels. A threshold of  $R_{rs}$ (660) = 0.04 sr<sup>-1</sup> was selected to identify the extremely turbid pixels ( $C_{SPM} > 100$  mg/L) according to field measurements, as  $R_{rs}$ (660) of all turbid samples were found beyond 0.04 sr<sup>-1</sup>. As a result,  $R_{rs}$ (745) and  $R_{rs}$ (865) of extremely turbid pixels were corrected by,

$$R_{rs}(745) = \sum_{n=0}^2 j_n * R_{rs}(660)^n \quad (4.5)$$

$$R_{rs}(865) = \sum_{n=0}^2 k_n * R_{rs}(745)^n \quad (4.6)$$

where  $j_n$  and  $k_n$  are empirical constants that calculated by the least-square fits between field-measured  $R_{rs}$ (660),  $R_{rs}$ (745), and  $R_{rs}$ (865), with  $R^2$  of 0.78 and 0.97 for Eq. (4.5) and Eq. (4.6), respectively. The best-fit values of  $j_{0-2}$  and  $k_{0-2}$  are 20.216, -0.705, 0.033, 9.005, 0.393, and 0.004, respectively.

#### 4.2.4 Retrieval models for SPM concentration

For a proper calibration and validation of the  $C_{SPM}$  retrieval models, we divided  $C_{SPM}$  and matched  $b_{bp}(555)$  data pairs into two optimal calibration and validation sets based on the Geo Cal/Val model (Salama et al., 2012), which ensured the selected Cal and Val datasets with same moments (mean and standard deviation). The Cal dataset includes 50 paired  $b_{bp}(555)$  and  $C_{SPM}$  data with  $C_{SPM}$  ranging between 0.4 and 2068.8 mg/L, while the Val dataset includes 50 data pairs with  $C_{SPM}$  between 0.7 and 1865.8mg/L. Note that both field-measured  $b_{bp}(555)$  and  $R_{rs}$ -derived  $b_{bp}(555)$  were used to maximize the number of paired measurements. Since the Val dataset is independent of the Cal dataset, validation results of the calibrated  $C_{SPM}$  retrieval models are expected to be convincing.

##### 4.2.4.1 The He model

The empirical model proposed by He et al. (2013), which was previously calibrated and validated in the Yangtze estuary, is also employed in this study for comparison. This model, hereafter denoted as the He model, is expressed as,

$$C_{SPM} = 10^{s_1 + s_2(R_{rs}(745)/R_{rs}(490))} \quad (4.7)$$

where  $s_1$  and  $s_2$  are empirical constants. These two constants were optimized by the nonlinear least-square fit using the Cal dataset. The optimized  $s_1$  and  $s_2$  are 1.137 and 1.080, respectively.

#### 4.2.4.2 The GOCI standard SPM model

The GOCI standard SPM model is described in the GDPS (GOCI Data Processing System, version 1.3) algorithm description (available at [http://222.236.46.37/GDPS/ATBD\\_master\\_v9.pdf](http://222.236.46.37/GDPS/ATBD_master_v9.pdf)), where two models are provided for relatively clear and extremely turbid waters. The same switching scheme as the one used in the atmospheric correction (described in Section 4.2.3) is adopted here. The two models are expressed as,

(i) For relatively clear waters ( $R_{rs}(660) < 0.04 \text{ sr}^{-1}$ ),

$$C_{SPM} = 10^{c_0 + c_1(R_{rs}(555) + R_{rs}(670)) - c_2(R_{rs}(490)/R_{rs}(555))} \quad (4.8)$$

(ii) For extremely turbid waters ( $R_{rs}(660) \geq 0.04 \text{ sr}^{-1}$ ),

$$C_{SPM} = 10^{c_3 + c_4(R_{rs}(745)/R_{rs}(555)) + c_5(R_{rs}(680)/R_{rs}(490))} \quad (4.9)$$

where  $c_{0-5}$  are empirical coefficients. In this study, these empirical constants were re-calibrated using the Cal dataset, and the best-fit values are 0.586, 13.503, 0.659, 1.77, 1.847, and -0.558 for  $c_{0-5}$ , respectively.

#### 4.2.4.3 The $b_{bp}$ - $C_{SPM}$ model

In this study, we propose a novel model to estimate  $C_{SPM}$  from remotely sensed  $b_{bp}$  at 555 nm, where  $b_{bp}(555)$  is first derived from  $R_{rs}$ , and then convert to  $C_{SPM}$  using a locally calibrated model. The proposed model employs a power-law function between  $C_{SPM}$  and an index of  $b_{bp}(555)$ . Model calibration and validation are described in Section 4.2.4.2.

The 2SeaColor model is employed to remotely retrieve  $b_{bp}$ , which includes an analytical forward model based on the solution of the two-stream radiative transfer equations and an inversion scheme (Salama and Verhoef, 2015). The reflectance results predicted by the 2SeaColor model are the infinite reflectance ( $r_\infty$ ) and the directional-hemispherical reflectance ( $r_{sd}^\infty$ ) of the semi-infinite medium, which are a function of  $a$  and  $b_b$  (see. Eq. (2.14) and Eq. (2.15)). Details regarding the 2SeaColor model can be found in Section 2.2. The  $b_{bp}$  is retrieved by the 2SeaColor model with improved parameterizations using spectral optimization (Yu et al., 2016). See Section 2.2.2 for the improved parameterization.

#### 4.2.5 Scattering correction for particulate absorption measured by the ac-s meter

Scattering issue contributes a large part to the uncertainties in measured absorption by the ac-s meter, because the instrument fails to collect all the scattered light, resulting in overestimations of measured absorption. Therefore, correction for the scattering effects in ac-s measurements is essential. However, the contribution by scattering effects could vary widely from location to location due to the natural variability of SPM (e.g., particle size, type, and concentration). In this study, three correction methods were performed and validated with laboratory measurements to ascertain the most appropriate method for the Yangtze estuary. These methods are,

- (i) The *proportional* method (Petzold, 1972),

$$a_{corr}(\lambda) = a_m(\lambda) - a_m(715) \frac{c_m(\lambda) - c_m(715)}{c_m(715) - a_m(715)} \quad (4.10)$$

where subscript  $m$  indicates ac-s measured data and  $a_{corr}$  is the corrected absorption coefficient.

- (ii) The *semi-analytical* method (Röttgers et al., 2013),

$$a_{corr}(\lambda) = a_m(\lambda) - [a_m(715) - a_{emp}(715)] \frac{(1/e_c)c_m(\lambda) - c_m(715)}{(1/e_c)c_m(715) - a_m(715)} \quad (4.11)$$

where  $a_{emp}(715) = 0.212 a_m(715)^{1.135}$  and  $e_c = 0.56$  (Boss et al., 2009).

- (iii) The *Boss\_13* method (Boss et al., 2013),

$$a_{corr}(\lambda) = a_m(\lambda) - \varepsilon(c_m(\lambda) - a_{corr}(\lambda)) \quad (4.12)$$

where Eq. (4.12) is solved by non-linear-squares fit with the assumption of  $a_{corr}(\lambda)$  follows  $a_{corr}(\lambda) = a_{corr}(440)e^{-S(\lambda-440)}$ , which is an over-constrained problem with three unknowns ( $a_{corr}(440)$ ,  $S$  and  $\varepsilon$ ) and 83 spectral measurements of  $a_m(\lambda)$  and  $c_m(\lambda)$  by the ac-s.

#### 4.2.6 Data analysis

To quantitatively evaluate the goodness-of-fit between derived and known values, we calculated the slope and intercept of the regression line and the determination coefficient  $R^2$  according to the model-II regression (Laws, 1997). The mean of absolute relative difference (rMAD) and the root mean square error (RMSE) were computed by Eq. (3.5) and Eq. (3.6).

## 4.3 Results

### 4.3.1 Quality assurance of *in-situ* measured IOPs

#### 4.3.1.1 Scattering correction of *ac-s* measurements

To evaluate the performance of the three scattering correction methods, comparisons between scattering corrected *ac-s* measured absorptions and laboratory-measured absorption were made at 440 nm and 650 nm, and the results are presented in Figure 4.5. The lab-measured total absorption coefficient is calculated as the combination of absorption by suspended particles, CDOM, and pure seawater. The absorption coefficients of pure seawater from 380 to 700 nm were taken from Pope and Fry (1997).

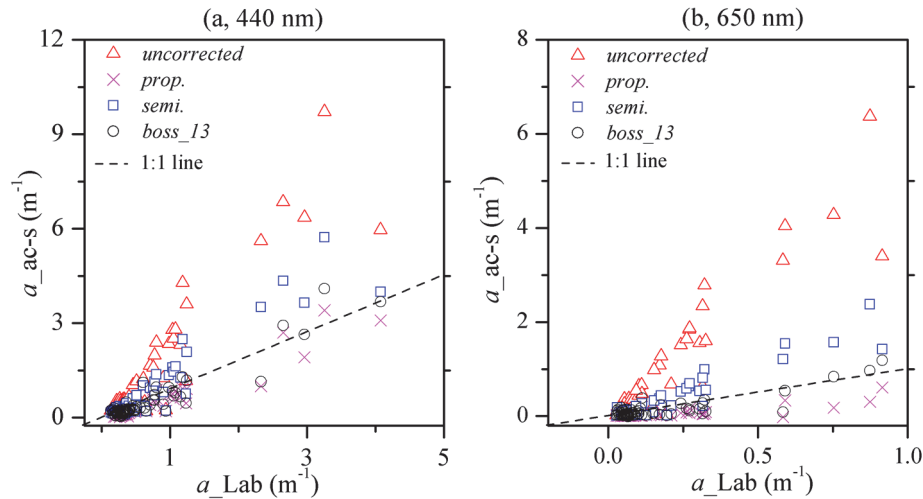


Figure 4.5. Comparison between laboratory and field measured absorption coefficients (uncorrected and corrected results by the three methods) at 440 nm (a) and 650 nm (b).

As shown in Figure 4.5, *ac-s* measurements without scattering correction overestimated the absorption coefficient significantly compared to the laboratory measurements at both 440 and 650 nm, which justified the need for scattering correction for *ac-s* measurements. From the three correction methods, the Boss\_13 model provides the most comparable results with laboratory measurements at both bands with all data points being well distributed alongside the 1:1 line. Whereas the proportion method underestimates the absorption at 650 nm and the semi-empirical model overestimates the absorption at both 550 and 650 nm, especially for the turbid waters. Table shows the statistic results of the three correction methods, where almost unity slopes and the smallest root mean

square errors (RMSE) were derived from the Boss\_13 model for both  $a(440)$  and  $a(650)$ . Therefore, the Boss\_13 method was adopted as the scattering correction method for the WET Labs ac-s meter in this study.

Table 4-1. Statistic results of scattering corrected ac-s measured absorptions by the three methods compared with laboratory measurements at 440 and 650 nm.

	$a(440)$				$a(650)$			
	<i>uncorrected</i>	<i>prop</i>	<i>semi</i>	<i>Boss13</i>	<i>uncorrected</i>	<i>prop</i>	<i>semi</i>	<i>Boss13</i>
<b>Slope</b>	2.478	0.886	1.484	1.030	6.374	0.506	2.364	1.133
<b>Intercept</b>	-0.413	-0.171	-0.215	-0.164	-0.245	-0.031	-0.087	-0.050
<b><math>R^2</math></b>	0.860	0.886	0.886	0.890	0.881	0.620	0.895	0.751
<b>RMSE</b>	1.240	0.343	0.462	0.284	1.100	0.157	0.280	0.105

#### 4.3.1.2 Inter-comparison of in-situ measured $b_{bp}$

Two concurrent measurements of  $b_{bp}$ , by the Wet labs ECO-bb9 and ECO-VSF, were cross-compared for the surface layer data, where a good agreement was observed, covering a wide range of  $b_{bp}$  from 0.02 to 2.43  $\text{m}^{-1}$  from the ECO-bb9 measurements (Figure 4.6). The robust linear relationship (slope = 1.005,  $R^2 = 0.99$ ) suggests an insignificant difference between the two measurements and confirms the good data quality of the field measured  $b_{bp}$ . Considering ECO-bb9 is more widely used by ocean color community and has more spectral bands available,  $b_{bp}$  data of the ECO-bb9 were employed for model calibration and validation in this study.

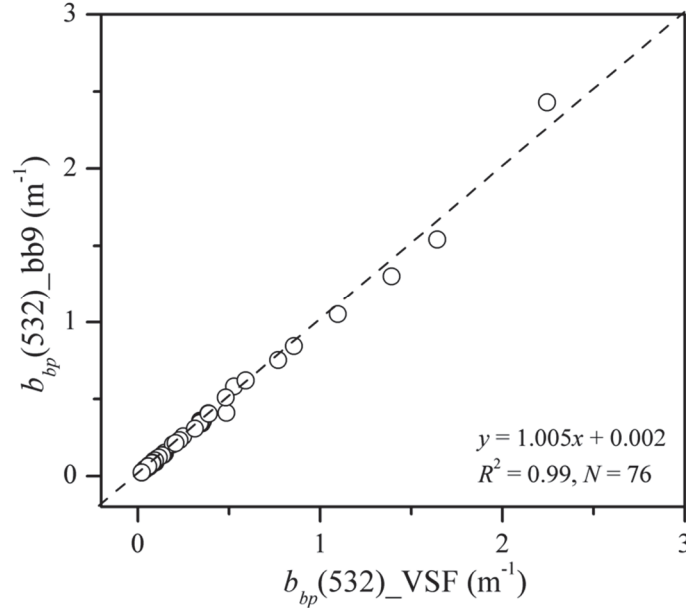


Figure 4.6. Cross-comparison between field-measured  $b_b$  by the ECO-VSF and ECO-bb9 for the surface data.

#### 4.3.2 Validation of retrieved IOPs by the 2SeaColor model

Figure 4.7 presents the comparison between retrieved IOPs from the 2SeaColor model and field-measured IOPs, take  $a(440)$  and  $b_b(532)$  for example. Note that pure seawater absorption and backscattering coefficients were added to the ac-s and ECO-bb9 measurements to compute the measured  $a(440)$  and  $b_b(532)$ . As shown in Figure 4.7, satisfactory retrievals are obtained with RMSEs of  $0.25 \text{ m}^{-1}$  and  $0.10 \text{ m}^{-1}$  for retrieved  $a(440)$  and  $b_b(532)$ , respectively. The slopes and the determination coefficients calculated from the linear regression also suggest the fairly accurate estimations by the 2SeaColor model. Note that the validation results of derived IOPs in this study are consistent with previous validation efforts of the 2SeaColor model for the COASTLOOC and NOMAD datasets (Figure 3.8).

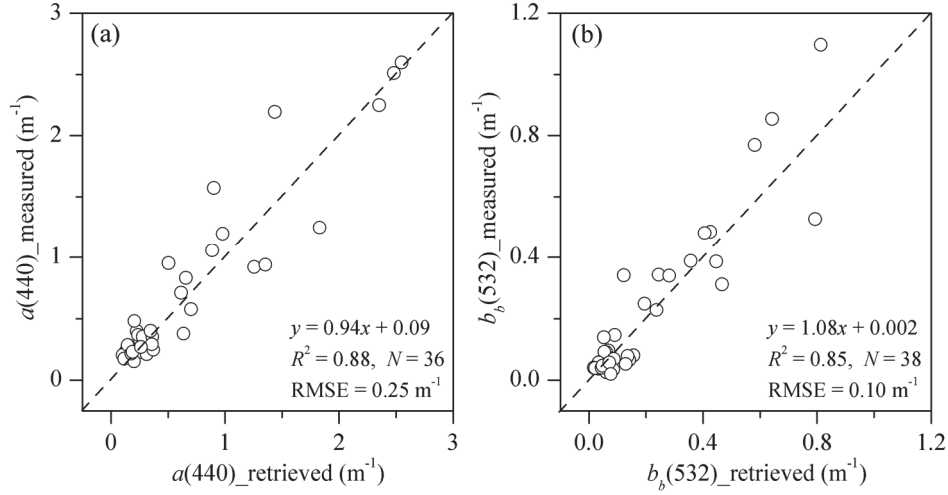


Figure 4.7. Validation of retrieved IOPs by the 2SeaColor model for (a) total absorption coefficient at 440 nm, and (b) total backscattering coefficient at 532 nm.

#### 4.3.3 Cal/Val of the SPM retrieval model

Figure 4.8 shows the scatter plots, in log scale, between  $b_{bp}(555)$  and  $C_{\text{SPM}}$ . It is demonstrated in Figure 4.8 that  $C_{\text{SPM}}$  is linearly correlated to  $b_{bp}(555)$  in relatively clear waters and their relationship becomes significantly curved in extremely turbid waters. The resultant relationships between  $b_{bp}(555)$  and  $C_{\text{SPM}}$  from Figure 4.8 for relatively clear and extremely turbid waters are formulated as,

$$C_{\text{SPM}} = 59.83 * b_{bp}(555) \quad (4.13)$$

$$C_{\text{SPM}} = 84.77 * b_{bp}(555)^{1.696} \quad (4.14)$$

Due to the wide dynamic range of measured  $C_{\text{SPM}}$  (0.4 ~ 2068.8 mg/L), we found that using a single function, linear or non-linear, to fit  $b_{bp}(555)$  and  $C_{\text{SPM}}$  will inevitably result in overestimations of  $C_{\text{SPM}}$  for clear waters or underestimations for extremely turbid waters. However, using two respective relationships to calibrate the  $b_{bp}(555)$ - $C_{\text{SPM}}$  model, such as a linear function for relatively clear waters and a power-law function for extremely turbid waters, requires first a robust switching scheme and could result in discontinuities in the estimated SPM (Shen et al., 2013).



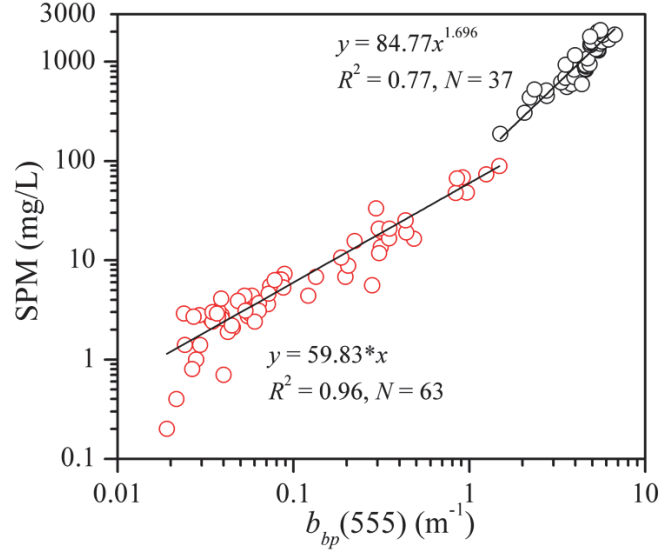


Figure 4.8. Relationship between  $b_{bp}(555)$  and  $C_{SPM}$ . The red circles mark samples with  $C_{SPM}$  less than 100 mg/L, while the black markers represent samples with  $C_{SPM} > 100$  mg/L.

To account for the wide-range of  $C_{SPM}$  and the turbidity-dependent relationship between  $b_{bp}(555)$  and  $C_{SPM}$ , we introduce a  $b_{bp}$  index based on sigmoid function (hereafter denoted as  $S_{index}$ ) to retrieve  $C_{SPM}$ . The  $S_{index}$  is defined as,

$$S_{index} = b_{bp} / (1 + \max(b_{bp}) - b_{bp}) \quad (4.15)$$

where the  $\max(b_{bp})$  represents the maximum  $b_{bp}(555)$  of the calibration dataset. The  $\max(b_{bp})$  from our dataset is about  $6.5 \text{ m}^{-1}$  corresponding to an  $C_{SPM}$  of 2000 mg/L. In this study, we select a maximum  $b_{bp}$  of  $10 \text{ m}^{-1}$  to cover the wide range of  $C_{SPM}$  in the Yangtze estuary. A detailed discussion is given in Section 4.4.2.1.

Figure 4.9a presents the scatter plots between  $C_{SPM}$  and the  $b_{bp}$  index ( $S_{index}$ ) for the Cal dataset, and a power-law function is employed to calibrate the  $C_{SPM}$ - $S_{index}$  relationship, which gives,

$$C_{SPM} = a * S_{index}^b \quad (4.16)$$

where  $a$  and  $b$  are empirical constants and derived as 1463.4 and 1.15 from the non-linear-squares fit with a determination coefficient ( $R^2$ ) of 0.98.

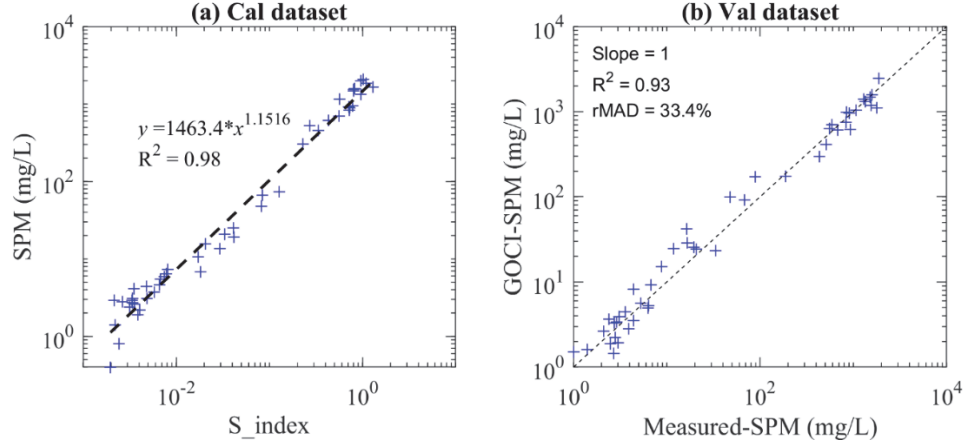


Figure 4.9. Calibration and validation of the  $S_{index}$ - $C_{SPM}$  model using the Cal and Val datasets.

To evaluate the performance of proposed  $S_{index}$ - $C_{SPM}$  model, the calibrated Eq. (4.16) was first applied to the Val dataset. As shown in Figure 4.9b, derived  $C_{SPM}$  matches well with measured  $C_{SPM}$  with an  $R^2$  of 0.93 and an rMAD of 33.4%. Comparison between our proposed model and the He model and the GOCI standard SPM model are presented in Figure 4.10. All three models have comparable performance for the extremely turbid waters, while for the relatively clear waters, the proposed  $S_{index}$ - $C_{SPM}$  model seems to be better in the retrieval accuracy.

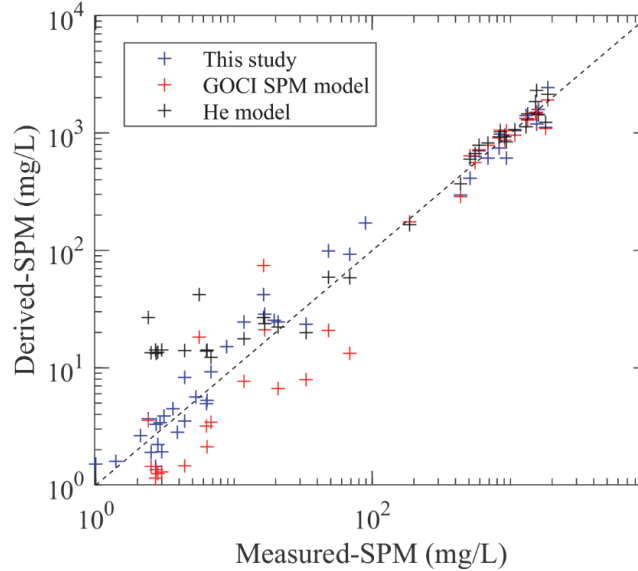


Figure 4.10. Validation results of the three  $C_{SPM}$  retrieval models using the Val dataset.

Statistical data derived from the three models are tabulated in Table 4-2. Considering the He model was mainly developed for turbid waters, evaluation of the three models was also made for extremely turbid samples with  $C_{SPM} > 100$  mg/L. As shown in Table 4-2, validation results show that our proposed model has a better retrieval accuracy of the whole  $C_{SPM}$  range with rMAD of 33.45% compared to rMAD of 50.2% for the GOCI SPM model and rMAD of 185.2% for the He model. The uncertainties of the He model are mainly due to the overestimation of  $C_{SPM}$  for less turbid waters, where an rMAD of 240.58% is calculated. As for the extremely turbid waters, the three models show comparable retrieval accuracy with the GOCI SPM model has the smallest rMAD of 12.7%.

Table 4-2. Statistical data of the three SPM retrieval models that support Figure 4.10.

	SPM < 100 mg/L			SPM > 100 mg/L			All		
	Slope	$R^2$	rMAD (%)	Slope	$R^2$	rMAD (%)	Slope	$R^2$	rMAD (%)
GOCI SPM model	0.94	0.08	85.70	0.89	0.85	12.72	0.95	0.95	50.20
He model	0.81	0.66	344.35	1.14	0.79	17.22	1.09	0.93	185.20
This study	1.80	0.93	44.72	1.10	0.79	15.29	0.99	0.93	33.45

#### 4.3.4 Mapping SPM concentrations in the Yangtze estuary

To map the spatial distribution of  $C_{SPM}$  in the Yangtze estuary, the proposed  $S_{index}$ - $C_{SPM}$  model was applied to atmospherically corrected GOCI image acquired at 02:16 UTC, March 8<sup>th</sup> of 2013, as shown in Figure 4.11a. SPM mapping products of the He model (Figure 4.11b) and the GOCI standard SPM model (Figure 4.11c) were also presented for proper comparison.

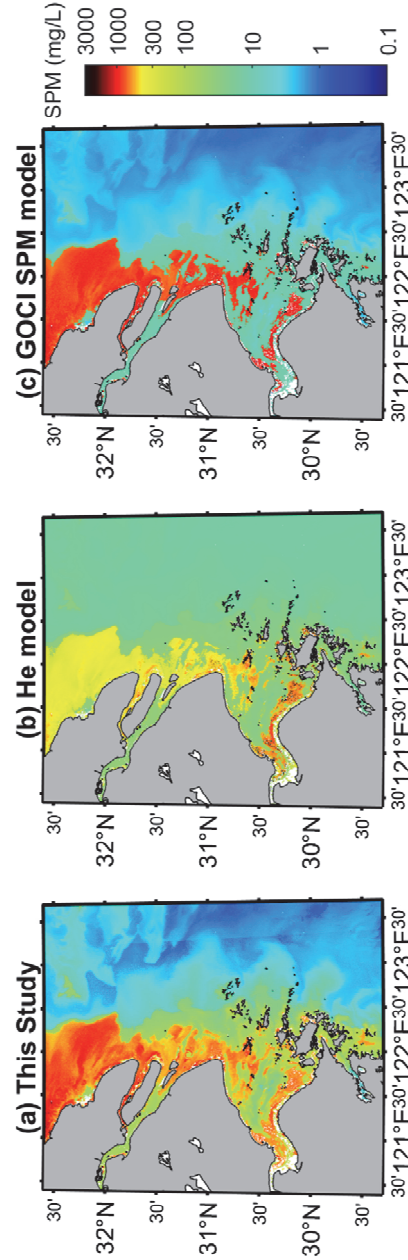


Figure 4.11. SPM Maps over the Yangtze estuary produced by (a) our proposed model, (b) the He model, and (c) the GOCI SPM model. The GOCI image was acquired at 02:16 UTC, March 8<sup>th</sup>, 2013.

As shown in Figure 4.11, the produced  $C_{SPM}$  maps by the three models are consistent in the distribution pattern and spatial gradient, *i.e.*, a decreasing trend of  $C_{SPM}$  from inshore to offshore waters and extremely turbid waters mainly exist

in the Hangzhou Bay, the North Branch, and the Subei bank. However, there also shows substantial differences in terms of the dynamic ranges of retrieved  $C_{SPM}$  from the three models, which allows immediate evaluation of the models' performance. As shown in Figure 4.11a and Figure 4.11c, our proposed model has comparable  $C_{SPM}$  range with the output of the GOCI SPM model for both clear and turbid waters, while the He model seems to less reasonable in terms of the produced  $C_{SPM}$  range. With field-measured  $C_{SPM}$  in some of the Yangtze estuary regions reported over 1000mg/L in winter, it is reasonable to assume that our proposed model and the GOCI standard model are closer to catching the actual variabilities of  $C_{SPM}$ .

To better evaluate the performance of these three models, retrieved  $C_{SPM}$  was validated with field-measured  $C_{SPM}$  for the match-ups (the black pins in Figure 4.1a), and the result is presented in Figure 4.12. The statistical data show that the proposed  $S_{index}$ - $C_{SPM}$  model provides an overall better estimation of  $C_{SPM}$ , especially for the match-ups with  $C_{SPM} > 10$  mg/L. rMAD of the retrieved  $C_{SPM}$  from the  $S_{index}$ - $C_{SPM}$  model is 62.3 %, which is improved with respect to the GOCI standard  $C_{SPM}$  model (rMAD = 74.9%) and the He model (rMAD = 180.2%).

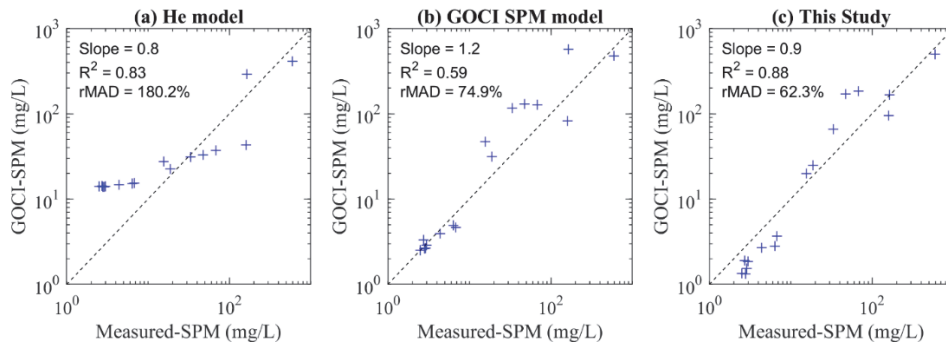


Figure 4.12. Validation of retrieved SPM for GOCI match-ups by the GOCI standard model, the He model, and our proposed model.

## 4.4 Discussion

### 4.4.1 Uncertainties in the atmospheric correction algorithm

Due to no GOCI data available for the cruise in February of 2011, we were not able to validate the performance of the revised AC in extremely turbid waters. As shown in Figure 4.3,  $R_{rs}(660)$  derived from the GOCI standard AC matches very well with *in-situ* measurements with an rMAD of 14.2%. Note that the maximum  $C_{SPM}$  of the 16 match-ups is 591 mg/L, which indicates that GOCI standard AC

could derive acceptable  $R_{rs}(660)$  for turbid waters, which is also reported in Doxaran et al. (2014). However, it is reasonable to assume that the empirical-in-nature relationships between  $R_{rs}(\text{NIR})$  and  $R_{rs}(660)$  for extremely turbid waters, of which majority data was collected in February of 2011, could probably result in uncertainties when applied to GOCI images acquired in 2013.

As shown in Figure 4.11, both our model and the GOCI SPM model show a comparable range of derived  $C_{\text{SPM}}$ , while  $C_{\text{SPM}}$  from the He model seems to be underestimated. The recalibrated He model has shown comparable performance when applied to the Val dataset for extremely turbid waters (see Figure 4.10 and Table 4-2), which means the underestimation in the SPM map is associated with  $R_{rs}(745)$  from the revised AC, if we assume GOCI derived  $R_{rs}(490)$  is of less uncertainties in the extremely turbid waters.

Although the AC may be imperfect, the mapping SPM products provide still useful insights into the applicability of the three models. The He model utilizes only two wavebands at 745 and 490nm, and uncertainties in  $R_{rs}(745)$  could heavily affect the retrieval accuracy of  $C_{\text{SPM}}$  using Eq. (4.7). In contrast, the GOCI SPM model involves four wavebands, three of which in the visible range, and the  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model is based on retrieved  $b_{bp}$ , which is derived by taking account of all the spectral information of  $R_{rs}$ . In the latter two cases, the uncertainty of  $R_{rs}(\text{NIR})$  would be partly compensated for if  $R_{rs}$  at visible bands is of good accuracy.

#### 4.4.2 The proposed $S_{\text{index}}\text{-}C_{\text{SPM}}$ model

##### 4.4.2.1 The sigmoid index of $b_{bp}$

The sigmoid function (Eq. (4.16)) is not simply the mathematical approach to avoid the discontinuities issue that could result from the two respective relationships between  $b_{bp}$  and  $C_{\text{SPM}}$  for clear and turbid waters, but also has physical meaning that considers the saturation of  $b_{bp}$  under extremely turbid conditions. The power-law relationship expressed in Eq. (4.14) suggests a logarithmic increase of  $b_{bp}$  with increasing turbidity, which is to the contrary of observations in natural waters. Remote sensing reflectance has been observed to saturate in the visible bands when  $C_{\text{SPM}}$  reach a certain level, and the saturation will shift to longer wavelengths with continuing increasing of  $C_{\text{SPM}}$  (Knaeps et al., 2012; Knaeps et al., 2015; Luo et al., 2018; Shen et al., 2010b). The key physics behind the  $R_{rs}$  saturation at the given waveband is that the wave-leaving signal, mainly due to the particulate backscattering, is less sensitive to the increasing turbidity. Also, for commercial devices that measure the

backscattering coefficient, saturation of the  $b_{bp}$  is also observed when the  $C_{SPM}$  extend a certain limit. Though current  $b_{bp}$  device, *e.g.*, WET Labs ECO-BB and HOBI Labs HydroScat, is not able to obtain reasonable reading in extremely waters due to complex multi-scattering events (Doxaran et al., 2016), it is, however, reasonable to assume that  $b_{bp}$  device won't be able to distinguish  $C_{SPM}$  between 4000 mg/L and 5000 mg/L.

Mathematically, the sigmoid function forces the  $b_{bp}$  index saturated, or with limited increment, for extremely turbid waters. For example, in order to increase derived  $C_{SPM}$  from 4000 mg/L to 5000 mg/L, the power-law function expressed in Eq. (4.14) requires derived- $b_{bp}$  increasing from 9.7 to 11  $m^{-1}$ , while increment of  $b_{bp}$  from 7.8 to 8.1  $m^{-1}$  will satisfy the  $C_{SPM}$  changes from the sigmoid function of Eq. (4.16). The determination should be as close as the “true” maximum  $b_{bp}$  for better model performance. However, it is hard to pre-determine the maximum  $b_{bp}$  for optically-complex environments like the Yangtze estuary, especially for remote sensing applications. The selected  $\max(b_{bp})$  of 10  $m^{-1}$  in this study for model calibration is arbitrary, which attempts to cover the maximum potential  $C_{SPM}$  that could exist in the Yangtze estuary, as field measured  $C_{SPM}$  could beyond 5000 mg/L (He et al., 2013).

#### 4.4.2.2 Uncertainties of the $S_{index}$ - $C_{SPM}$ model

The 2SeaColor model was previously validated in various types of waters with fairly accurate estimations of IOPs (Arabi et al., 2016; Salama and Verhoef, 2015; Yu et al., 2016). The consistently good performance of the model could enhance our confidence in adopting it to retrieve  $b_{bp}$  in extremely turbid waters. More importantly,  $b_{bp}$  in this study serves as an intermediate variable of  $R_{rs}$  to  $C_{SPM}$ , which means the uncertainties in the retrieved  $b_{bp}$  could be compensated for by regionally well calibrated  $b_{bp}$ - $C_{SPM}$  relationships.

The empirical constants of  $S_{index}$ - $C_{SPM}$  model expressed in Eq. (4.16) is fundamentally determined by the two respective relationships between  $b_{bp}$  and  $C_{SPM}$  in the Cal dataset for clear and turbid waters. Therefore, the calibrated  $S_{index}$ - $C_{SPM}$  model could be subjected to the uncertainties raised from the two respective  $b_{bp}$ - $C_{SPM}$  relationships when applied to different regions or the same region of different seasons.

For relatively clear waters, the slope of the linear regression between  $C_{SPM}$  and  $b_{bp}$ , which can be interpreted as the projected concentration per unit area (hereafter denoted as  $\eta$ ) (Mikkelsen, 2002), and its variation could be significant

in global coastal waters. The  $\eta$  is 59.83 g.m<sup>-2</sup> in the Yangtze estuary (Eq. (4.13)), while the  $\eta$  values are 112.71 g.m<sup>-2</sup> ( $R^2 = 0.62$ ,  $N = 56$ ) and 143.13 g.m<sup>-2</sup> ( $R^2 = 0.66$ ,  $N = 73$ ) for coastal waters in the Baltic Sea and the North Sea, respectively (figures not shown here). Calculations for the latter two  $\eta$  values are based on the COASTALOO dataset (Babin et al., 2003), with  $b_{bp}(555)$  retrieved from  $R_{rs}$  using the 2SeaColor model. The regional and temporal variability of the  $b_{bp}$ - $C_{SPM}$  relationship can be explained by the differences in particles type, composition (refractive index) and concentration, which determine the  $b_{bp}$  (Neukermans et al., 2012; Woźniak et al., 2010). Besides, the presence of large amounts of phytoplankton in water would have a notable impact on the linear relationship between  $b_{bp}$  and  $C_{SPM}$ , even though phytoplankton itself has an insignificant contribution to  $b_b$  (Babin et al., 2003; Boss et al., 2004). Therefore, the calibrated linear relationship between  $b_{bp}$  and  $C_{SPM}$  could be expected with larger errors for application in other estuarine and coastal waters.

For extremely turbid waters, the impact of particle type and composition on the  $b_{bp}$ - $C_{SPM}$  relationship, though less documented, was indirectly demonstrated in this study. The power law relationship between  $C_{SPM}$  and  $b_{bp}$  (*i.e.*, Eq. (4.14)), with an exponent  $>1$ , represents the decline of backscattering efficiency by suspended particles with increasing turbidity. This is mainly because fine fluvial sediments are expected to flocculate when they meet saline water and dispersal in the form of flocs in the Yangtze estuary (Guo and He, 2011; Shi, 2010). Due to fine particles backscatter more pronouncedly than coarse ones, the flocs in the extremely turbid waters in the Yangtze estuary, with the mean diameter up to an order of magnitude larger than the primary particles (Guo and He, 2011), would significantly reduce the backscattering efficiency. Flocculated particles are believed to be the dominant form of suspended sediments in most world's rivers and estuaries, and could account for more than 80% of the total suspended sediments in the estuary (Droppo and Ongley, 1994). Therefore, the proposed power-law relationship between  $b_{bp}$  and  $C_{SPM}$  could potentially be also applicable to other major turbid estuaries.

To evaluate the performance of the calibrated  $b_{bp}$ - $C_{SPM}$  relationships in other optically-complex estuaries, we applied the calibrated  $b_{bp}$ - $C_{SPM}$  models (includes Eq. (4.13) and Eq. (4.14)) and the  $S_{index}$ - $C_{SPM}$  model (Eq. (4.16)) to the independent SeaSWIR dataset without tuning the empirical constants.  $b_{bp}$  of the SeaSWIR dataset was first retrieved from convoluted  $R_{rs}$  at 8 GOCI bands using the 2SeaColor model. Figure 4.13 shows the comparison between measured  $C_{SPM}$  and retrieved  $C_{SPM}$  from the  $S_{index}$ - $C_{SPM}$  model (blue cross) and the  $b_{bp}$ - $C_{SPM}$  model



(black circle). Note that retrievals of the  $b_{bp}$ - $C_{SPM}$  model in Figure 4.13 employ Eq. (4.13) for samples with  $C_{SPM} < 100$  mg/L and Eq. (4.14) for samples with  $C_{SPM} > 100$  mg/L, respectively. As shown in Figure 4.13, the power-law  $b_{bp}$ - $C_{SPM}$  relationship for turbid waters seems to be robust for the SeaSWIR dataset with data points well-distributed along the 1:1 line, especially for samples with  $C_{SPM} > 200$  mg/L. However, the linear  $b_{bp}$ - $C_{SPM}$  relationship for relatively clear waters shows substantial errors, which echoes our previous discussion on the applicability of the  $b_{bp}$ - $C_{SPM}$  relationships in different types of water. The performance of the  $S_{index}$ - $C_{SPM}$  model for the whole SPM range shows a great improvement of  $C_{SPM}$  retrieval accuracy with rMAD of 30.8% compared to 45.2% of the  $b_{bp}$ - $C_{SPM}$  model. More encouragingly, the  $S_{index}$ - $C_{SPM}$  model presents strong resistance to uncertainties even for the relatively clear waters. The rMAD of the  $S_{index}$ - $C_{SPM}$  model for  $C_{SPM} < 100$  mg/L is 20.8% and is 42.8% for the linear  $b_{bp}$ - $C_{SPM}$  model.

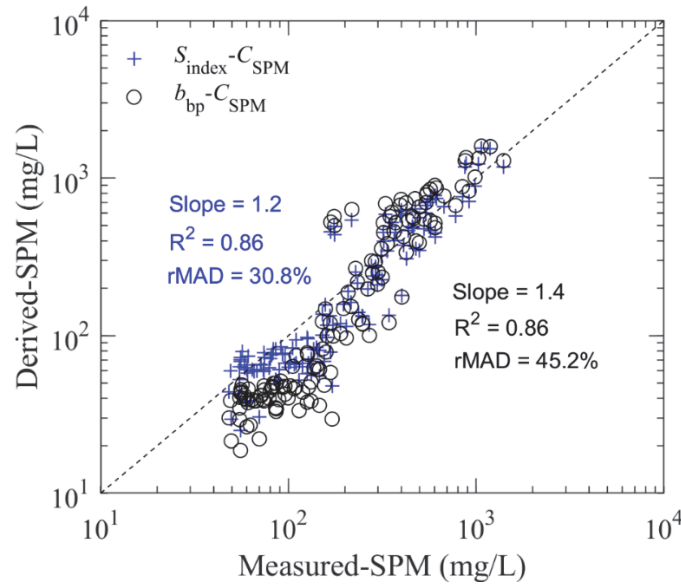


Figure 4.13. Retrieval of  $C_{SPM}$  for the SeaSWIR dataset using the  $b_{bp}$ - $C_{SPM}$  (Eq. (4.13) and Eq. (4.14) and the  $S_{index}$ - $C_{SPM}$  (Eq. (4.16)) relationships calibrated in this study.

#### 4.4.2.3 Discontinuity issue of the retrieved $C_{SPM}$

For the  $C_{SPM}$  retrieval models that employ a switching scheme to pre-classify water types, the predicted  $C_{SPM}$  are always involved with discontinuity issue (Shen et al., 2010b; Shen et al., 2013). For example, as shown in Figure 4.8, for a given  $b_{bp}$  of  $1.5 \text{ m}^{-1}$ , which roughly locates at the switching threshold of the two  $b_{bp}$ - $C_{SPM}$  relationships, estimated  $C_{SPM}$  could increase from 90 mg/L using the

linear model to 169 mg/L when the power-law function applies. For the GOCI SPM model, a threshold of  $R_{rs}(660) = 0.04 \text{ sr}^{-1}$  was used to initiate the switch between Eq. (4.8) and Eq.(4.9). As a result, produced  $C_{\text{SPM}}$  map by the GOCI SPM model present obvious discontinuities at the front of the turbid plume in the Subei Bank and the inshore area of the Yangtze estuary (Figure 4.11c). Contrarily, the proposed  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model, which uses a single equation to retrieve the wide-range  $C_{\text{SPM}}$ , avoid the discontinuity problem. As shown in Figure 4.11a, the spatial distribution of  $C_{\text{SPM}}$  is more realistic at the front of the turbid plume. Therefore, the  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model could be more suitable for  $C_{\text{SPM}}$  retrieval in the optically-complex environment where  $C_{\text{SPM}}$  varies in 3 to 4 orders of magnitude.

#### **4.4.3 Diurnal dynamics of SPM concentrations**

The GOCI image, available each hour continuously from 00:16 UTC to 07:16 UTC per day, greatly facilitates the investigation on diurnal dynamics of  $C_{\text{SPM}}$  at large spatial scale. The tidal type in the Yangtze estuary is semidiurnal, which means the GOCI data could be adequate to reveal the temporal dynamics of  $C_{\text{SPM}}$  during, at least, a half tidal cycle. To present the diurnal dynamics of SPM, the proposed  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model was applied to 8 GOCI images acquired on March 8<sup>th</sup>, 2013, as shown in Figure 4.14.

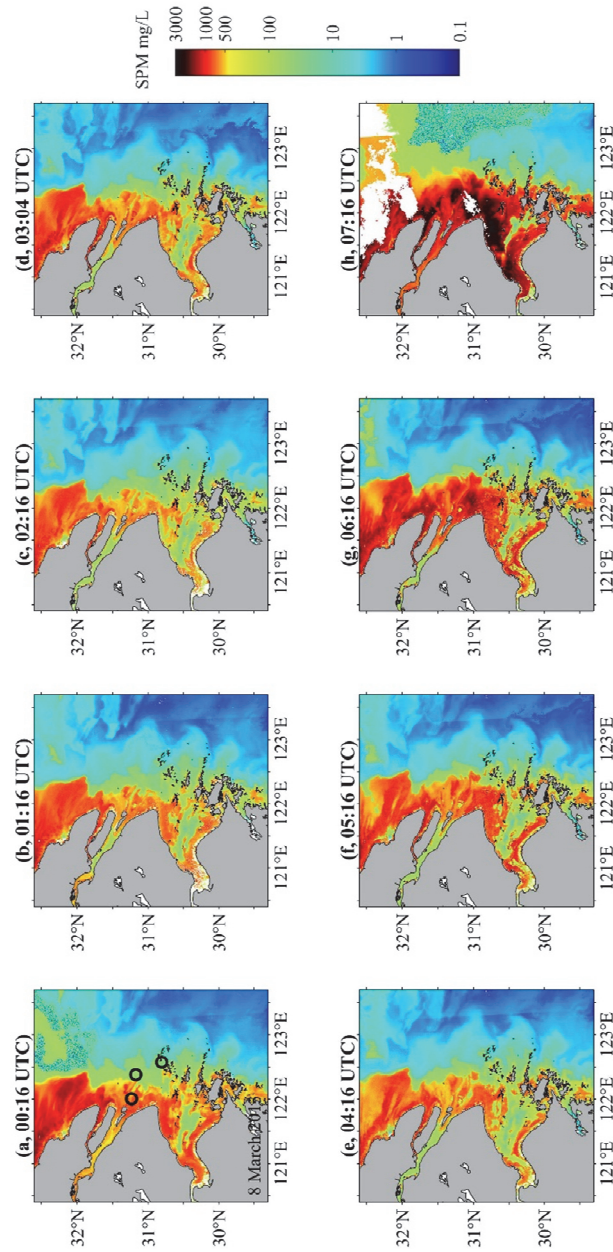


Figure 4.14. Diurnal variations of SPM concentrations in the Yangtze estuary on March 8th, 2013. The black circles in (a) represent the locations of the hydrologic stations.

It is clearly demonstrated that diurnal variation of  $C_{\text{SPM}}$  is noteworthy in the Yangtze estuary, where SPM concentrations could vary by 1 to 2 orders of magnitude for some areas. To better understand the dynamics of SPM during this

time period and the influence of the tide, tidal level data at three hydrologic stations (the black cycles in Figure 4.14a, also referred to St. 1, St. 2 and St.3 in Figure 4.1a) were compared with retrieved  $C_{SPM}$  from GOCI images for matched location and time, as presented in Figure 4.15.

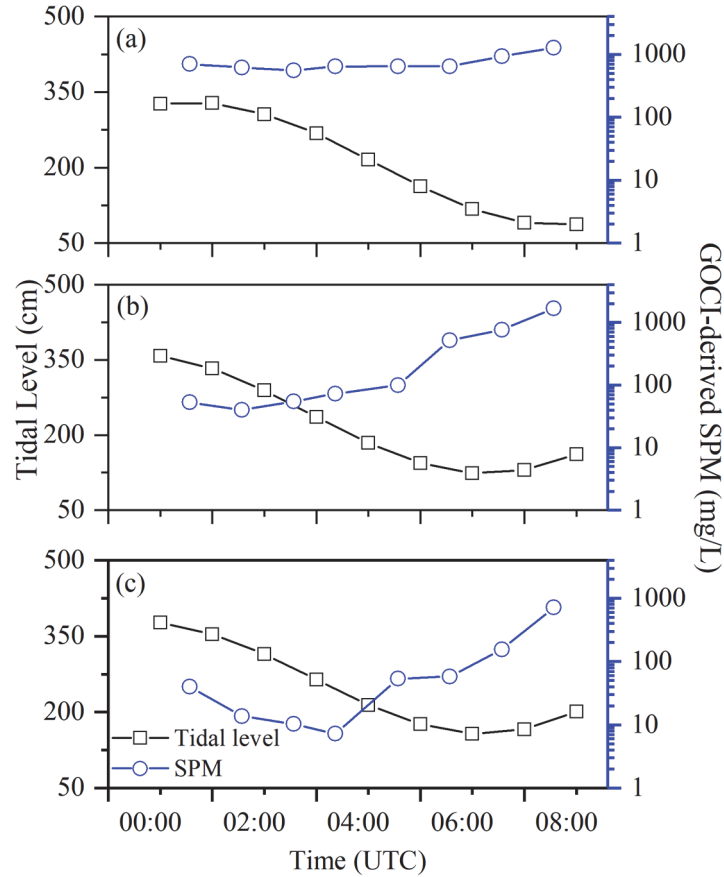


Figure 4.15. Diurnal dynamics of CSPM (the right ordinate axis) at three hydrometric stations corresponding to the variations of the tidal level (the left ordinate axis). Figs. a, b, and c are corresponding to the stations labeled St. 1, St.2, and St. 3 in Figure 4.1a, respectively.

As shown in Figure 4.15, SPM concentrations show a strong dependency on the tide, with increasing  $C_{SPM}$  during the ebb tide. This can be explained by the transportation of suspended sediments during the ebb tide, where ebb-dominated tidal flow transits from landward high- $C_{SPM}$  regions to the offshore carrying large amounts of suspended sediments (Liu et al., 2010). Additionally, tide-induced mixing and resuspension will be enhanced during the ebb tide, with more sediments from the bottom layer re-suspended to surface waters (Chen et al.,

2006). Sediment resuspension, as a key factor in dominating the SPM spatial-temporal variation in coastal waters, has also been observed in other turbid estuaries (Doxaran et al., 2009).

Figure 4.15 also shows that  $C_{SPM}$  increases more rapidly after the tidal current speed reaches its maximum, which is roughly in the middle phase of the ebb tide. This is mainly because the surface  $C_{SPM}$  is much slower in responding to sediment exchange, especially horizontally, than other layers in water column because of the larger depth (Yang et al., 2004). This time lag could be attributed to the fluid shear stress that determines the suspension and settling velocities of flocs in the estuary (Guo and He, 2011; Shi, 2010), which results in the phase difference between tidal current and surface  $C_{SPM}$ . The same reason explains why the maximum  $C_{SPM}$  is observed approximately 1-2 hours later after ebb slack water (the status when ebb tide reaches the lowest level), as observed in all three stations in Fig. 10. This phenomenon has also been reported from field observations in the Yangtze estuary by (Yang et al., 2004; Zhang et al., 2015), from whose study the time lag between the slack water and maximum  $C_{SPM}$  could vary from 1 to 3 hours from location to location. Vice versa, there could also be time lags for the surface  $C_{SPM}$  decreases to the minimum after the flood slack water. For instance, as shown in Figure 4.15, it takes 2 to 3 hours for seawater to dilute and compensate the tidal-induced SPM mixing and resuspension before  $C_{SPM}$  reaches the minimum.

It is also expected that winds, river discharges, and ocean currents could also have significant effects on the spatial-temporal distribution of SPM in the Yangtze estuary (Chen et al., 2003). However, their impact on a diurnal scale could rather be limited, and tide could be the most dominant factor in determining the diurnal dynamics of  $C_{SPM}$  (He et al., 2013). The results presented here indicate that reasonable and realistic diurnal dynamics of  $C_{SPM}$  in the Yangtze estuary could be reconstituted from GOCI imagery by our proposed model. Further efforts could be a focus on using GOCI imagery to better understand and monitor SPM dynamics in the Yangtze estuary at large temporal-spatial scales.

## 4.5 Conclusions

In this study, we proposed a novel model to estimate wide-range SPM concentrations in the Yangtze estuary from remotely retrieved  $b_{bp}(555)$  using the 2SeaColor model. This model, using a sigmoid function to take account of the different relationships between  $b_{bp}(555)$  and  $C_{SPM}$  for relatively clear and

extremely turbid waters, was calibrated for data collected in the Yangtze estuary with  $C_{\text{SPM}}$  ranging from 0.4 to 2068.8 mg/L. Validation results with both the Yangtze dataset and an independent SeaSWIR dataset show the robustness of the proposed  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model in retrieving  $C_{\text{SPM}}$  for optically-complex coastal waters with a wide range of  $C_{\text{SPM}}$ .

The  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model was later applied to GOCI images and reproduced reasonable and repeatable  $C_{\text{SPM}}$  maps for both clear and turbid waters, where switching-based model shows conspicuous discontinuities at the front of the turbid plume. The model was further implemented to a time-series of GOCI imagery and produced realistic diurnal dynamics of  $C_{\text{SPM}}$ , which can be well explained by the hydrologic process. The encouraging performance of the proposed  $S_{\text{index}}\text{-}C_{\text{SPM}}$  model shows the potential of monitoring SPM in the optically complex water in the Yangtze estuary with high accuracy.

## **Chapter 5      Light absorption properties of CDOM in the Yangtze estuarine and coastal waters: an alternative approach for DOC estimation\***

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\* This Chapter is based on:

Yu, X., Shen, F., & Liu, Y. (2016b). Light absorption properties of CDOM in the Changjiang (Yangtze) estuarine and coastal waters: An alternative approach for DOC estimation. *Estuarine, Coastal and Shelf Science*, 181, 302-311

## **5.1 Introduction**

Colored dissolved organic matter (CDOM) is an important component of natural waters and plays various roles in the functioning of ecological processes and biogeochemical cycles in coastal ecosystems. Light absorption property of CDOM mainly depends on the origins of CDOM, as terrestrial source contains more humic acid and large size of dissolved organic matter (DOM) molecules while aquatic source contains more fulvic acid and small size of DOM molecules (Helms et al., 2008). Riverine discharge is always considered as the main source of CDOM in most coastal waters, resulting in a robust correlation between salinity and CDOM absorption coefficient (Fichot and Benner, 2011; Granskog, 2012; Xie et al., 2012). The aquatic CDOM is mainly produced locally by phytoplankton degradation and bacterial decomposition. To get better insights of CDOM sources, spectral slope ( $S$ , in  $\text{nm}^{-1}$ ) of CDOM absorption was recommended as it is related to the ratio of fulvic to humic acids and the molecular weight (MW) of fulvic acids (Carder et al., 1989; Keith et al., 2002). Moreover,  $S$  calculated from a narrow waveband in the ultraviolet spectral range (*e.g.*, 275-295 nm) is more sensitive to CDOM sources than the  $S$  determined from a broader waveband (Asmala et al., 2012; Fichot and Benner, 2011, 2012). This is mainly because the  $S_{275-295}$  is negatively related to DOM molecular weight (MW) (Helms et al., 2008).

For the past few decades, the global carbon cycle has been one of the hottest researches in the field of global climate change due to  $\text{CO}_2$ , the major greenhouse gas (Canadell and Dickinson, 2006). Dissolved organic carbon (DOC), on the other hand, represents 97% of the organic carbon in the ocean (Hansell and Carlson, 1998), is an essential part of the global carbon cycle. Phytoplankton converts inorganic carbon to organic carbon through photosynthesis (Longhurst and Harrison, 1989), which represents a substantial source of DOC in the oceans. Other sources of DOC in the oceans mostly come from riverine discharge and zooplankton activity (Kuliński and Pempkowiak, 2008). However, traditional DOC measurements are always limited due to discrete sampling and measurement. Therefore, a prompt and continuous estimation from earth observation data on large spatial scales would be strongly recommended (Liu et al., 2013; Mannino et al., 2008).

CDOM is the optically active component of DOM and therefore can be potentially linked to DOC concentration. Although CDOM only represents a portion of the entire DOC pool, optical properties of CDOM have been found to



be empirically related to DOC concentration in some estuarine and coastal regions (Asmala et al., 2012; Ferrari, 2000; Fichot and Benner, 2011, 2012; Matsuoka et al., 2012a; Rochelle-Newall et al., 2014; Spencer et al., 2007). There are mainly two methods to estimate DOC concentration from CDOM optical properties. The first approach is linking DOC concentration to CDOM absorption coefficient at a reference wavelength (*e.g.*, 250nm or 350nm) by single linear regression (Baker and Spencer 2004; Matsuoka et al. 2012a; Spencer et al., 2009a). However, linear models rely heavily on the conservative mixing of DOC and CDOM, and sometimes generate large uncertainties and are not practical in some of the coastal waters, such as the Mississippi River estuary (Chen and Gardner, 2004), the Pearl River estuary (Chen et al., 2004) and Finnish estuaries (Asmala et al., 2012). The other approach is using the non-linear model developed by (Fichot and Benner, 2012) to estimate DOC from spectral slope  $S_{275-295}$  and DOC-normalized CDOM absorption coefficient. The non-linear model is supposed to be more suitable for coastal waters, especially for river-influenced coastal waters with the lignin as the important chromophore of CDOM pool and the main terrigenous component of DOC.

The hydrodynamic environment is unique in the Yangtze estuarine and coastal waters due to joint influences by runoffs, tides and coastal circulations, which makes the mixing behavior between DOC and CDOM more complicated and variable in seasonal and regional scales. For example, the conservative mixing behavior was observed in the Yangtze estuary, and DOC concentration was linearly estimated from CDOM absorption coefficient (Liu et al., 2013; Liu et al., 2014; Zhang et al., 2013). However, the linear relationship was not always practical in this region (Cauwet and Mackenzie, 1993; Liu et al., 2013). Hence, a better understanding of the variation of DOC and CDOM relationship and more detailed information on the spatial-temporal variation of CDOM optical properties in the Yangtze estuary are required.

In this study, we first investigate the seasonal and vertical distributions of CDOM absorption coefficient that were less documented in the Yangtze estuarine and coastal waters. We further discuss the potential sources of CDOM in this region and develop an empirical model to estimate DOC concentration from CDOM optical properties. The developed relationship was found robust in the Yangtze estuarine and coastal waters, providing an alternative approach to estimate DOC concentration in large spatial scales from remotely retrieved CDOM optical properties.

## **5.2 Materials and methods**

### **5.2.1 Sampling and field measurement**

Water samples were collected during five cruises in the Yangtze estuarine and coastal waters from 2011 to 2013, and the sampling stations are presented in Figure 5.1. A total number of 114 CDOM samples were collected in summer from the surface layer in July 2011 (32 samples), July 2012 (29 samples) and August 2013 (53 samples). Another 61 CDOM samples were collected in winter in February 2012 (37 samples) and March 2012 (24 samples). Salinity (in PSU, practical salinity units) was recorded by the CTD device (SeaBird Electronics INC) synchronously to CDOM sampling, with salinities at four of the sampling stations were missing record. DOC samples were collected simultaneously with surface CDOM samples in July of 2011 (31 samples) and August 2013 (53 samples). During the cruise in July 2011, we also collected CDOM samples at the water depth of 5 m (31 samples) and 10 m (29 samples) below the surface layer to investigate the vertical variation of CDOM optical properties.

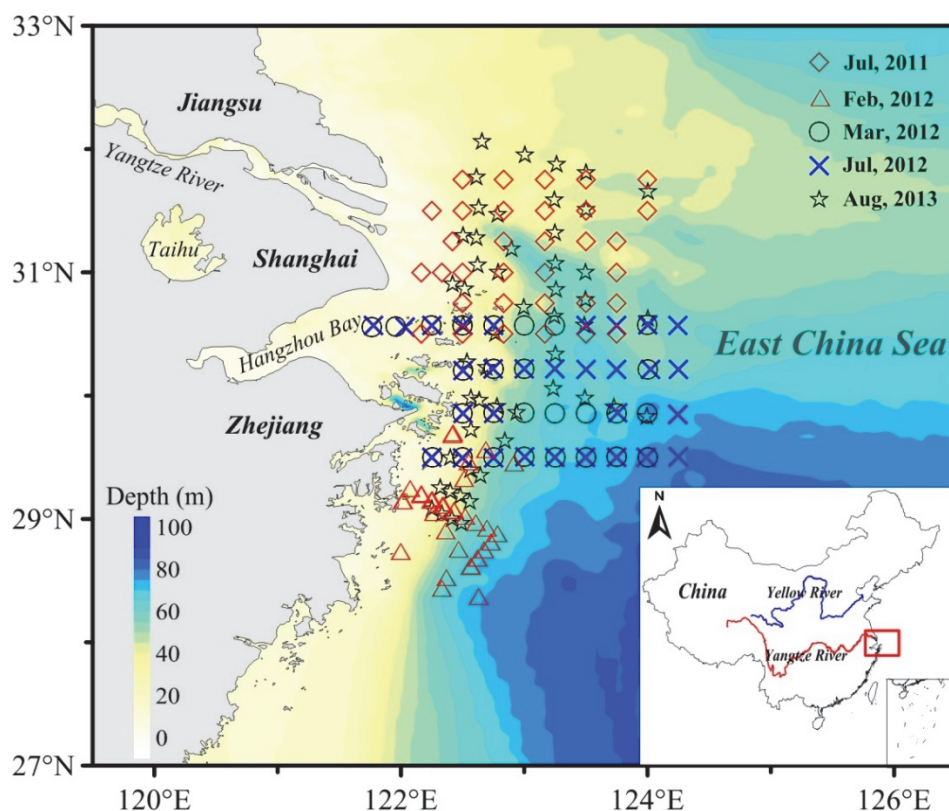


Figure 5.1. Location of sampling stations in the Yangtze estuarine and coastal waters. Samples were collected from five cruises in summer (July 2011, July 2012 and August 2013) and winter (February and March of 2012).

CDOM samples were collected according to the NASA protocol of marine optical measurements (Mueller and Fargion, 2002). Water samples were gravity-filtered on shipboard using 0.22 $\mu$ m polycarbonate membrane (Millipore,  $\Phi$ 47mm) under low vacuum. The membranes were soaked in 10% HCl for 15 mins and then rinsed by Milli-Q water three times before the filtration. The filtered CDOM samples were collected in borosilicate glass vials and then stored in the -40°C refrigerator during the cruises. All vials were pre-soaked in 10% HCl for 24 h, rinsed by Milli-Q water three times, and pre-combusted in an oven at 450°C for 5 h.

DOC samples were filtered to ampoule bottles by a disposable pinhole filter with an aperture of 0.45 $\mu$ m immediately after the water samples were collected. Ampoule bottles were pre-combusted in an oven at 500°C for 4 h. After the filtration, the ampoule bottles were sealed by fusing the bottleneck and then stored in the - 40°C refrigerator.

### 5.2.2 Measurements in lab

To minimize the freezing influence on the CDOM absorption spectra, CDOM samples were unfrozen and warmed in fully dark condition to room temperature immediately prior to measurement. The absorbance spectra were recorded by the PerkinElmer Lambda 1050 UV/VIS spectrophotometer. CDOM absorbance was recorded from 250 to 900 nm with 2 nm interval fitted with a 10cm quartz cuvette and referenced to the Milli-Q water. The absorbance of CDOM ( $D(\lambda)$ ) was converted to absorption coefficient ( $a'_g(\lambda)$ , in  $\text{m}^{-1}$ ) according to,

$$a'_g(\lambda) = 2.303 \times D(\lambda)/l \quad (5.1)$$

where  $l$  is the cuvette path length (*i.e.*, 0.1 m) (Pegau et al. 1997). In order to remove the scattering contribution from the measurements, the CDOM absorption coefficient was normalized to zero at 700 nm (Bricaud et al., 1981),

$$a_g(\lambda) = a'_g(\lambda) - a'_g(\lambda_0) \times \lambda/\lambda_0 \quad (5.2)$$

where  $a_g(\lambda)$  is the CDOM absorption coefficient after the scattering correction and used for data analysis in this study. Figure 5.2 shows the absorption spectra of all CDOM samples.

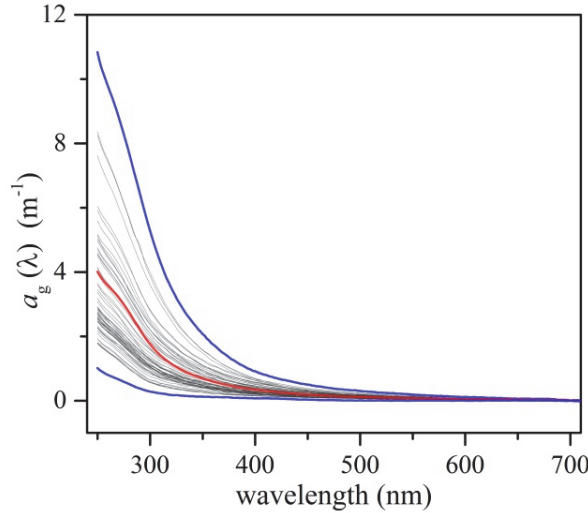


Figure 5.2. CDOM absorption spectra of all measurements. The two blue spectra represent samples with maximum and minimum absorption coefficients, while the red spectrum is the mean absorption coefficient of all samples.

The spectral slope ( $S$ ,  $\text{nm}^{-1}$ ) is derived from CDOM absorption spectra by fitting the absorption data to the exponential equation (*i.e.*, Eq. (5.3)) using the optimization method (Bricaud et al., 1981; Green and Blough, 1994),

$$a_g(\lambda) = a_g(\lambda_0)e^{-S(\lambda-\lambda_0)} \quad (5.3)$$

where  $\lambda$  is wavelength (nm), and  $\lambda_0$  is the reference wavelength (nm). As shown in Eq. (5.3),  $S$  is depended on the wavelength interval over which it is calculated. We chose 440nm as the reference wavelength to calculate the spectral slope for the broad range from 250 and 700nm, which is hereafter denoted as  $S_{440}$ .  $S_{275-295}$  and  $S_{350-400}$  were also calculated for the narrow wavelength interval from 275 to 295 nm and 350 to 400 nm, respectively.  $S_R$  was then defined as the ratio of  $S_{275-295}$  and  $S_{350-400}$ .

DOC concentration (in  $\mu\text{mol/L}$ ) was determined in duplicate with a Shimadzu TOC-VCPH Total Organic Carbon analyzer using potassium hydrogen phthalate as the standard for the DOC calibration curve. The typical precision of this instrument is less than 2% in term of the coefficient of variation. DOC samples were acidified by adding 100  $\mu\text{L}$  2  $\mu\text{mol/L}$  phosphoric acid before the measurement to remove the inorganic carbon. In addition, the Florida Strait seawater reference, provided by Dr. Hansell's laboratory at the University of Miami, was determined every day to ensure the data quality of our measurements. DOC concentration was calculated as the average value of the duplicate measurements and the DOC-normalized CDOM absorption coefficient ( $a_g^*(\lambda_0)$ , in  $\text{L } \mu\text{mol}^{-1} \text{ m}^{-1}$ ) was defined as the ratio of  $a_g(\lambda_0)$  and DOC concentration.

### 5.2.3 Data analysis

To quantitatively evaluate the linear relationship between salinity and CDOM absorption coefficient and spectral slopes, we calculate the Pearson's coefficient  $r$  and probability  $p$  of the linear regression equation. For the goodness-of-fit between derived and known DOC concentration, statistical parameters such as the determination coefficient  $R^2$ , the mean of absolute relative difference (rMAD) and the root mean square error (RMSE) are determined. The rMAD and the RMSE are determined by Eq. (3.5) and Eq. (3.6).

## 5.3 Results

### 5.3.1 CDOM light absorption and its spatial-temporal distribution

Based on the CDOM measurements from the five cruises over the Yangtze estuarine and coastal water, the spatial distribution of CDOM absorption coefficient ( $a_g(355)$ ) at the surface layer is presented in Figure 5.3.

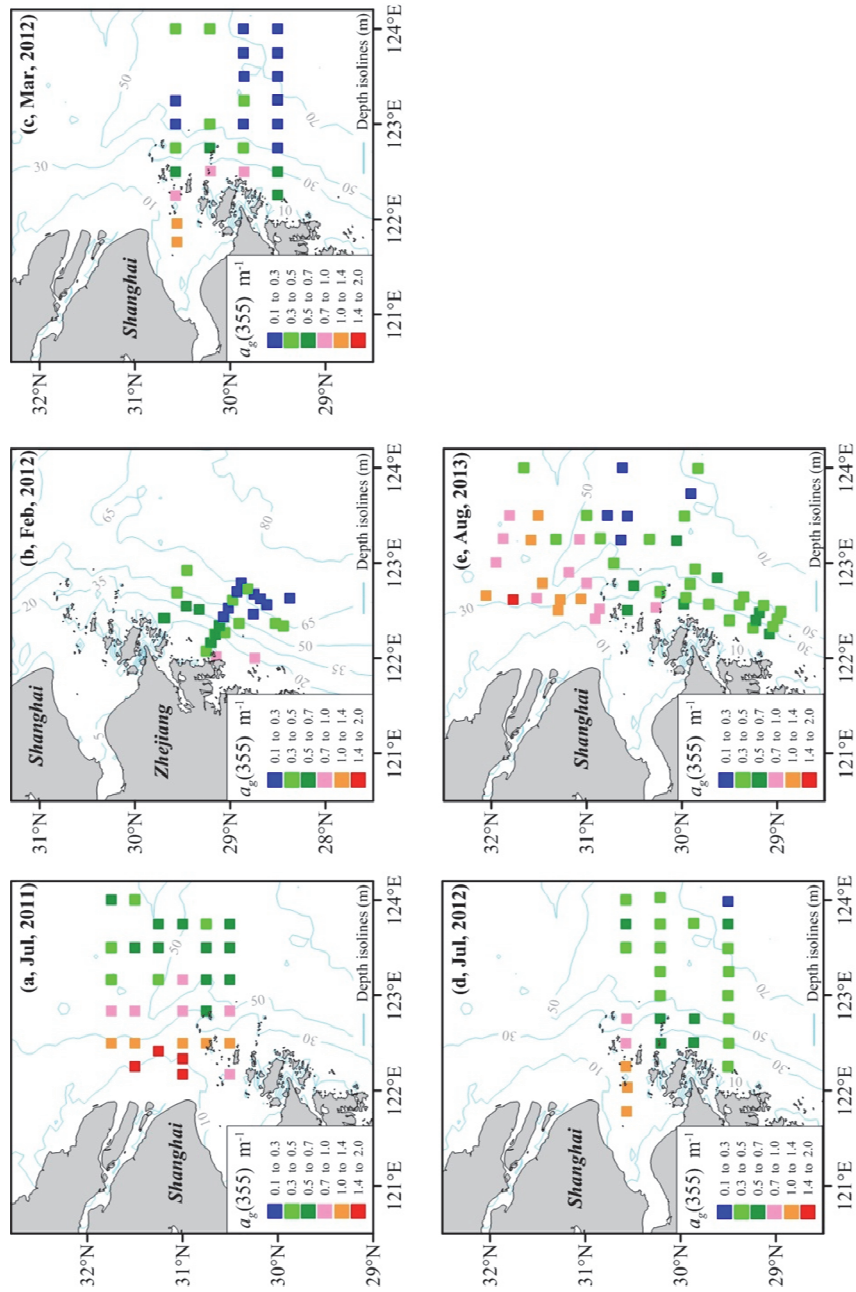


Figure 5.3. Spatial distribution of surface over the Yangtze estuarine and coastal region. Data were collected from five cruises in July of 2011 (a), February (b), March (c), July (d) of 2012 and August of 2013 (e).

Figure 5.3 shows a noticeable spatial-temporal variation of surface  $a_g(355)$  over the Yangtze estuarine and coastal waters. The value of surface  $a_g(355)$  ranged

from 0.42 to 1.91  $\text{m}^{-1}$  in July 2011, from 0.11 to 0.84  $\text{m}^{-1}$  in February 2012, from 0.13 to 1.21  $\text{m}^{-1}$  in March 2012, from 0.26 to 1.36  $\text{m}^{-1}$  in July 2012 and from 0.23 to 1.45  $\text{m}^{-1}$  in August 2013, respectively. The surface  $a_g(355)$  exhibited a gradual decrease from nearshore to offshore, formed a northwest-southeast gradient, and eventually kept steady and low (less than 0.5  $\text{m}^{-1}$ ) in regions far from the river mouth. Most samples in March and July 2012 were collected at same sites (See Figure 5.3c and Figure 5.3d), from which we can observe that CDOM absorption coefficients in summer (*e.g.*, in July) were higher than that in winter (*e.g.*, in March). The annual variation of  $a_g(355)$  over the study area in summer from 2011 to 2013 is also observed according to Figures 5.3a, 5.3d, and 5.3e. The average values of  $a_g(355)$  are 0.89  $\text{m}^{-1}$ , 0.52  $\text{m}^{-1}$  and 0.60  $\text{m}^{-1}$  in July 2011, July 2012 and August 2013, respectively.

As shown in Figure 5.4, the variations of  $a_g(355)$  in July 2011 in a vertical direction among different stations are rather flexible, and no clear pattern can be found. In general, water is more stratified at inshore stations compared to offshore stations. The largest variations of  $a_g(355)$  depth-profile are found at stations adjacent to the Yangtze River mouth (*e.g.*, stations within longitude of 122.5°E in Figures 5.4c, 5.4e and 5.4f), while vertical variations at most offshore stations are rather small, especially for stations with longitude larger than 123°E.

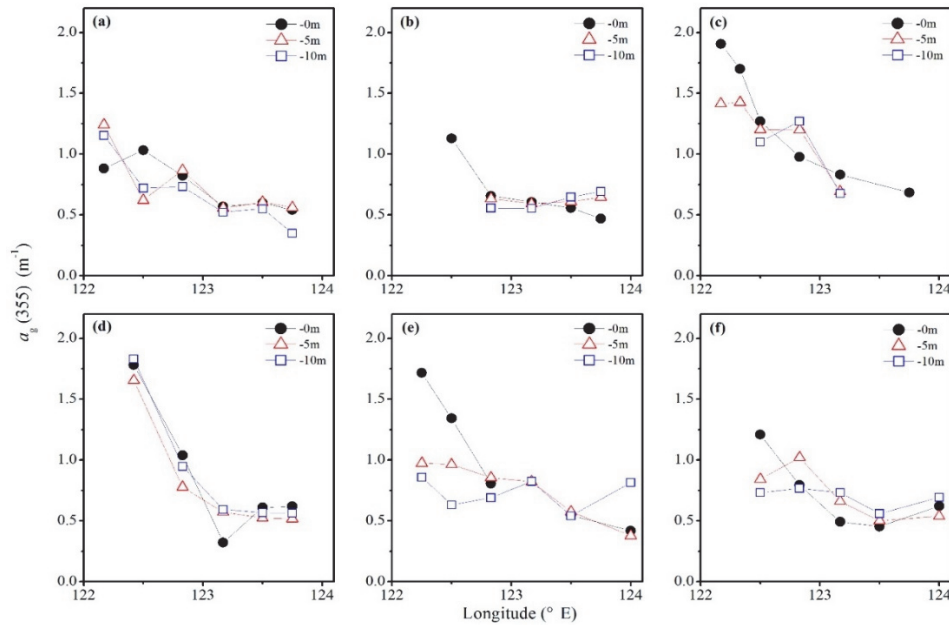


Figure 5.4. Vertical distribution of  $a_g(355)$  along geographic longitude increasing in July

2011. Panels 4a - 4f are corresponding to the six axial transections in Figure 5.3a from the bottom to the top.

### 5.3.2 CDOM absorption coefficient and salinity

Salinity is a perfect indicator of the influence of river runoff in the estuarine regions. The variation of salinity is mainly due to the mixture of fresh water from river and ocean water. Freshwater is almost extinguished from salinity while the surface salinity in Open Ocean is typically between 32 and 37 PSU (Lagerloef et al., 1995). To better explore the relationship between CDOM absorption coefficient and salinity, CDOM measurements are divided into subsets based on two criteria. The first classification is based on sampling seasons, *i.e.*, winter and summer. The other one is according to the salinity, from which three subsets are classified with two thresholds of 20 PSU and 32 PSU. We consider that samples with salinity lower than 20 PSU are riverine-dominant while marine-dominant samples have salinities greater than 32 PSU. Figure 5.5 shows the relationships between salinities and  $a_g(355)$ , as well as their linear fitting parameters.

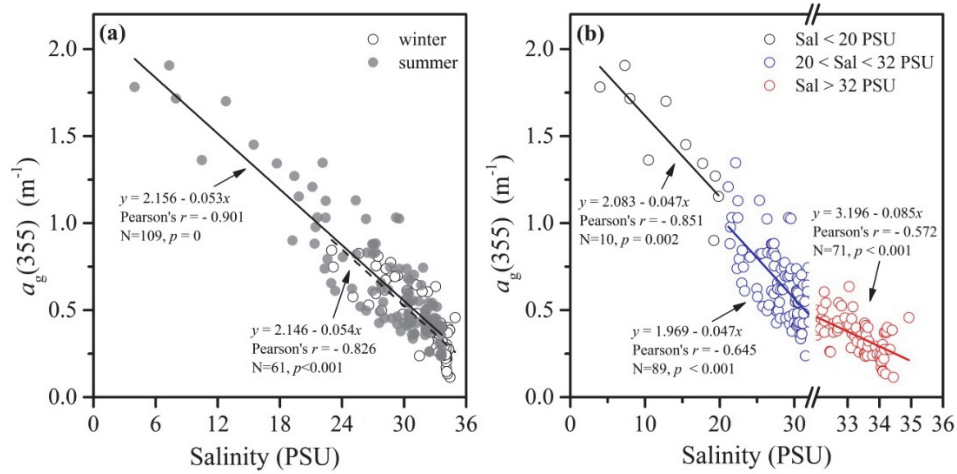


Figure 5.5. Relationships between salinity and  $a_g(355)$ . All CDOM measurements are divided into (a) two subsets for samples collected in winter and summer and (b) three subsets according to the salinity with two thresholds at 20 and 32 PSU.

From Figure 5.5 we can find that  $a_g(355)$  is negatively correlated with salinity in our study area. However, this relationship is more robust in summer (Figure 5.5a, Pearson's coefficient  $r = -0.901$ ,  $N = 109, p = 0$ ) than that in winter (Figure 5.5a, Pearson's coefficient  $r = -0.826$ ,  $N = 61, p < 0.001$ ). It also can be found that the negative relationship is more robust for samples with salinity lower than 20 PSU



(Figure 5.5b, Pearson's coefficient  $r = -0.851$ ,  $N = 10$ ,  $p = 0.002$ ) compared to samples with higher salinities.

### 5.3.3 Spectral slope of CDOM absorption and salinity

Spectral slope of CDOM light absorption also varied significantly in spatial and temporal scales in this study. The range of  $S_{440}$  for all samples was from 0.009 to 0.019  $\text{nm}^{-1}$  with the average value of 0.015  $\text{nm}^{-1}$ . Seasonal variations of  $S_{440}$  are also observed as the average value of  $S_{440}$  in July 2012 was  $0.017 \pm 0.001 \text{ nm}^{-1}$  while the average value in March 2012 was  $0.0145 \pm 0.0017 \text{ nm}^{-1}$ . On the other hand,  $S_{275-295}$  ranges from 0.015 to 0.032  $\text{nm}^{-1}$  for all observations with the average value of 0.021  $\text{nm}^{-1}$ . However, opposite to the seasonal variation of  $S_{440}$ , the average value of  $S_{275-295}$  in March 2012 ( $0.023 \pm 0.004 \text{ nm}^{-1}$ ) was higher than that in July 2012 ( $0.018 \pm 0.002 \text{ nm}^{-1}$ ).

There is no clear relationship between  $S_{440}$  and salinity in this study (figure not shown here), while both  $S_{275-295}$  and  $S_R$  increased along the increasing salinity gradient as shown in Figure 5.6. The patterns between salinity and  $S_{275-295}$  and  $S_R$  are consistent with previous studies (Fichot and Benner, 2012; Helms et al., 2008). Moreover, the linearity is more robust in low salinity samples ( $< 20 \text{ PSU}$ ) for both  $S_{275-295}$  and  $S_R$ .

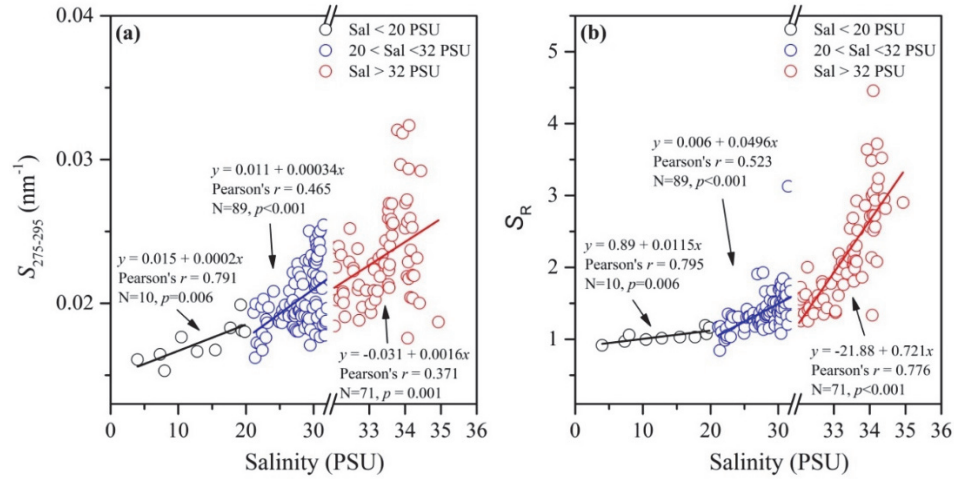


Figure 5.6. Scatter plots between salinity and  $S_{275-295}$  (a) and  $S_R$  (b). All observations are divided into three subsets using thresholds of salinity at 20 PSU and 32 PSU.

### 5.3.4 CDOM properties and DOC

DOC concentrations ranged from 60.54 to 217.06  $\mu\text{mol/L}$  for measurements from July 2011 and August 2013 with an average value of 115.41  $\mu\text{mol/L}$ . DOC

was unrelated to both salinity and the CDOM absorption coefficient in this study, presenting a non-conservative mixing behavior between CDOM and DOC. To estimate DOC from CDOM properties, we adopted the approach of Fichot and Benner (2012) and developed a non-linear relationship between  $S_{275-295}$  and DOC-normalized CDOM absorption coefficient at 355nm (*i.e.*,  $a_g(355)/\text{DOC}$ , denoted as  $a_g^*(355)$ ). All the DOC samples were divided into two optimal calibration and validation sets according to the Cal/Val model (Salama et al., 2012). The Cal and Val sets were selected such that they have the same moments (mean and standard deviation). An exponential function was used to fit  $a_g^*(355)$  and  $S_{275-295}$  for the calibration dataset (Figure 5.7a), resulting in a coefficient of determination ( $R^2$ ) of 0.78. The fitting equation is expressed as,

$$a_g^*(355) = a_g(355)/\text{DOC} = \exp(-2.45 - 130.79S_{275-295}) \quad (5.4)$$

The fitting equation from Fichot and Benner (2012) is also provided in Figure 5.7a for comparison. The developed model was applied to the Val dataset to estimate DOC from CDOM absorption coefficient and  $S_{275-295}$  as shown in Figure 5.7b. The uncertainty of this model is clarified with the  $R^2$  of 0.746, an RMSE of 20.99  $\mu\text{mol/L}$  and an rMAD of 14.46%.

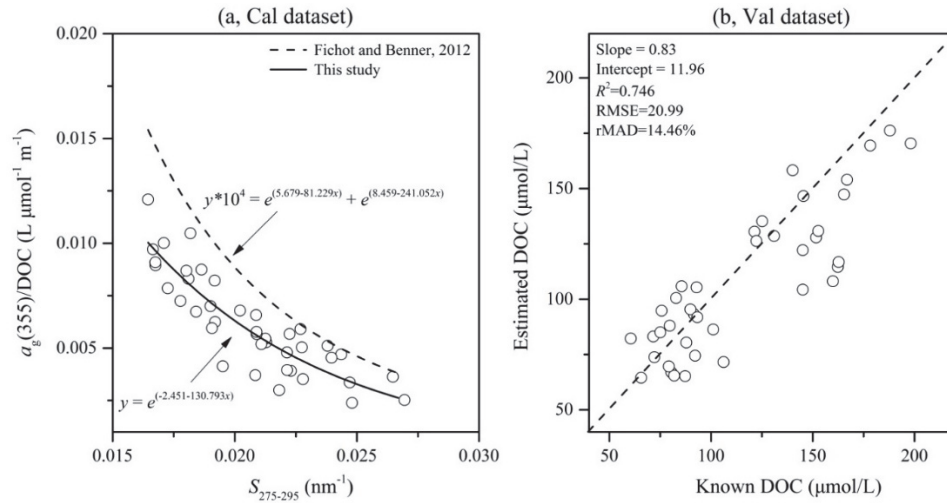


Figure 5.7. Calibration and validation for DOC estimation model using the Cal and Val datasets. (a) A non-linear relationship (black line) is used to fit  $a_g(355)/\text{DOC}$  and  $S_{275-295}$  using the Cal dataset. A reference relationship from Fichot and Benner (2012) is also provided (dash line). (b) Validation of the developed model using the Val dataset.

## 5.4 Discussion

### 5.4.1 CDOM optical properties among world's estuarine and coastal waters

Table 5-1 presents the comparison of several investigated parameters around the world's estuarine and coastal waters, including the ranges of CDOM absorption coefficient, spectral slope  $S_{440}$ , salinity, and DOC concentration. In general, the level of  $a_g(355)$  in this study was on par with results from Pearl River estuary and Mackenzie River estuary but was much lower than that in some other major estuaries as shown in Table 5-1.

The abundance of DOM in the estuarine and coastal water is always influenced by vegetation, climate, the type of rock and soil in the river basins. As shown by the Terra MODIS NDVI images, the coverage of vegetation in the Yangtze River basin is lower compared to other drainage basins of major world rivers, which is probably due to rapid urbanization and massive deforestation in the past decades. Therefore, the low degree of vegetative cover in the Yangtze River basin could be an important factor in determining the relatively low content of soil organic matter in the drainage basins, resulting in the relatively low concentration of CDOM in the Yangtze estuary. The low degree of vegetative cover in the drainage basin could also be the key factor leading to the relatively low level of CDOM in Pearl River estuary (Chen et al., 2004; Hong et al., 2005). In addition, extensive use of inorganic fertilizers in Chinese agriculture reduces one of the important sources of DOM in the soils and rivers that come from traditional manure and organic fertilizer (Chen et al., 2004).

On the other hand, the recharge types of the stream could also have a huge impact on the level of CDOM in the river. Compared to streams originated from marshland, peatland, and rainforest, streams originated from glacier and mountain area usually have a low degree of DOM and their recharges into the Yangtze river will dilute the concentration of DOM (Pan et al., 2012). The Yangtze River is originated from the Tanggula Mountains that are underlain by glaciers in large expanses. Dilution by stream recharged from glacial melting could also be an important factor in determining the low concentration of CDOM in the estuary. The other factor that results in the low degree of CDOM in Yangtze estuary could be attributed to the high turbidity in the Yangtze River and estuary. Compared to other major rivers in the world, the Yangtze River has a relatively high sediment load with an average annual discharge of 0.39 billions

of tons of sediments from 1951 to 2010, observed at Datong Station (Changjiang Sediment Bulletin of 2012, available in Chinese at <http://www.cjh.com.cn/fileupload/2013-09-13/130913171010455.pdf>). The high load of suspended matter prevents the light penetrating the water and limits the photosynthesis of phytoplankton in the upper layer. Therefore, the aquatic source of CDOM from photosynthesis could be reduced significantly. However, the relative importance of each of these driven factors could vary spatially and seasonal and cannot be quantified with further investigation.

Table 5-1. Comparisons of CDOM absorption  $a_g(355)$ ,  $S_{440}$  and salinity among different estuarine and coastal waters.

Region	$a_g(355) / m^{-1}$	$S_{440} / nm^{-1}$	Sal (PSU)	DOC ( $\mu mol/L$ )	Season	References
Amazon River and Estuary	0.14~3.12	0.014~0.033	0~36.4	NA	winter	(Green and Blough 1994)
Orinoco River plume	0.43~4.17	0.012~0.022	19.7~33.7	70~276	fall	(Del Castillo et al. 1999)
St. Lawrence River Estuary	0.2~3.8 $a_g(365)$	0.015	0~35.5	NA	summer	(Nieke et al. 1997)
Mackenzie River Estuary	0.6~2.02 $a_g(320)$	0.018~0.025	24.5~30.1	73.3~205.8	fall	(Retamal et al. 2007)
Scheldt estuary	0.97~4.22 $a_g(375)$	0.017~0.019	0.7~29.6	NA	winter	(Astoreca et al. 2009)
Chesapeake Bay	2.34~4.29 $a_g(375)$	0.018~0.019	2.7~19.9	NA	summer	
	0.4~4.6	0.015~0.024	0~33	60~290	spring	(Rochelle-Newall and Fisher 2002)
Southern Beaufort Sea	0.018~1.08 $a_g(440)$	0.015~0.023	0~35.0	49~460	summer	(Matsuoka et al. 2012a)
Pearl River estuary	0.24~1.93	0.014~0.018	0~32.5	NA	winter	(Hong et al. 2005)
Pearl River estuary	0.34~1.40		0~35.0	86~250	summer	(Chen et al. 2004)
Yangtze Estuary <sup>a</sup>	0.1~3.2	0.017~0.020	0~32.0		summer	(Guo et al. 2007)
	0.10~2.82	0.017~0.02	0.12~29.4	41.7~126.2	spring	(Sun et al. 2014b)
	0.43~3.74	0.018~0.02	0.2~26.5	139.4~220	summer	(Zhang et al. 2013)
	0.05~1.15	0.012~0.02	19~34.5	63.7~129.5	winter <sup>b</sup>	
Yangtze Estuary <sup>c</sup>	0.11~1.20	0.008~0.018	18.7~34.9		winter	This study
	0.23~1.91	0.012~0.025	4.0~33.6	60.5~217	summer	

a) Include samples collected from the river mouth;

b) Winter subset includes data collected from the PN section in March of 2011;

c) Only data from surface layer is counted.

### 5.4.2 Spatial-temporal variation of CDOM in the Yangtze estuary

The exhibited distribution of surface  $a_g(355)$  in the Yangtze estuarine and coastal waters is mainly because the Yangtze River runoff contributes most parts to the CDOM concentration in this area. The influence of runoff is weakening with the increasing distance away from the river mouth, resulting in a decline of terrestrial material along the downstream transit and forming the northwest-southeast decreasing gradient of CDOM. This kind of spatial distribution was in accordance with previous studies in this area (Lei et al., 2012; Sun et al., 2014b). The gradient of CDOM from the inner shelf to outer shelf was also observed in other river-influenced estuarine and coastal waters such as the St. Lawrence Estuary (Nieke et al., 1997), the Southern Bight of the North Sea (Warnock et al., 1999) and two estuaries of Canadian Arctic (Retamal et al., 2007). Note that the spatial distribution of CDOM in February 2012 (Figure 5.3b) was not a regular gradient. A more likely explanation related to this fact is that samples were collected during a tidal cycle. The tide type in the Hangzhou Bay and adjacent water is typical semidiurnal tide, and sampling interval in February cruise was less than 1 hour. Therefore, CDOM could probably be remobilized from the interstitial waters of bottom sediments and injected into overlying waters during the reverse phase of the tide.

The seasonal difference in CDOM between March 2012 (Figure 5.3c) and July 2012 (Figure 5.3d) in the Zhejiang coastal waters was mainly because of the seasonal changes in Yangtze runoff. Yangtze runoff was almost doubled from March to July in 2012 (See also the Changjiang Sediment Bulletin of 2012), resulting in a higher level of CDOM in summer. However, the seasonal variation in CDOM is not obvious for the two stations in the Hangzhou Bay (Figure 5.3c and Figure 5.3d), which is most likely due to the resuspension of bottom sediment resuspension. Sediment resuspension in the Hangzhou Bay in winter is mainly due to the distinct structure of the bay and the mutual effects by wind and tide (Chen et al., 2003). During the resuspension, bottom sediments with terrestrial matter injected into overlying waters and thus increase the CDOM concentration (Bodineau et al., 1998), resulting in the almost same level of CDOM at these two sites in March and July.

The vertical distribution of  $a_g(355)$  in July of 2011 (Figure 5.4) was more stratified in coastal water and peaked at the surface layer, which might be due to hydrological dynamics and the mixing behavior between fresh water and salt

water. For example, water density has been reported closely correlated to DOC in the vertical direction in the Yangtze estuary (Liu et al., 2014). Similar phenomena were also observed in previous studies (Del Castillo and Miller 2011; Gardner et al., 2005). The stratification was not significant in most offshore stations above the 10 m depth might because the runoff influence was weak, and water was well mixed.

The interannual variation of CDOM absorption coefficient is also observed by comparing our study with reported results by Guo et al. (2007) and Zhang et al. (2013). Note that the ranges of  $a_g(355)$  in Guo et al. (2007) and Zhang et al. (2013) as shown in Table 5-1 include samples collected from the river mouth with high absorption. In general,  $a_g(355)$  in summer from our study and these published results are comparable, and the interannual variation of CDOM in the Yangtze estuarine and coastal waters is rather small. Therefore, field measurements of CDOM properties in this study (*e.g.*, average  $a_g(355)$  and  $S_{400}$ ) could be useful in reducing the uncertainties of retrieval models on inherent optical properties (IOPs). This is mainly because CDOM contributes a large portion to the total absorption coefficient in most estuaries and is therefore sensitive in IOPs-retrieval models (Garver and Siegel, 1997; Lee et al., 2007; Lee et al., 2002; Maritorena et al., 2002; Yang et al., 2013). However, CDOM parameterization (Eq. (5.3)) in these models always requires empirical coefficient (*e.g.*,  $S_{440}$ ) and a fine-tuned value from field measurements would meet the demand.

#### **5.4.3 Implications: CDOM sources**

The negative correlation between salinity and  $a_g(355)$  (shown in Figure 5.5a) suggests that terrestrial inputs were the main source of CDOM in the study area, especially for the riverine end observations with low salinity samples (Figure 5.6b). The lower Pearson's  $r$  ( $= -0.572$ ) between salinity and  $a_g(355)$  in high salinity water ( $> 32$  PSU) shows that river influence got weakened in marine end observations, resulting in a decline of the proportion of terrigenous CDOM.

The same conclusion can be found by the relationship between salinity and  $S_{275-295}$  the dimensionless ratio  $S_R$ . The smaller  $S_{275-295}$  and  $S_R$  are corresponding to the higher MW, which is mostly determined by humic acid from the terrestrial source (Helms et al., 2008). Therefore, as shown in Figure 5.6, we can find that the DOM shifted from high MW to low MW with increasing salinity, corresponding to a decreasing proportion of terrigenous CDOM. The consistent increase of  $S_{275-295}$  along the salinity gradient could be attributed to the photobleaching of

terrigenous CDOM, which reduced the molecular size of DOM during downstream transit and thereby increased the  $S_{275-295}$  (Helms et al., 2008). Additionally, responses of CDOM absorption at 275nm and 295nm to the exposure to solar radiation are different, where a much greater fractional decrease in  $a_g(295)$  was found than  $a_g(275)$ , resulting in a larger  $S_{275-295}$  in upstream (Del Vecchio and Blough, 2002, 2004; Fichot and Benner, 2012). Another factor determined the high  $S_{275-295}$  in marine end waters could be attributed to the aquatic CDOM, especially the protein-rich plankton CDOM production, which is proved to have a significant effect on the  $S_{275-295}$  (Fichot and Benner, 2012).

Compared to the  $S_{275-295}$ , the  $S_R$  seems to be a more effective indicator for the proportion of terrigenous DOM in this study, since the negative relationships between the  $S_R$  and the salinity are more robust for the three types of waters (*i.e.*, riverine-dominated water, mixed water, and marine water) as shown in Figure 5.6b. This is mainly due to the values of spectral slope  $S_{350-400}$  generally decreased with increasing salinity (blue circles in Figure 5.8), which is consistent with (Helms et al., 2008). The increase of  $S_R$  is mainly due to the increase of  $S_{275-295}$ , while the decreasing trend of  $S_{350-400}$  enhanced the sensitivity of  $S_R$  to the terrigenous DOM.

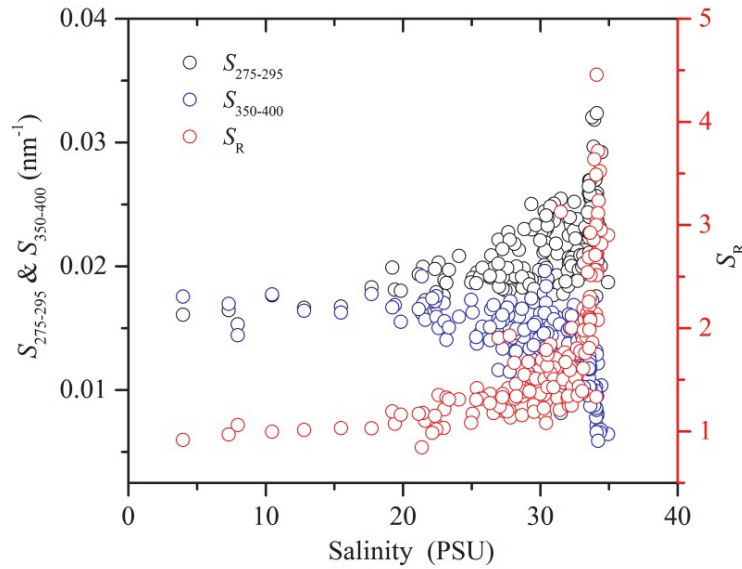


Figure 5.8. The trends of the spectral slopes ( $S_{275-295}$  and  $S_{350-400}$ ) and the slope ratio  $S_R$  to the increasing salinity.

#### **5.4.4 Implications: DOC source and estimation**

The range of DOC concentration in this study (*i.e.*, 60.54 to 217.06  $\mu\text{mol/L}$ ) is comparable with the investigations from Zhang et al. (2013) in the same study area, with reported DOC ranging from 76.4 to 264.4  $\mu\text{mol/L}$ . Note that results in Zhang et al. (2013) include several samples inside the river mouth that were rich in DOC concentration. On the other hand, DOC concentrations in this study agreed with CDOM absorption, are much lower than other major estuarine and coastal waters in the world. For instance, DOC in the Mississippi River estuary was as high as 446  $\mu\text{mol/L}$  (Chen and Gardner, 2004), while averaged at  $350 \pm 75$   $\mu\text{mol/L}$  in the Amazon River estuary (Ward et al., 2015), and even reached 800  $\mu\text{mol/L}$  in the Congo River estuary (Spencer et al., 2009b). It is most likely the same driven factors to the distribution of CDOM that result in the low concentration of DOC in the Yangtze estuary.

The mixture behavior between DOC and CDOM is non-conservative, which was also observed in the previous study in the Yangtze estuary (Cauwet and Mackenzie, 1993) and other river-dominated estuarine and coastal waters such as the Mississippi River estuary (Chen and Gardner, 2004), the Pearl River estuary (Chen et al., 2004) and Finnish estuaries (Asmala et al., 2012). However, the fact is that a conservative mixture behavior between DOC and CDOM was also reported in the Yangtze estuarine and coastal waters (Zhang et al., 2013), indicating the mixing behavior between the two parameters is highly variable in seasonal and regional scales. This could be potentially due to the asynchronous changes between DOC and CDOM (Vodacek et al., 1997), resulting in a variation of the CDOM portion in the DOM. For example, the discrepancy between DOC and CDOM in their stabilities and residence period is always significant in estuarine and coastal waters, as DOC are commonly unstable and easily decomposed compared to the CDOM in coastal waters (Spitzzy and Ittekkot, 1986).

As shown in Figure 5.7, the developed non-linear model presents a robust method to estimate DOC concentration from the  $S_{275-295}$  and  $a_g(355)$ . Note that despite the  $S_R$  is a more indicative parameter to terrestrial DOM in this study, it seems to be independent to DOC, which is consistent with results from (Helms et al., 2008). The developed model can be considered as satisfied, and the validation result is not biased since the adopted Cal/Val model (Salama et al., 2012) allows both calibration and validation datasets are independent to each other. It is worthy to note that the relationship between  $a_g^*(355)$  and  $S_{275-295}$  could work well for river-influenced coastal waters with the lignin as the important chromophore of CDOM



pool and the main terrigenous component of DOC (Fichot and Benner, 2012), which could probably be the reason for the good performance of this non-linear relationship in the Yangtze estuary.

However, due to the nature of this non-linear relationship is empirical, applying this relationship to different regions could require a fine-tuning on the empirical coefficients. For instance, the coefficients developed by Fichot and Benner (2012) will overestimate DOC in the Yangtze estuarine and coastal waters as shown in Figure 5.7a. On the other hand, as shown in this study, the developed relationship worked well for observations collected from both July 2011 and August 2013, indicating that the fine-tuned empirical coefficients could potentially work well for the coefficients-developed region in different temporal scales. Moreover, with the regionally applicable DOC-CDOM relationship, DOC can be potentially estimated from earth observation data since CDOM absorption coefficient can be retrieved from remote sensing reflectance ( $R_{rs}$ ) with acceptable accuracy (Lee et al., 2002; Matsuoka et al., 2012b; Siswanto et al., 2011).

## 5.5 Conclusions

Investigation in this study shows that the CDOM properties in the Yangtze estuarine and coastal waters varied significantly at spatial and temporal scales. CDOM absorption coefficient,  $a_g(355)$ , at the surface layer, generally decreased with the increasing distance from the river mouth and was overall much higher in summer (July 2012) than in winter (March 2012). Vertical variation of  $a_g(355)$  was more profound in near-shore waters than that in offshore waters. The analysis in this study indicates that terrestrial input was the main source of CDOM in the investigated region, but the proportion of terrigenous CDOM declined with increasing salinity, which was probably due to photobleaching during downstream transit. Sediment resuspension could also be an important source of CDOM in the Hangzhou Bay in winter. The aquatic source could contribute to a large proportion of CDOM in marine waters with high salinities.

The slope ratio  $S_R$  was a more indicative parameter to terrigenous DOM compared to the  $S_{275-295}$  in the Yangtze estuarine and coastal waters, but it was independent to DOC. A non-linear empirical relationship between  $a_g^*(355)$  and  $S_{275-295}$  was developed to estimate DOC from CDOM optical properties, which was proved robust in the Yangtze estuarine and coastal waters. This relationship provides an alternative approach to estimate DOC in large spatial scale from

remotely retrieved CDOM optical properties over the Yangtze estuarine and coastal water.

## **Chapter 6    Synthesis**

The purpose of this chapter is to synthesize the key findings resulting from this study and to put them into a broader perspective. This chapter begins with the research findings from Chapter 3, Chapter 4, and Chapter 5. It is followed by putting the scientific work of this thesis in a broader context and closing with an outlook on future efforts in ocean color remote sensing in turbid coastal waters.

## **6.1 Research Findings**

To summarize, in this study, the author employed the analytical 2SeaColor model, with his contribution of improved parameterizations, to derive the inherent optical properties and diffuse attenuation coefficient in the optically-complex Yangtze estuary. The derived particulate backscattering coefficient was later adopted to empirically estimate the concentration of SPM, in which a novel model was proposed to retrieve wide-range  $C_{SPM}$  without the need for a switching scheme. The last effort the author made was to investigate the seasonal variability of CDOM absorption properties and develop a locally calibrated model to estimate the concentration of DOC. The key findings resulting from the work mentioned above are discussed in the following text.

- 1) For the retrieval of IOPs and  $K_d$  using the 2SeaColor model as discussed in Chapter 2, the assumption of negligible absorption by non-water in NIR band does not hold for optically complex waters, such as the Yangtze estuary. Figure 3.4 shows that the 2SeaColor model with improved parameterization improves the retrieval accuracy compared with the original 2SeaColor model, in which model the assumption was adopted.
- 2) The 2SeaColor model, with improved parameterizations, should be recommended for  $K_d$  retrieval in optically complex waters instead of the empirical and semi-analytical models, which is mainly attributed to the fact that the other two types of models employ empirical coefficients that require adjustments or fine-tuning for different types of water in order to achieve the minimum retrieval error. To the contrary, the 2SeaColor model, without the need of fine-tuning of empirical constants, predicts the optimum solution by spectral optimization. However, for optically-simple waters, such as Open Ocean or offshore water, a regionally calibrated empirical model could be more suitable, with fast data process and comparable, if not better, performance with the semi-analytical and analytical models (Table 3-2).
- 3) In Chapter 4, the author shows that the relationship between  $C_{SPM}$  and  $b_{bp}$  changed from linear in clear waters to significantly curved in extremely turbid waters, which indicates that the retrieval models that assume constant

specific backscattering coefficient would face larger errors when applied to turbid waters. The change in the relationship also points out that using a single function to calibrate the relationship for a wide-range  $C_{SPM}$  could result in significant uncertainties while using two respective models will introduce discontinuity issues. The same situation should also apply to  $R_{rs}$ -based SPM retrieval models. The sigmoid function used in this study could provide a hint for the future development of wide-range SPM retrieval models.

- 4) Also, in Chapter 4, it is clearly observed that the GOCI standard AC product significantly underestimates  $R_{rs}$  at the two NIR bands (Figure 4.4). Even with the empirical correction (Eq. (4.5) and Eq. (4.6)), the  $R_{rs}$ (NIR) could be still underestimated. However, the imperfect AC provides important information regarding the SPM retrieval models, which is that a multi-band algorithm could compensate for the uncertainties caused by  $R_{rs}$ (NIR) if the visible bands are relatively accurate. The He model employs  $R_{rs}$  only at 745 nm and 490 nm, and the mapping product shows significant underestimation compared with the GOCI model (uses three visible bands plus 745 nm) and the 2SeaColor model (uses all the 8 bands).
- 5) The  $b_{bp}$  index using the sigmoid function seems to be robust when applied to other coastal waters, while the  $b_{bp}$ - $C_{SPM}$  relationships show relatively large variations. However, the physics or mathematics behind the model's robustness in less turbid waters requires further investigation. Nevertheless, the validation against the independent SeaSWIR dataset shows that the proposed  $S_{index}$ - $C_{SPM}$  model is applicable to other turbid coastal waters.
- 6) The  $b_{bp}$  derived from the 2SeaColor model was not validated in the extremely turbid waters (concentration > 100 mg/L). However, its robust correlation with  $C_{SPM}$  indicates that by using the spectral optimization technique, the resulting  $b_{bp}$  contains the integral information from all wavebands. The error in the derived  $b_{bp}$  from the "ground-truth" value, if there is any, does not matter that much in terms of retrieving the  $C_{SPM}$ . On the other hand, the semi-analytical retrieval model using a single band to derive  $b_{bp}$  was not suitable for developing the  $b_{bp}$ - $C_{SPM}$  relationship in extremely turbid waters.
- 7) In Chapter 5, the comparisons between CDOM absorption properties, especially the spectral slope of  $a_g(\lambda)$ , in global coastal waters demonstrate that spatial-temporal variation of  $S$  is significant. Therefore, for IOPs retrieval models, the author would not recommend using the constant  $S$  for model development.
- 8) The relationships between CDOM and salinity show very limited variation in winter and summer in terms of the slope from the regression analysis (Figure

5.5), which indicates that the CDOM could be potentially used to estimate salinity in coastal environments where the salinity satellite (such as the SMOS) has large uncertainties. The relationship between DOC and CDOM demonstrates the potential of remotely retrieving DOC from satellite imagery, but it does involve special requests that UV bands should be included in order to implement the Eq. (5.4).

## **6.2 Sustainable Development Goals**

This thesis contributes to our understanding of SPM and CDOM variations, in diurnal and seasonal scales, respectively, in the Yangtze estuary as related to their sources and the environmental conditions. As SPM and CDOM, reduce light transparency, contain particulate and dissolved carbon, and indicate land-based pollution, the scientific work presented here could be a direct contribution to the United Nations' Sustainable Development Goals (SDGs).

In particular, the SDGs Goal 6 indicators 6.3 and 6.6 require protecting and restoring water-related ecosystems, including the estuaries, and improving water quality globally by 2020 and 2030, respectively. The SDGs Goal 14 indicators 14.1 and 14.2, on the other hand, demand sustainably managing and protecting the coastal ecosystem to avoid significant adverse impacts and preventing and significantly reducing marine pollution of all kinds, especially from land-based activities, by 2020 and 2025, respectively. Overall, problems that threaten the sustainable development of the aquatic ecosystem, to some extent, are all related to human activities. To better understand how the ecosystem responds associated with human activities and provide expert advice to society for ecosystem governance, the scientist must rest on a solid scientific basis and long-term time series observations.

The scientific work of this dissertation shows that remote sensing technique could provide a synoptic view of the variations of water constituents, which are closely associated with the biogeochemical substances in the water. Thereby, the variations of water constituents, if correctly interpreted, can serve as proper indexes of the health status of the ecosystem regarding the biological productivity, nutrient pollution, and hazardous chemicals and materials. Current operational ocean color satellites, along with their predecessor CZCS since 1978, provide time series and near real-time images that can be employed to analysis, monitor and predict the water quality at large temporal and spatial scales for the aquatic ecosystem. Interpretation of the long-term time series satellite data will provide

a scientific basis for governments and social communities to take action to achieve a healthy and productive ecosystem.

### 6.3 Outlook

Research of ocean color remote sensing in turbid coastal waters is no doubt still full of great challenges for many years ahead. Dedicated efforts and innovative thoughts are required to promote the research and overcome the challenges and problems. Based on the research the author has done in this dissertation, the author thinks that there is still plenty of space for development and there are more topics raising from this study that could be worth studying in the near future.

- 1) The GOCI images, with great spatial and temporal resolutions, could serve as a tool in monitoring the optical properties in the Yangtze estuary if the AC can be improved for NIR bands for turbid waters.
- 2) The sigmoid function could be incorporated directly with  $R_{rs}$  to see if a continuous and accurate  $C_{SPM}$  can be derived for wide-range  $C_{SPM}$  data.
- 3) All the semi-analytical models and the 2SeaColor model do not consider the impact of the particle phase function on the accuracy of retrieved IOPs, especially the  $b_b$ . Further efforts could at least investigate the impact of the phase function on retrieved IOPs using simulated data generated by, for example, the Hydrolight® model.
- 4) The similar absorption spectral features between CDOM and NAP make it difficult to distinguish the two in the semi-analytical model. Partitioning the combined absorption by CDOM and NAP to their individual absorptions requires empirical relationships, which is less reported for the Yangtze estuary.
- 5) With the proposed model in this study, it is possible to investigate the monthly and annually variation of IOPs,  $K_d$  and SPM. It would be of interest to understand how the estuary ecosystem has changed in the last decades and how sensitive it responds to extreme climate events and global climate change.





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## *Bibliography*

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## Biography



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### Peer-reviewed journal papers:

**Yu, X.**, Salama, M.S., Shen, F., & Verhoef, W. (2016). Retrieval of the diffuse attenuation coefficient from GOCI images using the 2SeaColor model: A case study in the Yangtze Estuary. *Remote Sensing of Environment*, 175, 109-119

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